

BOUNDARIES WITHOUT BOUNDARIES

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Work with:

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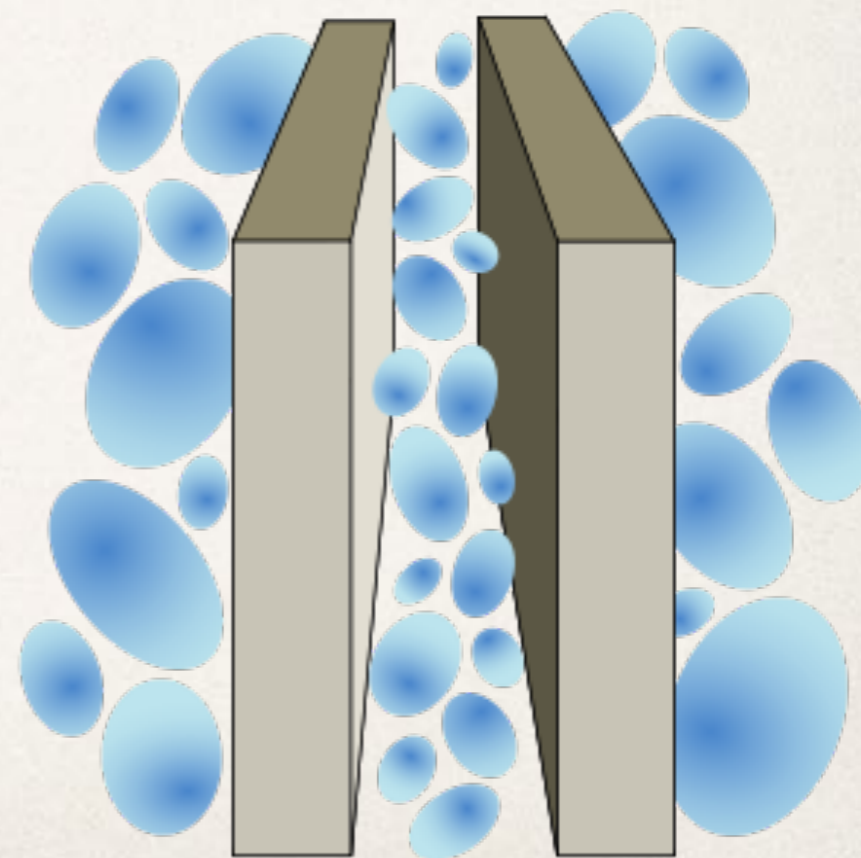
Fields, Particles and Information Geometry

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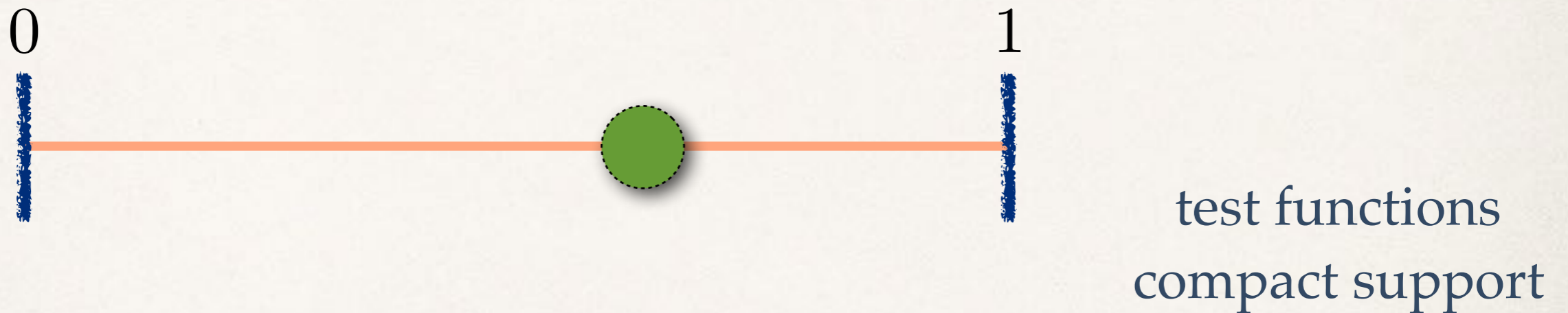
DIAS, Dublin, 24 January 2018

Outline

- Quantum boundary conditions
- Boundaries by reduction: from a circle to an interval
- A quantum particle on a circle: Laplacian
- Unitary folding
- Momentum on a half-line
- Momentum on an interval



A quantum particle on a segment



$$H = T = \frac{p^2}{2m} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \quad D(T) = \mathcal{D}(0, 1)$$

What happens at the boundary?

$$D(T^*) = \{\psi \in L^2(0, 1), \psi'' \in L^2(0, 1)\} = H^2[0, 1]$$

Quantum Boundary Conditions

$$T_U \psi = -\frac{\hbar^2}{2m} \psi'' \quad U \in U(2)$$

$$i(I + U) \varphi = (I - U) \dot{\varphi}$$

$$\varphi = \begin{pmatrix} \psi(0) \\ \psi(1) \end{pmatrix} \quad \dot{\varphi} = \begin{pmatrix} \psi'(0) \\ -\psi'(1) \end{pmatrix}$$

Quantum Boundary Conditions

Self-adjoint extensions



Unitary matrices at the boundary

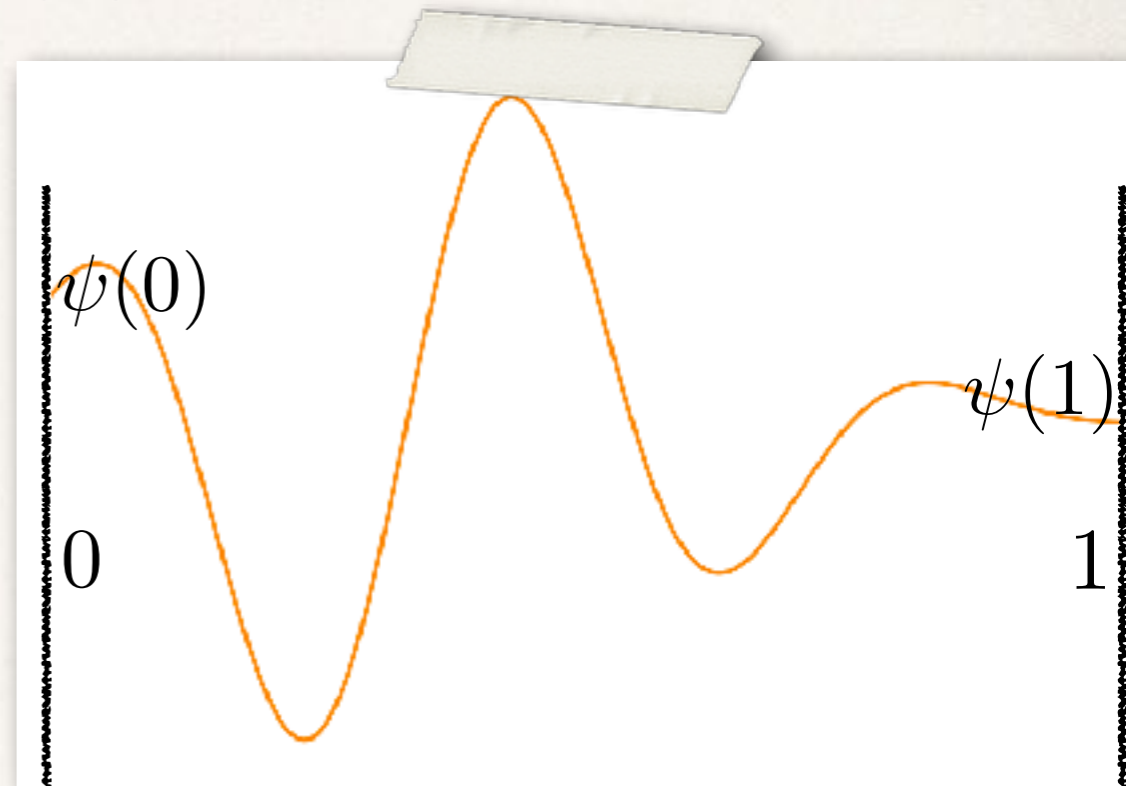
$$i(I + U)\varphi = (I - U)\dot{\varphi} \quad U \in U(2)$$

Examples

$$i(I + U) \begin{pmatrix} \psi(0) \\ \psi(1) \end{pmatrix} = (I - U) \begin{pmatrix} \psi'(0) \\ -\psi'(1) \end{pmatrix}$$

Dirichlet

$$U = I \quad \psi(0) = 0 = \psi(1)$$



Robin

$$U = -e^{-i\alpha} I$$

$$\alpha \in \mathbb{R} \quad \begin{cases} \psi'(0) = -\tan \frac{\alpha}{2} \psi(0), \\ \psi'(1) = \tan \frac{\alpha}{2} \psi(1) \end{cases}$$

