

# Gamma rays from molecular clouds (part 1)



**Stefano Gabici**  
**APC, Paris**



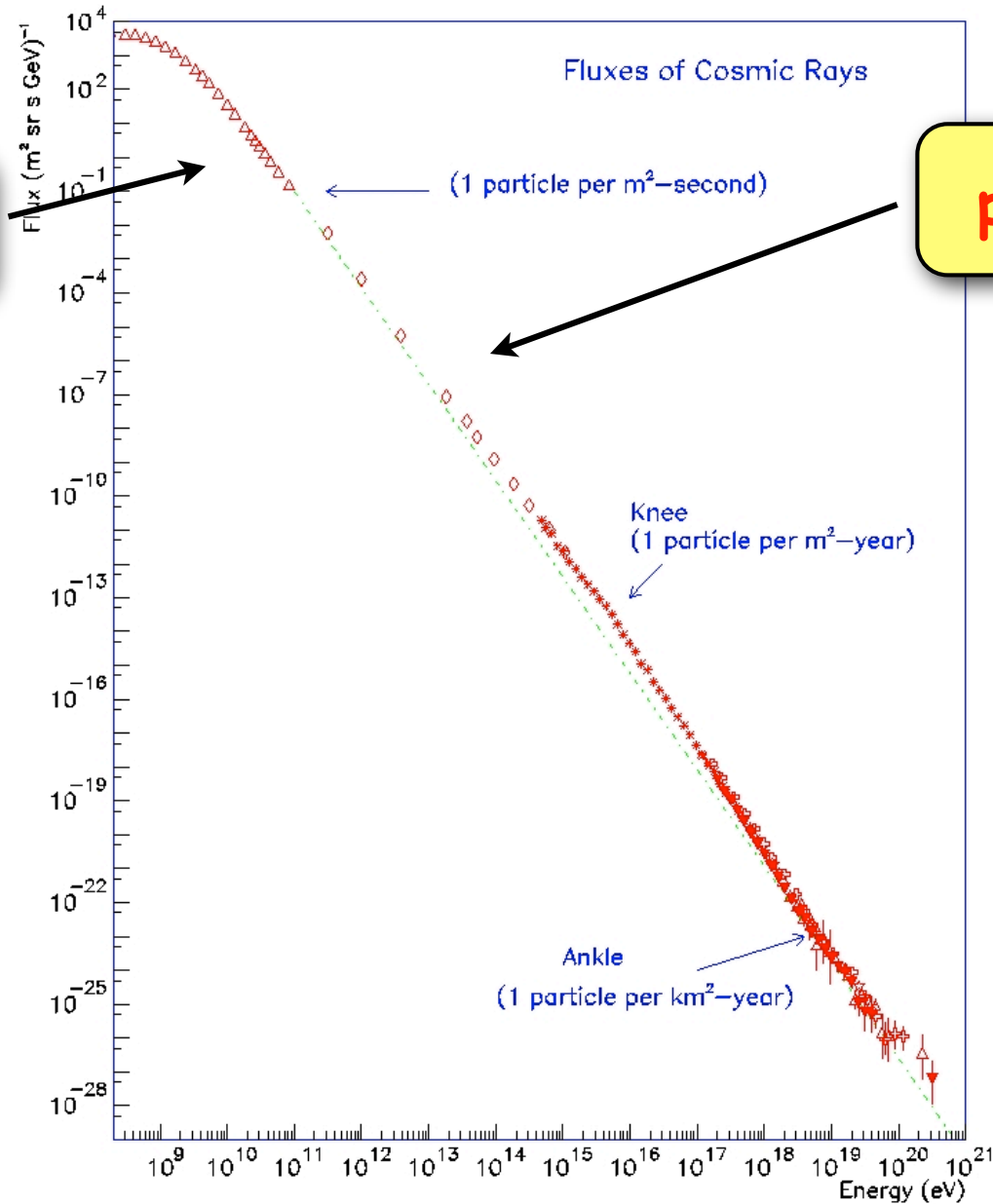
[www.cnrs.fr](http://www.cnrs.fr)