Neumann

$$U = -I$$

$$\psi'(0) = 0 = \psi'(1)$$

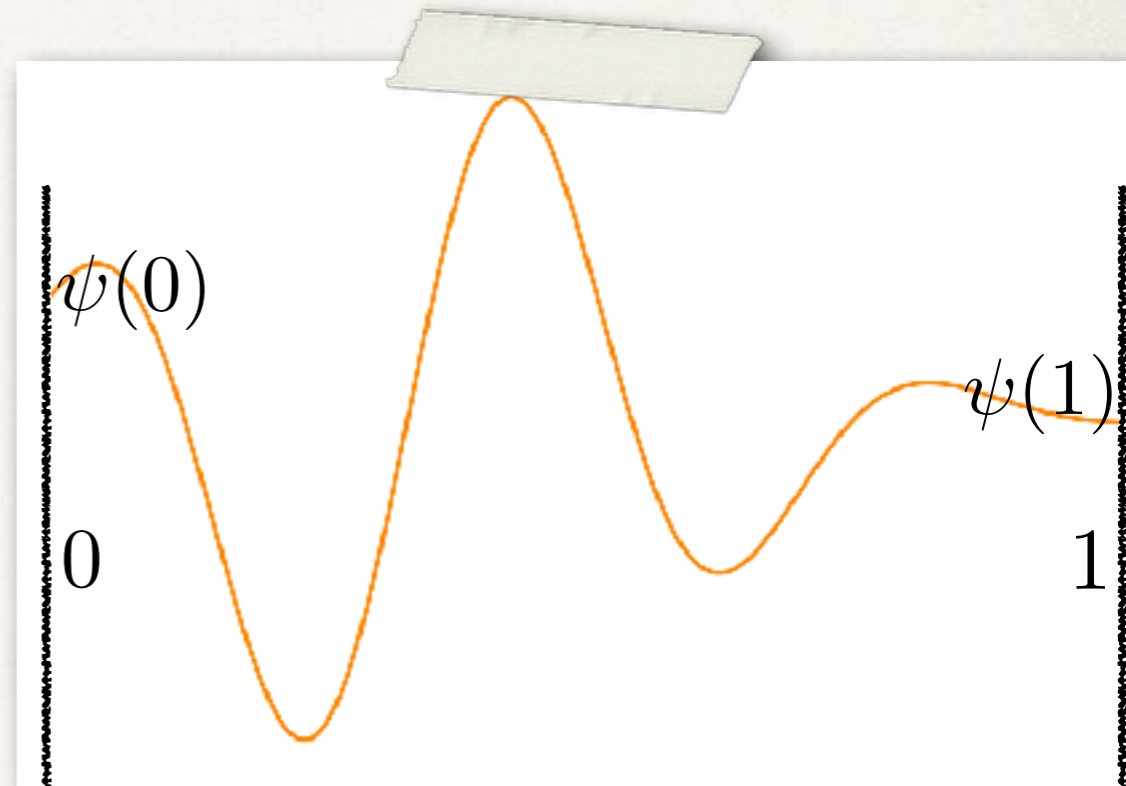
Examples

$$i(I + U) \begin{pmatrix} \psi(0) \\ \psi(1) \end{pmatrix} = (I - U) \begin{pmatrix} \psi'(0) \\ -\psi'(1) \end{pmatrix}$$

Dirichlet+Robin

$$U = \begin{pmatrix} 1 & 0 \\ 0 & -e^{-i\alpha} \end{pmatrix} \begin{cases} \psi(0) = 0 \\ \psi'(1) = \tan \frac{\alpha}{2} \psi(1) \end{cases}$$

$\alpha \in \mathbb{R}$



Antiperiodic

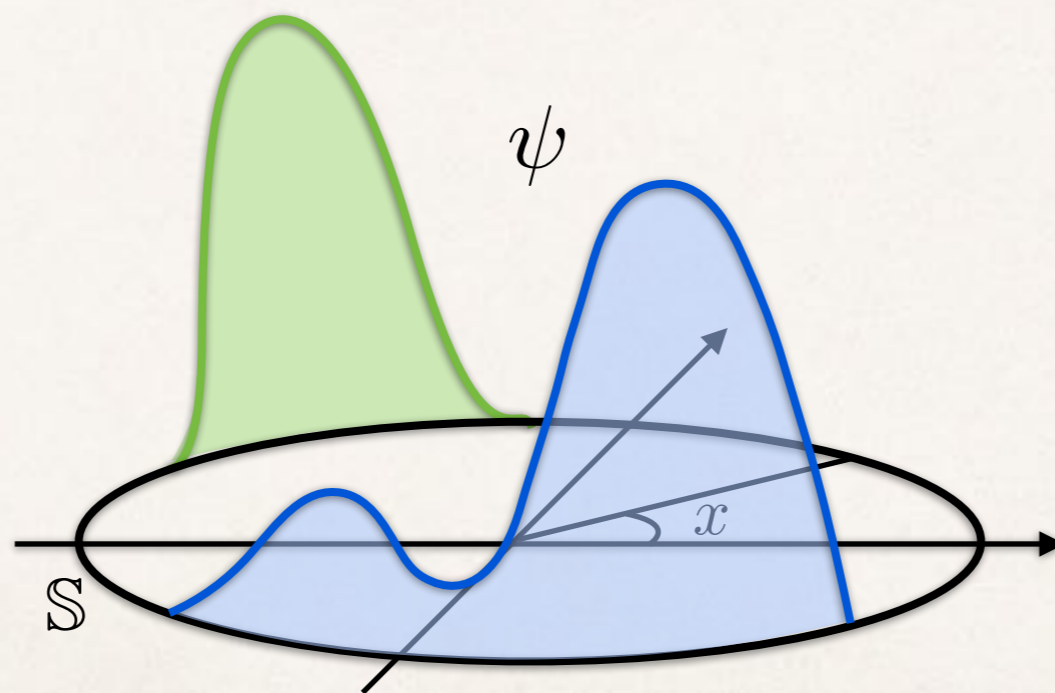
$$U = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{cases} \psi(0) = -\psi(1) \\ \psi'(0) = -\psi'(1) \end{cases}$$

Periodic

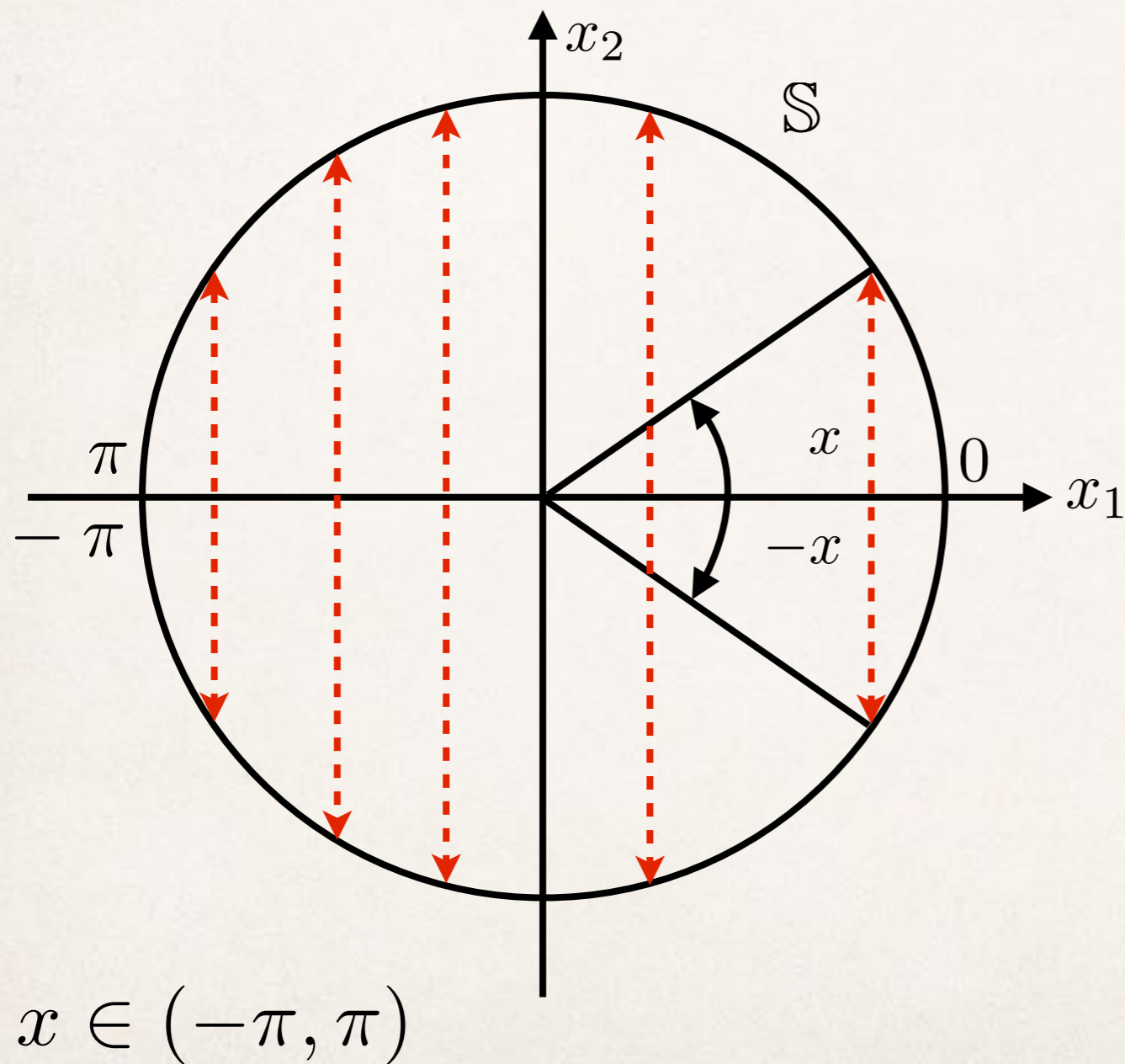
$$U = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix} \begin{cases} \psi(0) = \psi(1) \\ \psi'(0) = \psi'(1) \end{cases}$$

Boundaries without boundaries

Generate **quantum boundary conditions** starting from a particle on a manifold **without boundaries**?



Reduction by symmetry



$$\Pi : \mathbb{S} \rightarrow \mathbb{S},$$

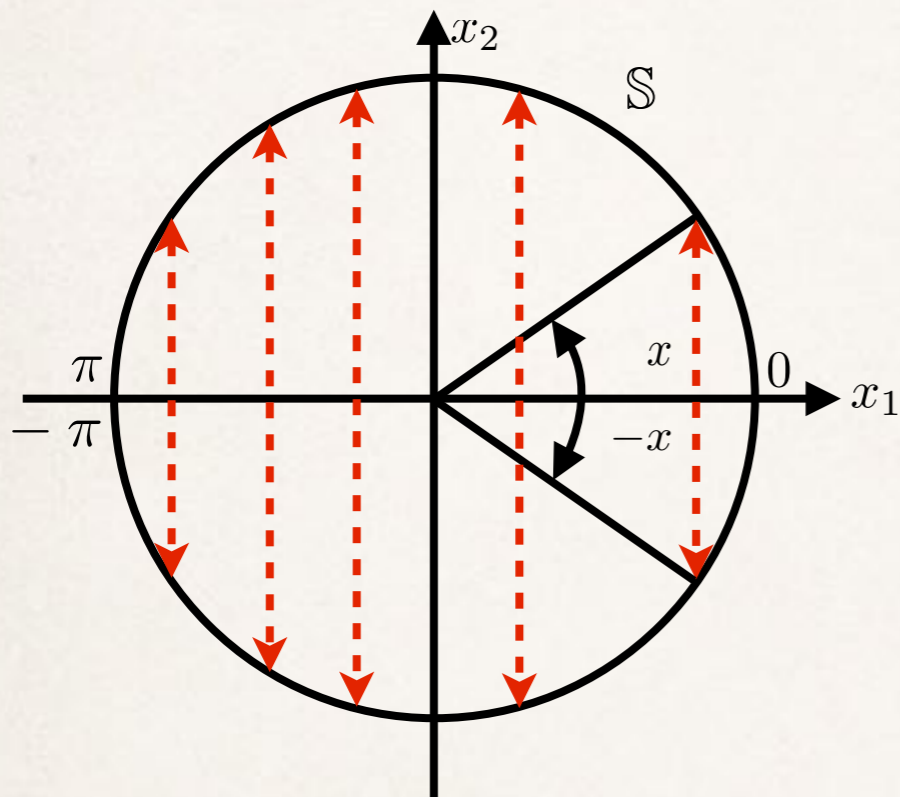
$$\Pi(x_1, x_2) = (x_1, -x_2)$$

$$\Pi(x) = -x$$

$$\mathbb{S}_+ = \mathbb{S}/\Pi = [0, \pi]$$

Manifold with boundary

QBC by reduction



$$\Pi(x) = -x$$

$$P : L^2(\mathbb{S}) \rightarrow L^2(\mathbb{S})$$

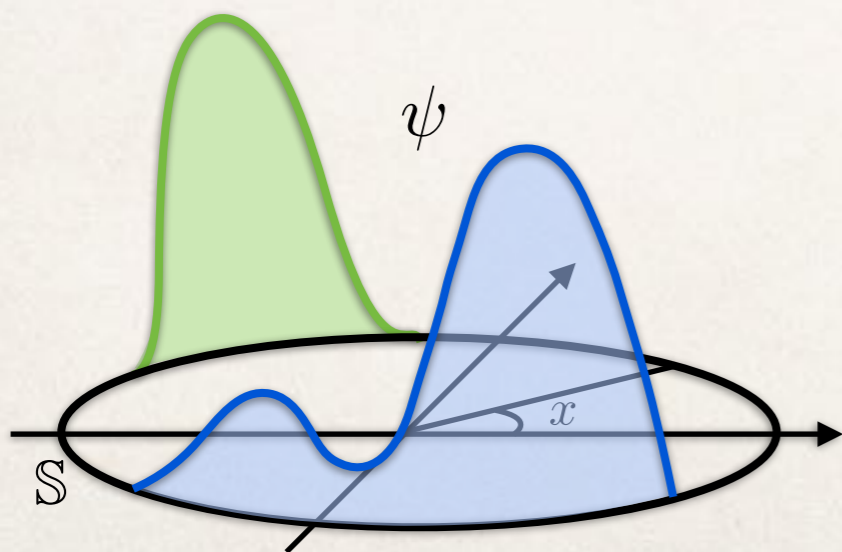
$$(P\psi)(x) = \psi(\Pi(x)) = \psi(-x)$$

parity operator

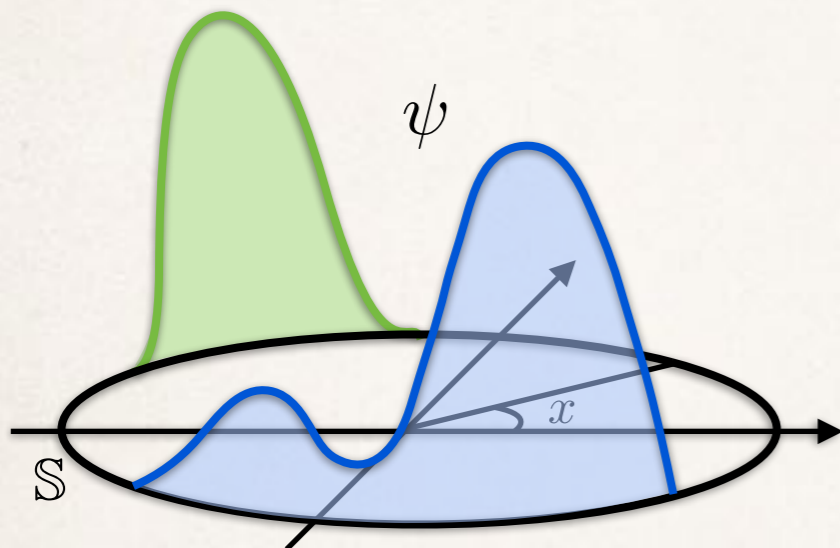
$$L^2(\mathbb{S}) = \mathcal{H}_+ \oplus \mathcal{H}_-$$

$$\mathcal{H}_\pm = \{\psi \in L^2(-\pi, \pi) \mid \psi(-x) = \pm\psi(x)\}$$

even/odd functions



QBC by reduction



Free particle Hamiltonian

$$H = \frac{p^2}{2m} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$$

domain of self-adjointness

$$H^2(\mathbb{S}) = \{\psi \in H^2[-\pi, \pi] : \psi(-\pi) = \psi(\pi), \psi'(-\pi) = \psi'(\pi)\}$$

second Sobolev space

$$H P = P H$$

commutes with parity

$$L^2(\mathbb{S}) = \mathcal{H}_+ \oplus \mathcal{H}_-$$

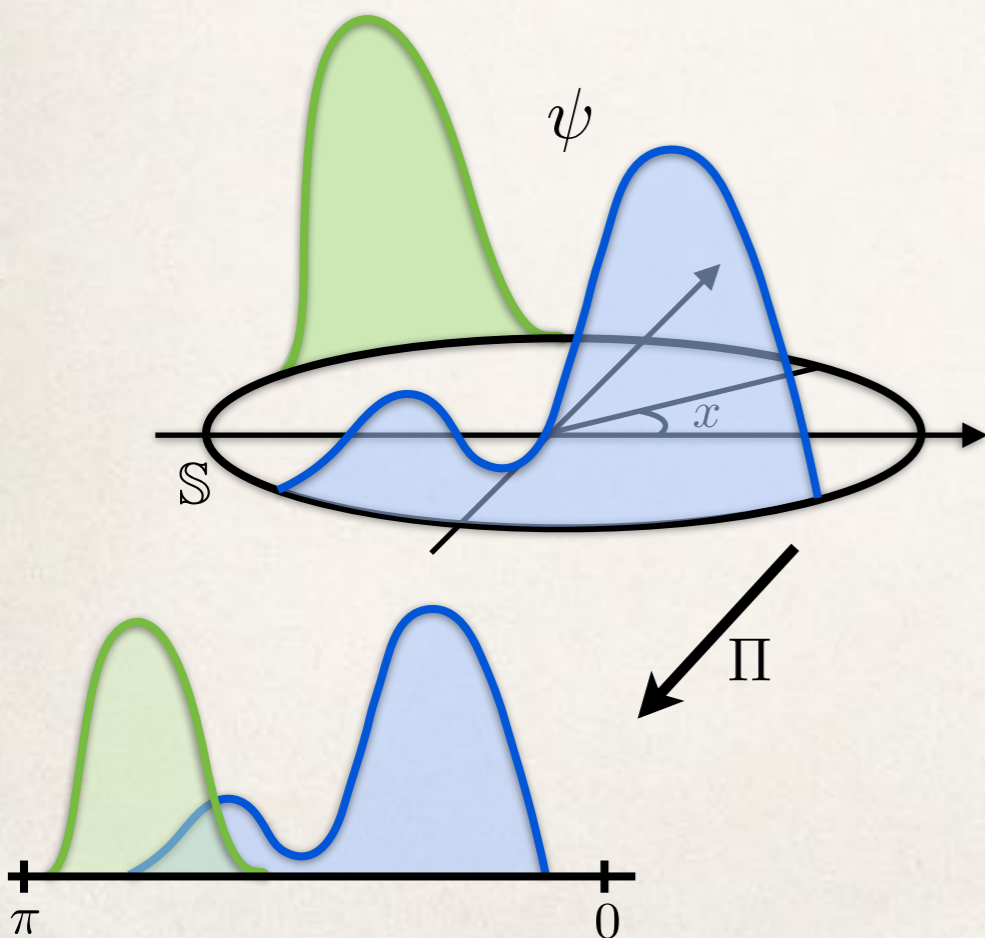
invariant subspaces

QBC by reduction

$$H^2(\mathbb{S}) = \{\psi \in H^2[-\pi, \pi] : \psi(-\pi) = \psi(\pi), \psi'(-\pi) = \psi'(\pi)\} \quad \text{circle}$$

restriction to even functions:

$$D(H|_{\mathcal{H}_+}) = H^2(\mathbb{S}) \cap \mathcal{H}_+ = \{\psi \in H^2[-\pi, \pi] \cap \mathcal{H}_+ : \psi'(-\pi) = 0 = \psi'(\pi)\}$$



$$U_+ : \mathcal{H}_+ \rightarrow L^2(0, \pi) \quad \text{interval}$$

$$\phi(x) = (U_+ \psi)(x) = \sqrt{2} \psi(x) \quad x \in [0, \pi]$$

$$H_+ := U_+ H|_{\mathcal{H}_+} U_+^\dagger = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \quad \text{on } L^2(0, \pi)$$

$$D(H_+) = \{\psi \in H^2[0, \pi] : \psi'(0) = 0 = \psi'(\pi)\}$$

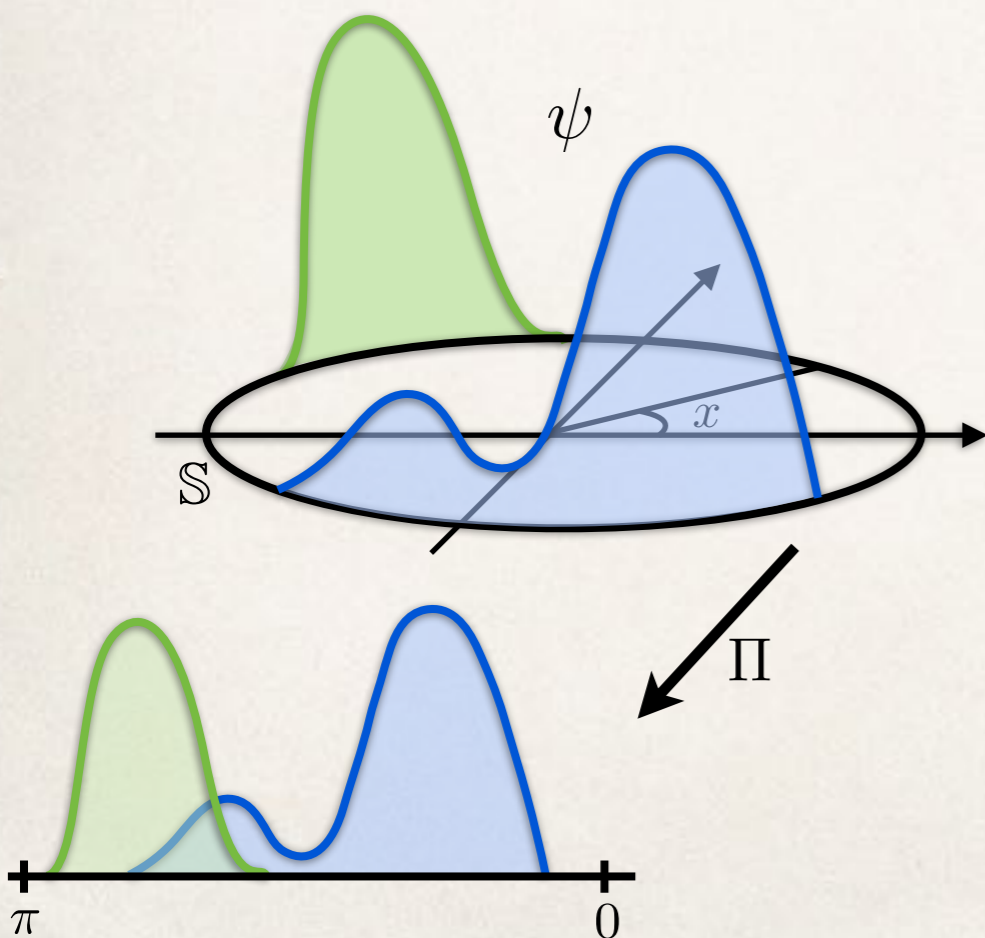
Neumann

QBC by reduction

$$H^2(\mathbb{S}) = \{\psi \in H^2[-\pi, \pi] : \psi(-\pi) = \psi(\pi), \psi'(-\pi) = \psi'(\pi)\} \quad \text{circle}$$

restriction to odd functions:

$$D(H|_{\mathcal{H}_-}) = H^2(\mathbb{S}) \cap \mathcal{H}_- = \{\psi \in H^2[-\pi, \pi] \cap \mathcal{H}_- : \psi(-\pi) = 0 = \psi(\pi)\}$$



$$U_- : \mathcal{H}_- \rightarrow L^2(0, \pi) \quad \text{interval}$$

$$\phi(x) = (U_- \psi)(x) = \sqrt{2} \psi(x) \quad x \in [0, \pi]$$

$$H_- = U_- H|_{\mathcal{H}_-} U_-^\dagger = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \quad \text{on } L^2(0, \pi)$$

$$D(H_-) = \{\psi \in H^2[0, \pi] : \psi(0) = 0 = \psi(\pi)\}$$

Dirichlet

QBC by reduction

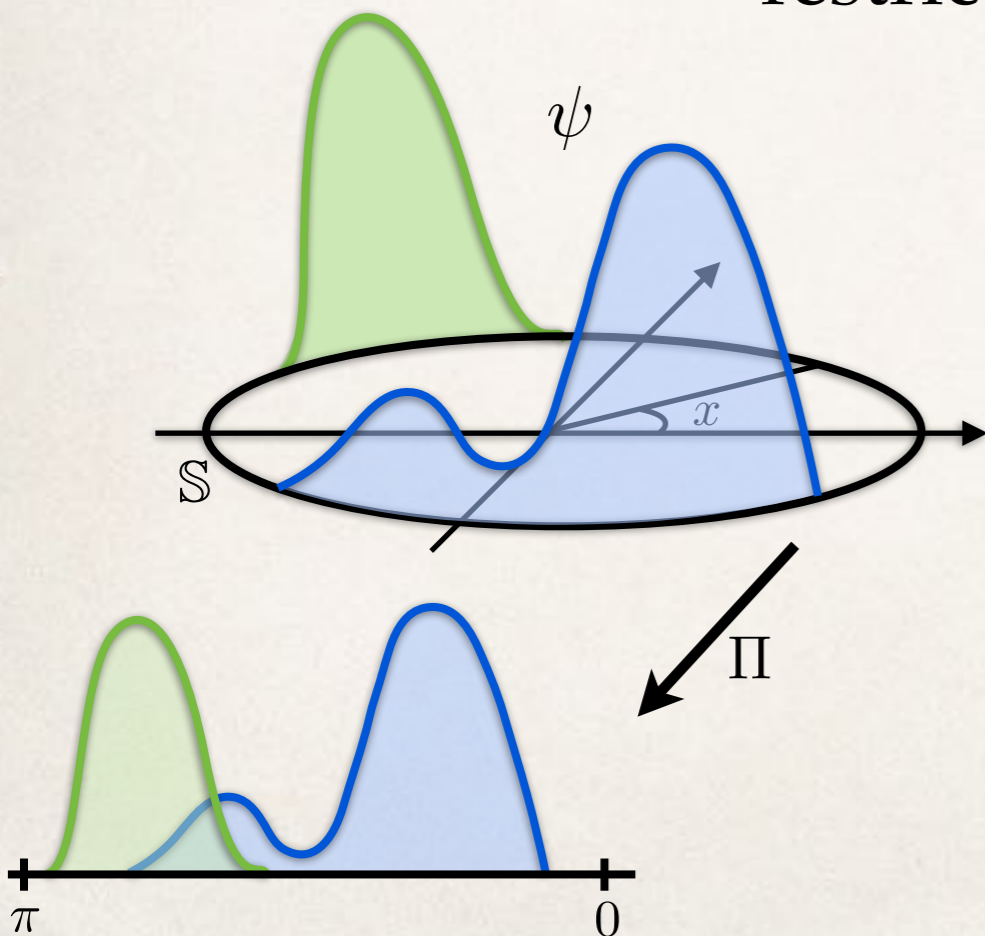
$$H = \frac{p^2}{2m} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$$

restriction to **even** functions

Neumann

restriction to **odd** functions

Dirichlet



Other boundary conditions?

General QBC

spatial metric

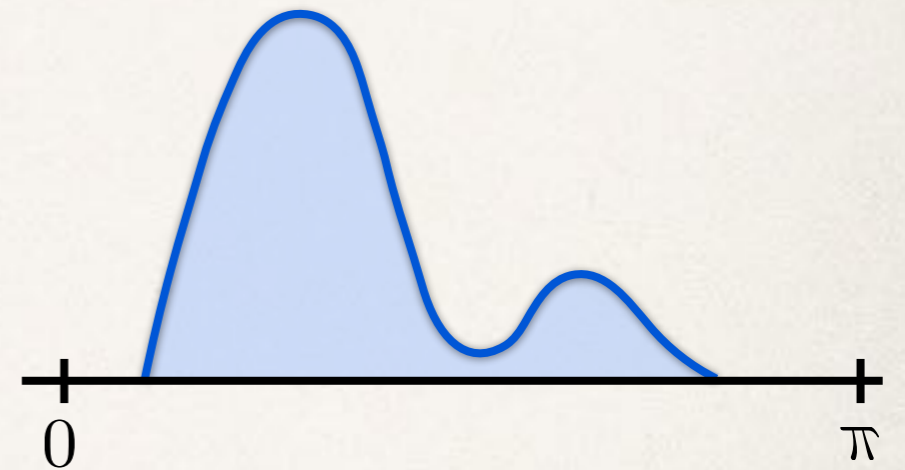
$$ds^2 = dx^2$$

Hamiltonian

$$H = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$$

Neumann bc

$$\psi'(0) = 0 = \psi'(\pi)$$



Change of coordinates on $[0, \pi]$

$$y = F(x) = \int_0^x f(t) dt \quad f \geq 0 \quad \int_0^\pi f(t) dt = \pi$$

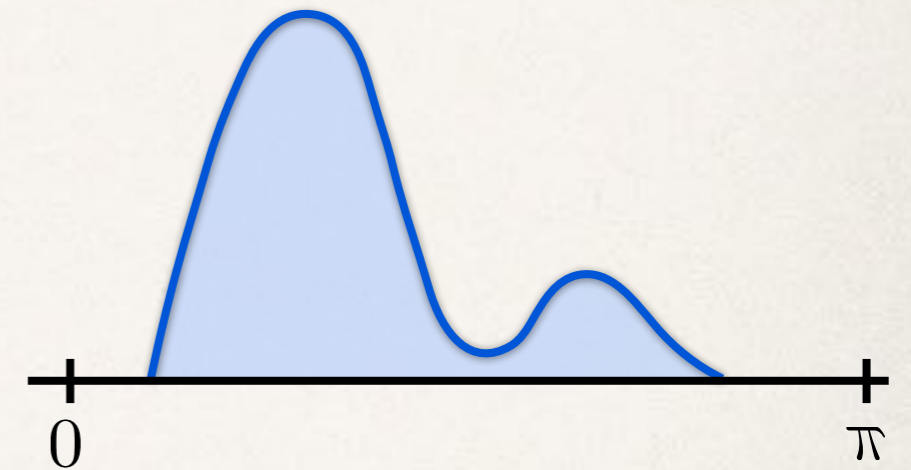
General QBC

spatial metric

$$ds^2 = \left(\frac{dx}{dy} \right)^2 dy^2 = \frac{1}{[f(y)]^2} dy^2$$

unitary transformation

$$\phi(y) = (U_f \psi)(y) = \frac{1}{\sqrt{g(y)}} \psi(F^{-1}(y)), \quad g(y) = f(F^{-1}(y))$$



Transformed Hamiltonian

$$H_f = U_f H U_f^\dagger = -g^2 \frac{\hbar^2}{2m} \frac{d^2}{dy^2} - gg' \frac{\hbar^2}{m} \frac{d}{dy} + V, \quad V = \frac{\hbar^2}{8m} [(g')^2 + 2gg'']$$

General QBC

boundary conditions

$$\begin{cases} \phi'(0) = \nu_0 \phi(0), \\ \phi'(\pi) = -\nu_\pi \phi(\pi), \end{cases}$$

$$\nu_0 = -\frac{1}{2} \frac{f'(0)}{[f(0)]^2},$$

$$\nu_\pi = \frac{1}{2} \frac{f'(\pi)}{[f(\pi)]^2}$$

Robin!

But transformed Hamiltonian \neq kinetic energy

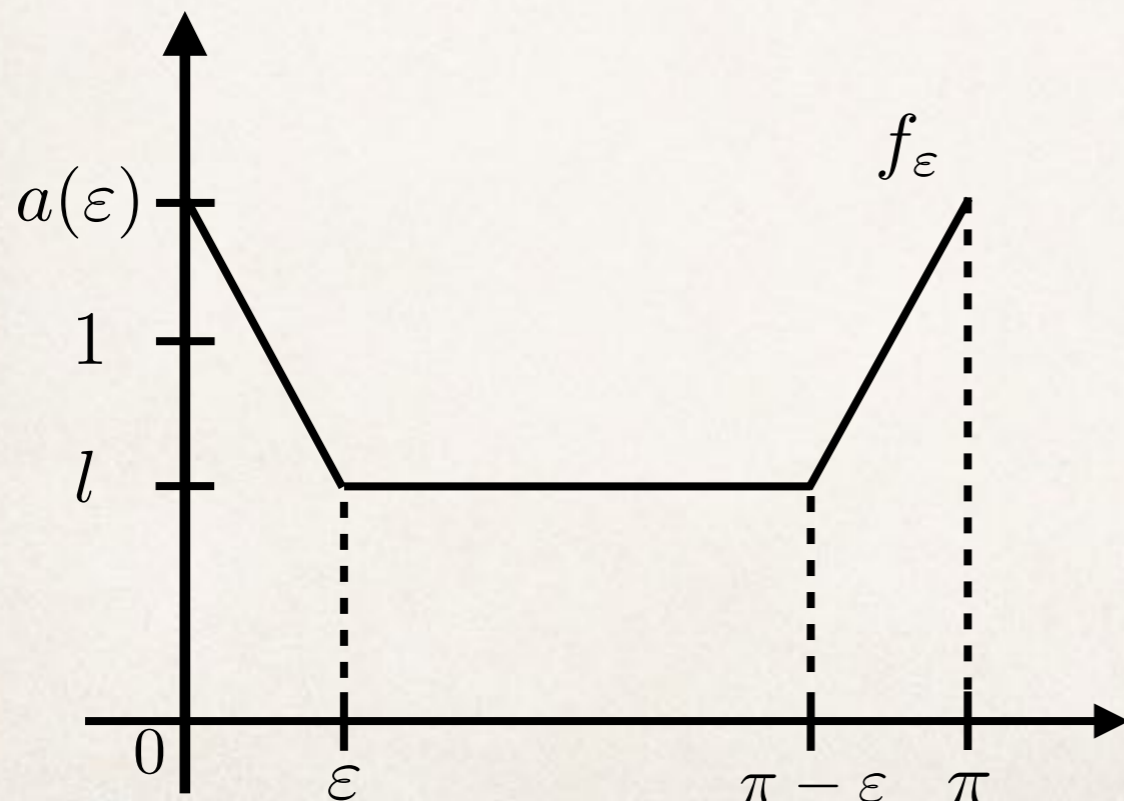
Change the metric only in a boundary layer!

From Neumann to Robin

$$H_f = U_f H U_f^\dagger = -g^2 \frac{\hbar^2}{2m} \frac{d^2}{dy^2} - gg' \frac{\hbar^2}{m} \frac{d}{dy} + V,$$

$$V = \frac{\hbar^2}{8m} [(g')^2 + 2gg'']$$

$$\nu_0 = -\frac{1}{2} \frac{f'(0)}{[f(0)]^2}, \quad \nu_\pi = \frac{1}{2} \frac{f'(\pi)}{[f(\pi)]^2}$$



sequence of functions

$$f_\epsilon(x) \rightarrow k \quad x \in (0, \pi)$$

$$H_{f_\epsilon} \rightarrow -\frac{\hbar^2}{2m} k^2 \frac{d^2}{dy^2} = -\frac{\hbar^2}{2M} \frac{d^2}{dy^2}$$

$$\lim_{\epsilon \rightarrow 0} \nu_0 = \mu_0 > 0,$$

$$\lim_{\epsilon \rightarrow 0} \nu_\pi = \mu_\pi > 0.$$

Symmetry?

Crucial ingredient for reduction:

$$H P = P H$$

Hamiltonian commutes with parity

$$L^2(\mathbb{S}) = \mathcal{H}_+ \oplus \mathcal{H}_-$$

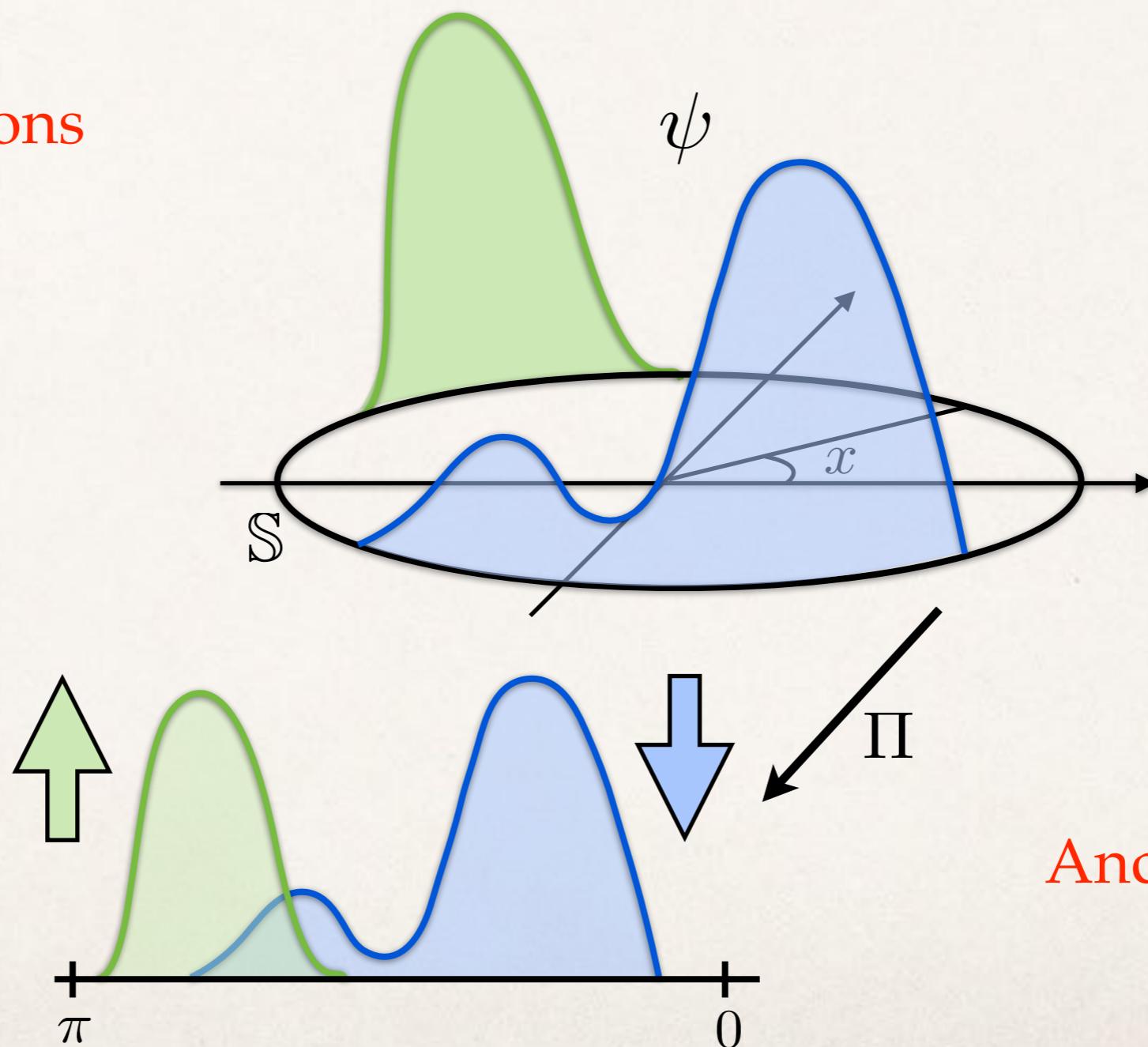
invariant subspaces

What if $[H, P] \neq 0$ and there are no invariant subspaces?

Replace projections by unitaries: **folding**

Folding

Replace projections
by unitaries



Ancillary spin

Momentum on the half-line

$$p = -i\hbar \frac{d}{dx} \quad \text{on } L^2(\mathbb{R}_+)$$

it admits **no** self-adjoint extensions!

Momentum on the half-line

Consider the momentum on the **full line**

$$p = -i\hbar \frac{d}{dx}$$

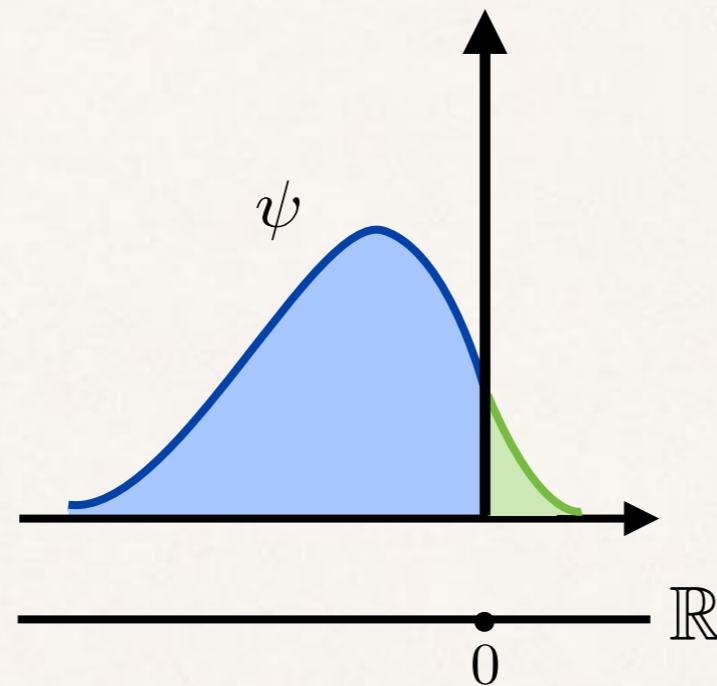
$D(p) = H^1(\mathbb{R}) = \{ \psi \in L^2(\mathbb{R}) \mid \psi' \in L^2(\mathbb{R}) \}$ first Sobolev space

It is self-adjoint.

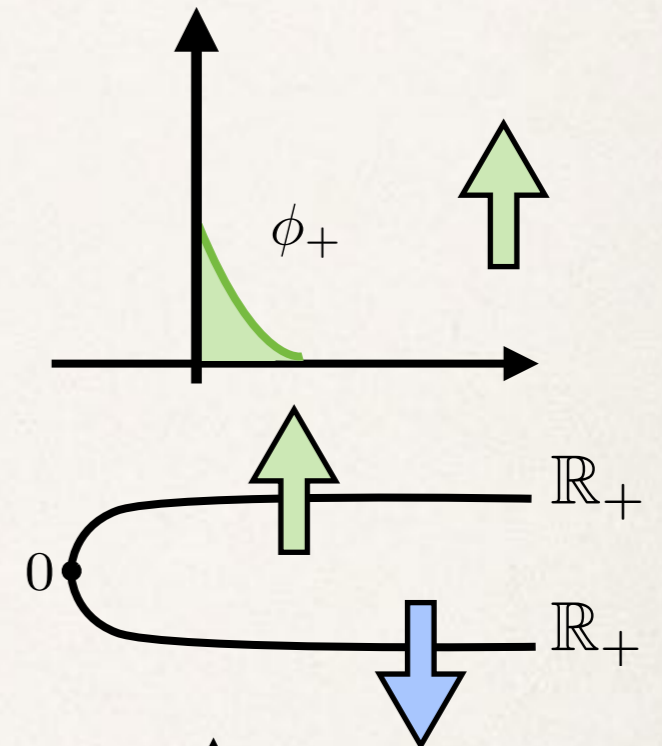
Need a **unitary** mapping onto the half-line

Folding

Ancillary spin

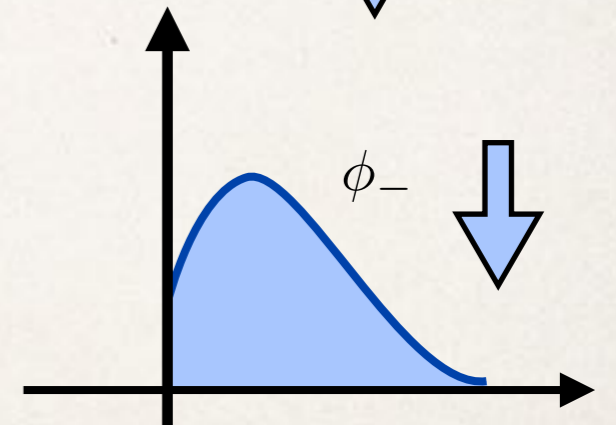


U



Unitary:

$$U : L^2(\mathbb{R}) \rightarrow L^2(\mathbb{R}_+) \otimes \mathbb{C}^2$$



$$\psi(x) \mapsto \Phi(y) = \begin{pmatrix} \phi_+(y) \\ \phi_-(y) \end{pmatrix} = (U\psi)(y) = \begin{pmatrix} \psi(y) \\ \psi(-y) \end{pmatrix}$$

Momentum on the half-line

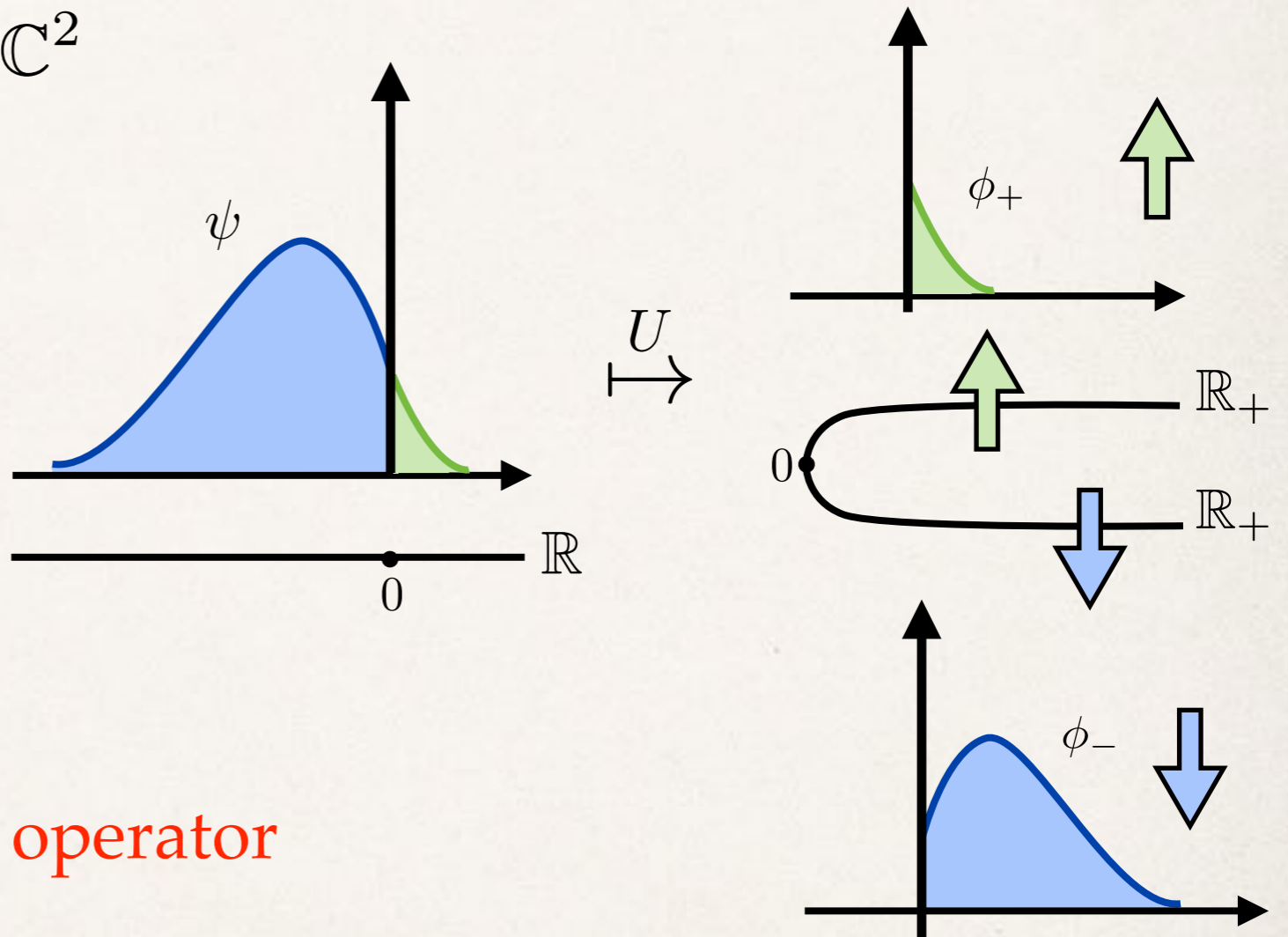
$$U : L^2(\mathbb{R}) \rightarrow L^2(\mathbb{R}_+) \otimes \mathbb{C}^2$$

Image of the momentum:

$$\tilde{p} = U p U^\dagger$$

$$\tilde{p} = -i\hbar \frac{d}{dy} \otimes \sigma_z \quad \text{Dirac operator}$$

$$D(\tilde{p}) = \{ \Phi \in H^1(\mathbb{R}_+) \otimes \mathbb{C}^2 \mid \Phi(0) = \sigma_x \Phi(0) \} \quad \text{QBC}$$

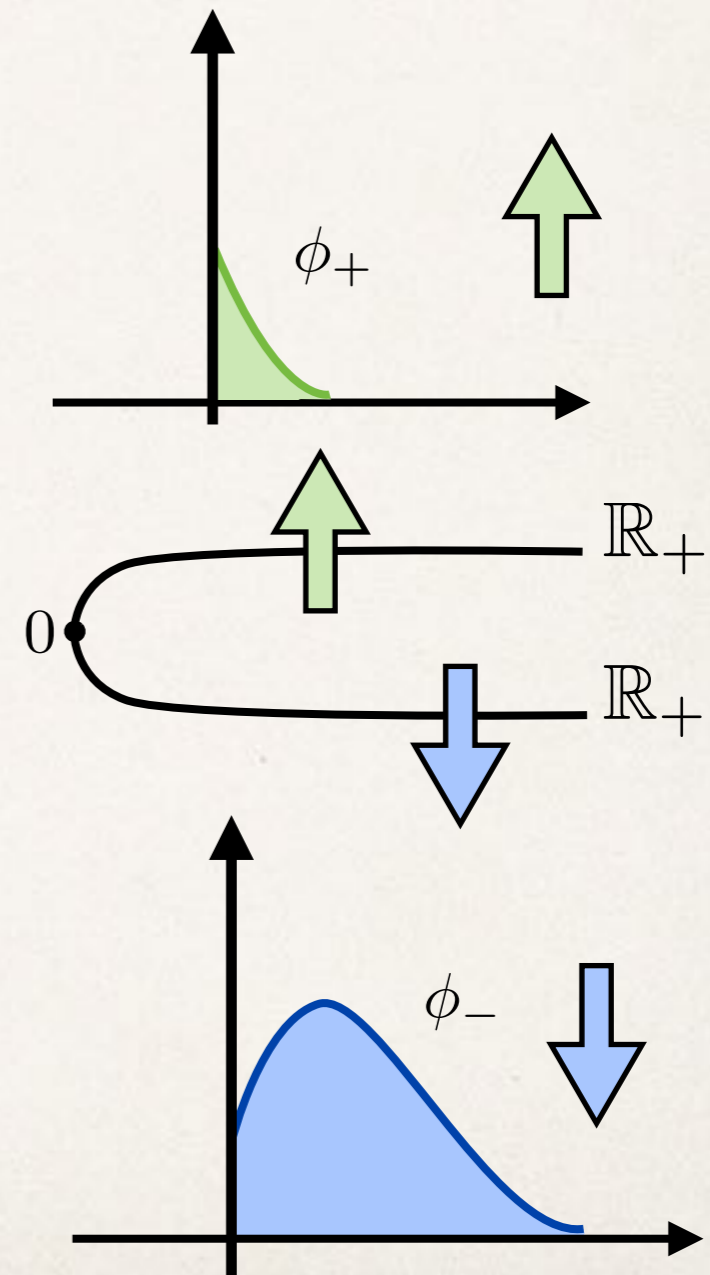


Momentum on the half-line

$$\tilde{p} = -i\hbar \frac{d}{dy} \otimes \sigma_z$$

$$\Phi(0) = \sigma_x \Phi(0)$$

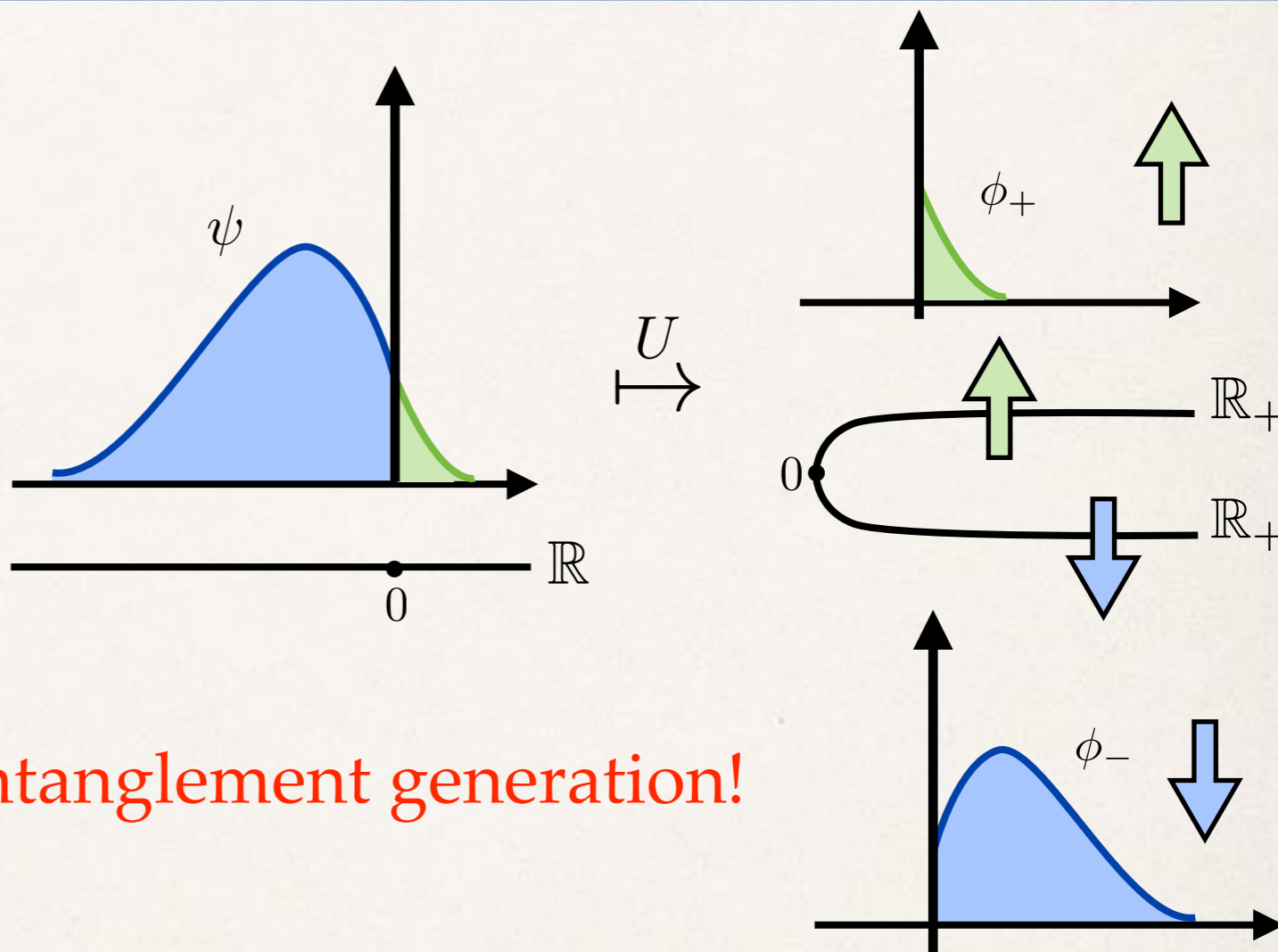
- **spin-1/2 particle:** Both momentum and spin flips when bouncing off the boundary. Helicity is conserved;
- **spinless particle + detector:** the detector clicks when the particle bounces off the boundary.



Entanglement vs self-adjointness

Factorized initial state

$$\phi \otimes \frac{|\uparrow\rangle + |\downarrow\rangle}{\sqrt{2}}$$



Entanglement generation!

Evolution:

$$e^{-itp \otimes \sigma_z} \left(\phi(x) \otimes \frac{|\uparrow\rangle + |\downarrow\rangle}{\sqrt{2}} \right) = \phi(x - t) \otimes \frac{|\uparrow\rangle}{\sqrt{2}} + \phi(x + t) \otimes \frac{|\downarrow\rangle}{\sqrt{2}}$$

Momentum on an interval

Consider the momentum on the circle

$$p = -i\hbar \frac{d}{dx}$$

$$D(p) = H^1(\mathbb{S}) = \{ \psi \in H^1[-\pi, \pi] \mid \psi(-\pi) = \psi(\pi) \}$$

first Sobolev space

It is self-adjoint.

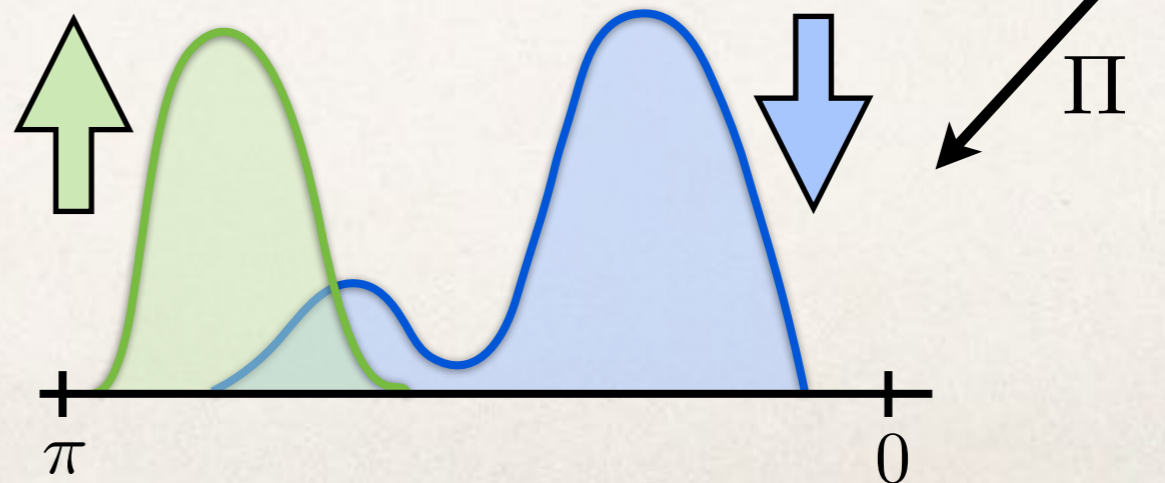
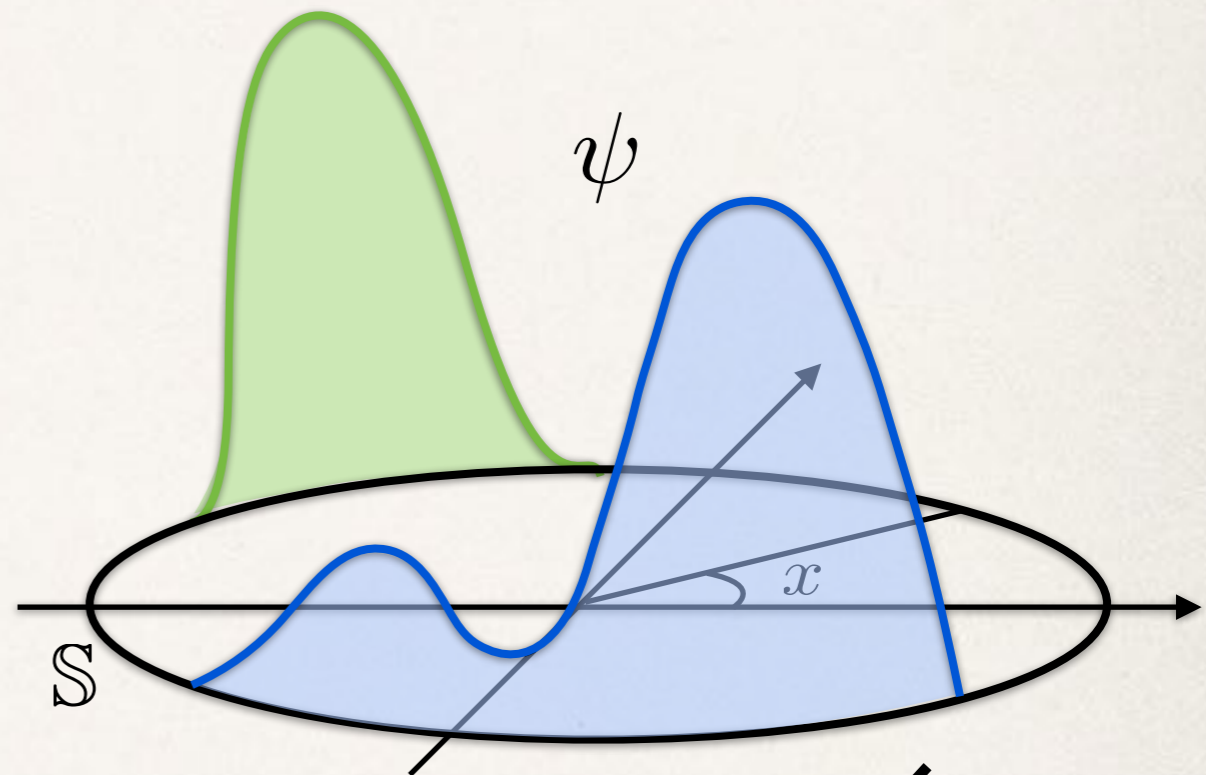
Need a **unitary** mapping onto the interval

Momentum on an interval

Unitary

$$U : L^2(\mathbb{S}) \rightarrow L^2(0, \pi) \otimes \mathbb{C}^2$$

$$\psi(x) \mapsto (U\psi)(y) = \begin{pmatrix} \psi(y) \\ \psi(-y) \end{pmatrix}$$



Ancillary spin

Momentum on an interval

Unitary

$$U : L^2(\mathbb{S}) \rightarrow L^2(0, \pi) \otimes \mathbb{C}^2$$

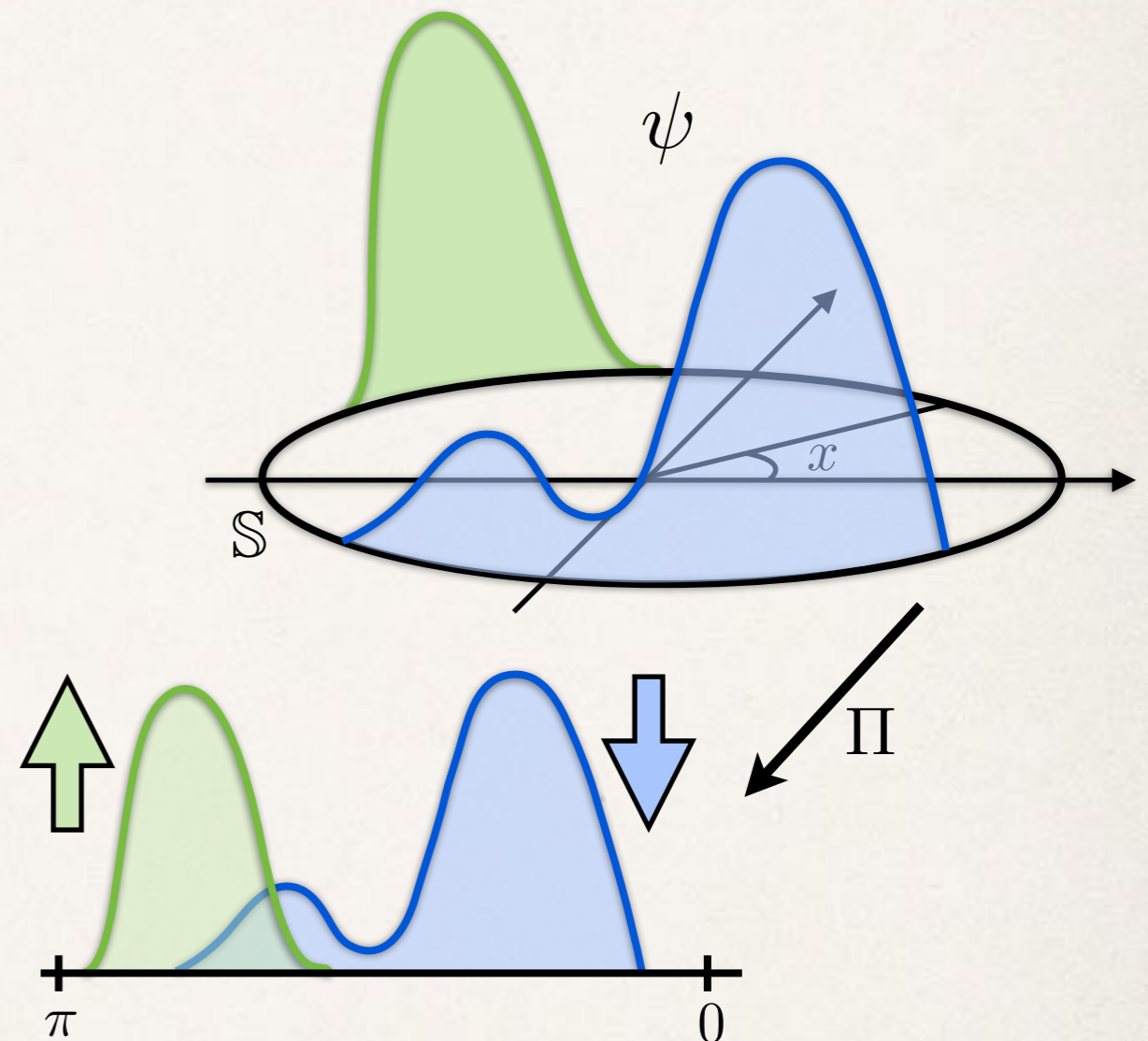
Image of the momentum

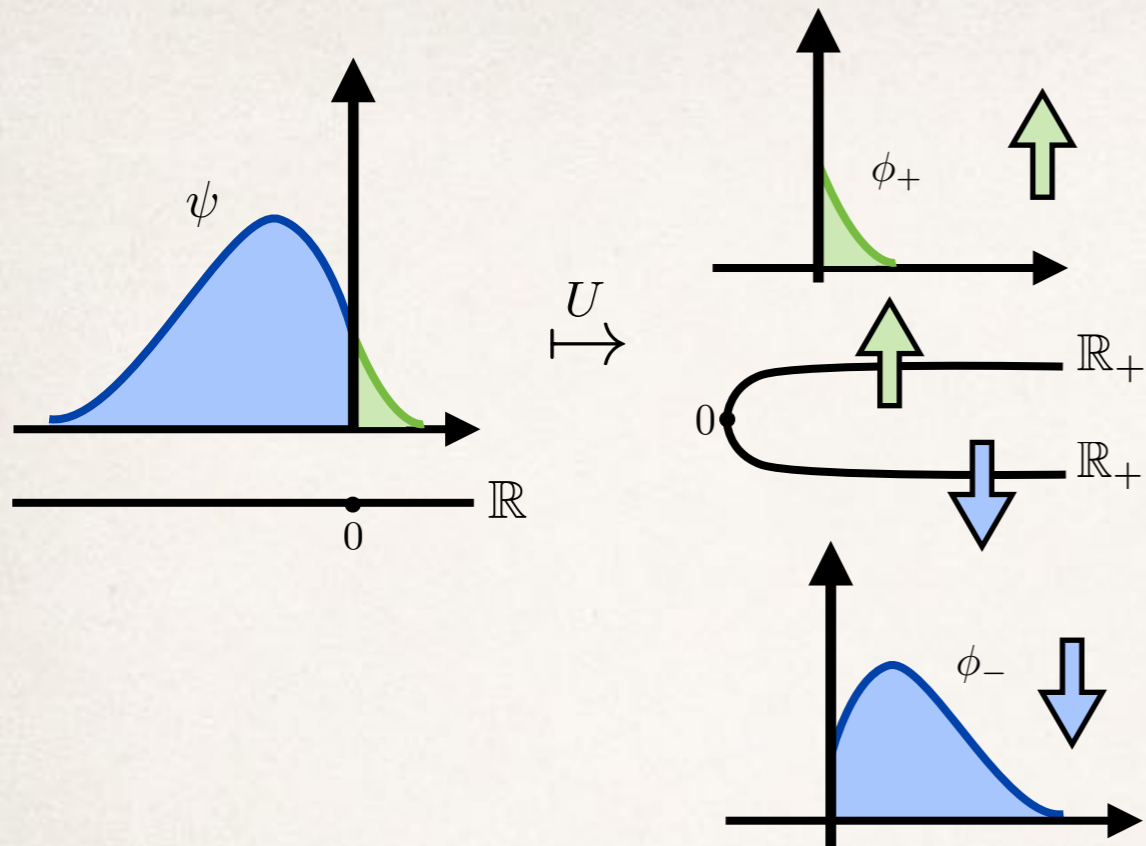
$$\tilde{p} = U p U^\dagger$$

$$\tilde{p} = -i\hbar \frac{d}{dy} \otimes \sigma_z$$

Quantum boundary conditions

$$D(\tilde{p}) = \{ \Phi \in H^1[0, \pi] \otimes \mathbb{C}^2 \mid \Phi(0) = \sigma_x \Phi(0), \Phi(\pi) = \sigma_x \Phi(\pi) \}$$





Happy birthday
Bal!