# Overview

- ☑ Brief introduction on **cosmic rays**
- ☑ The **supernova remnant hypothesis** for the origin of cosmic rays
- ☑ Why **gamma ray astronomy**?
- ☑ Why **molecular clouds**?
- ☑ Part 1: **Passive Molecular Clouds** -> how to use MCs as cosmic ray  
barometers
- ☑ Part 2: **Illuminated Molecular Clouds** -> how to use MCs to locate CR  
sources and constrain CR propagation

# The (local) Cosmic Ray spectrum

1 eV/cm<sup>3</sup>



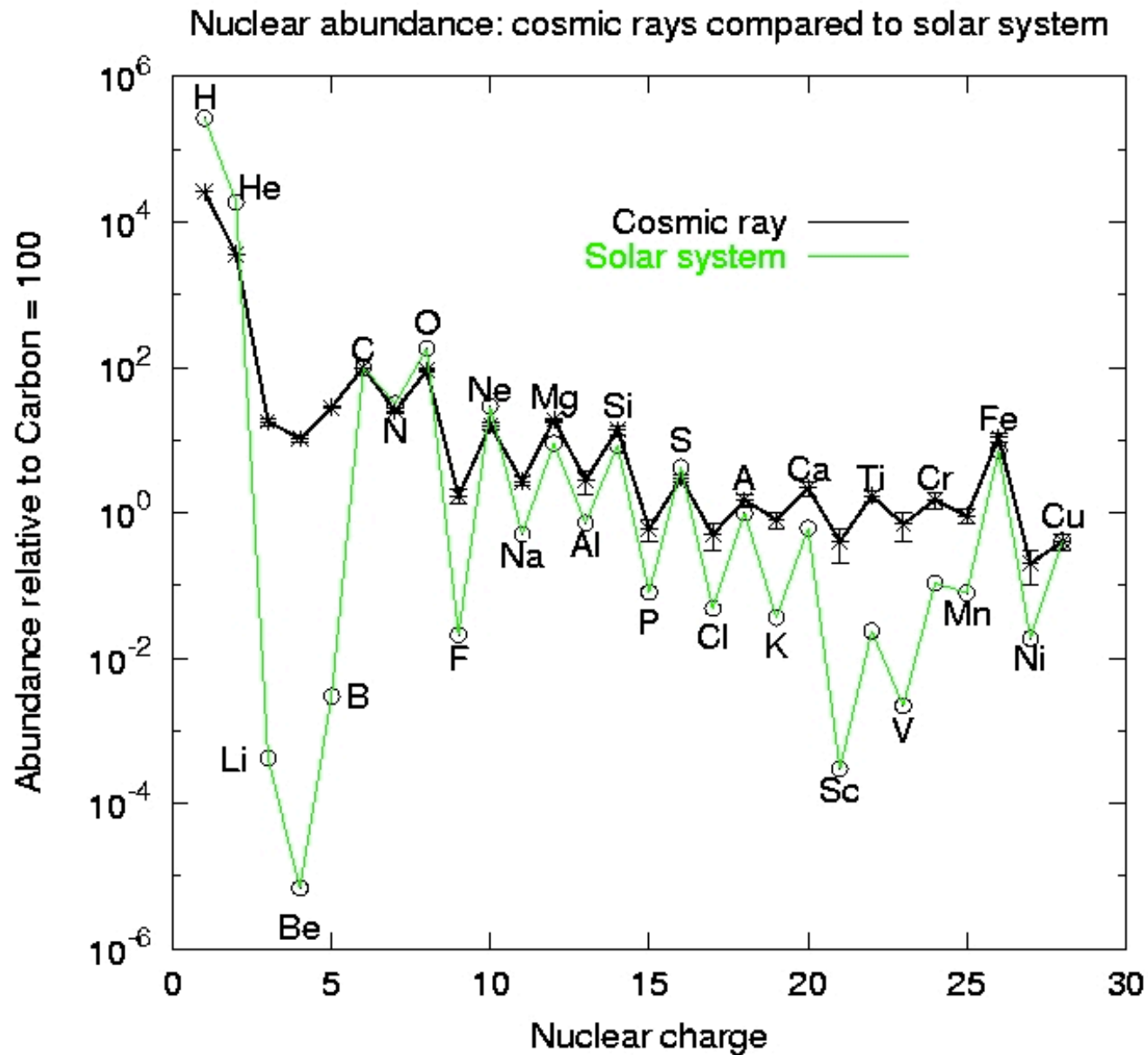
power law

See class by  
T. Bell on  
shock  
acceleration

isotropic

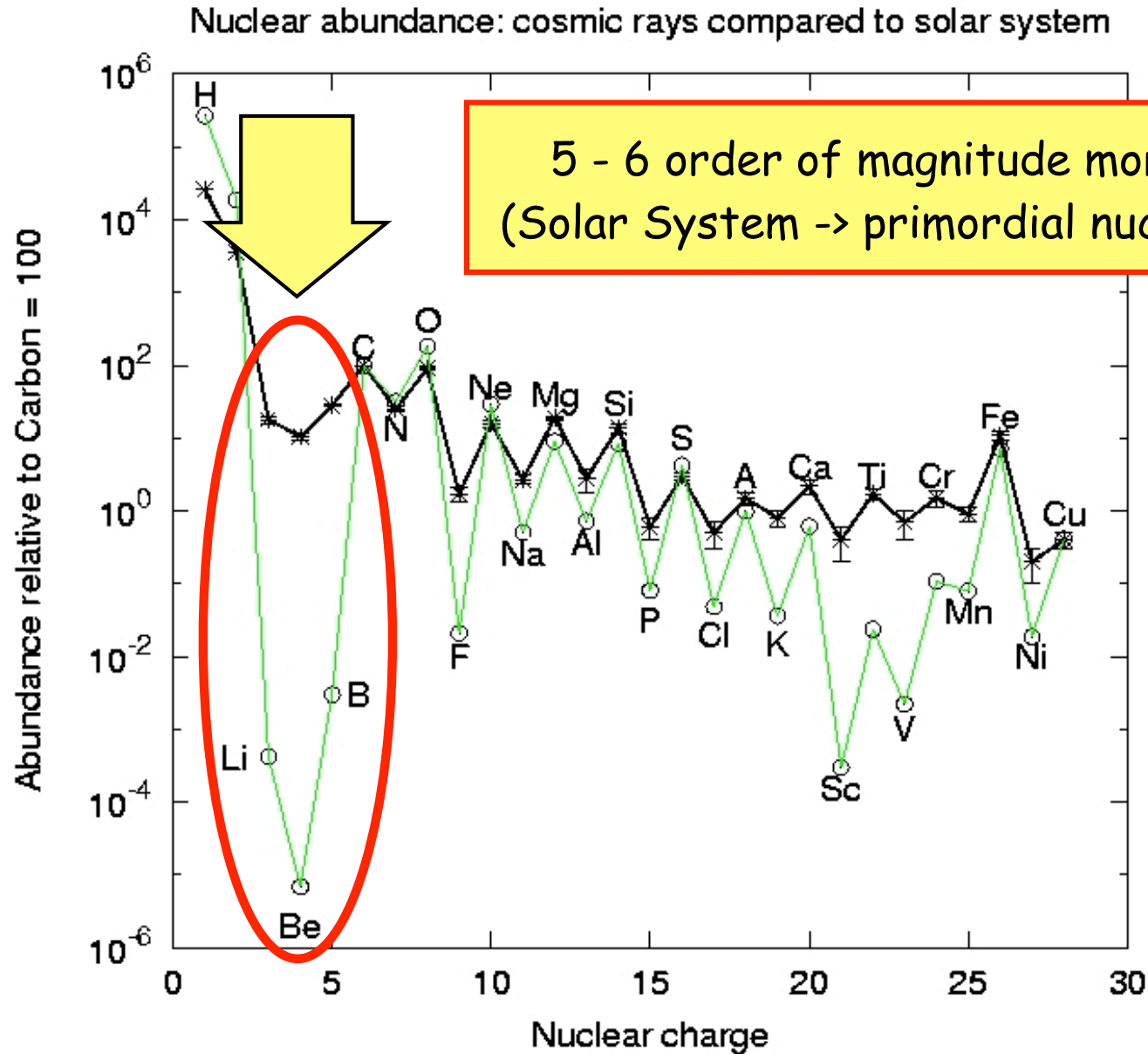
See class by  
A. Watson

# Cosmic Ray composition

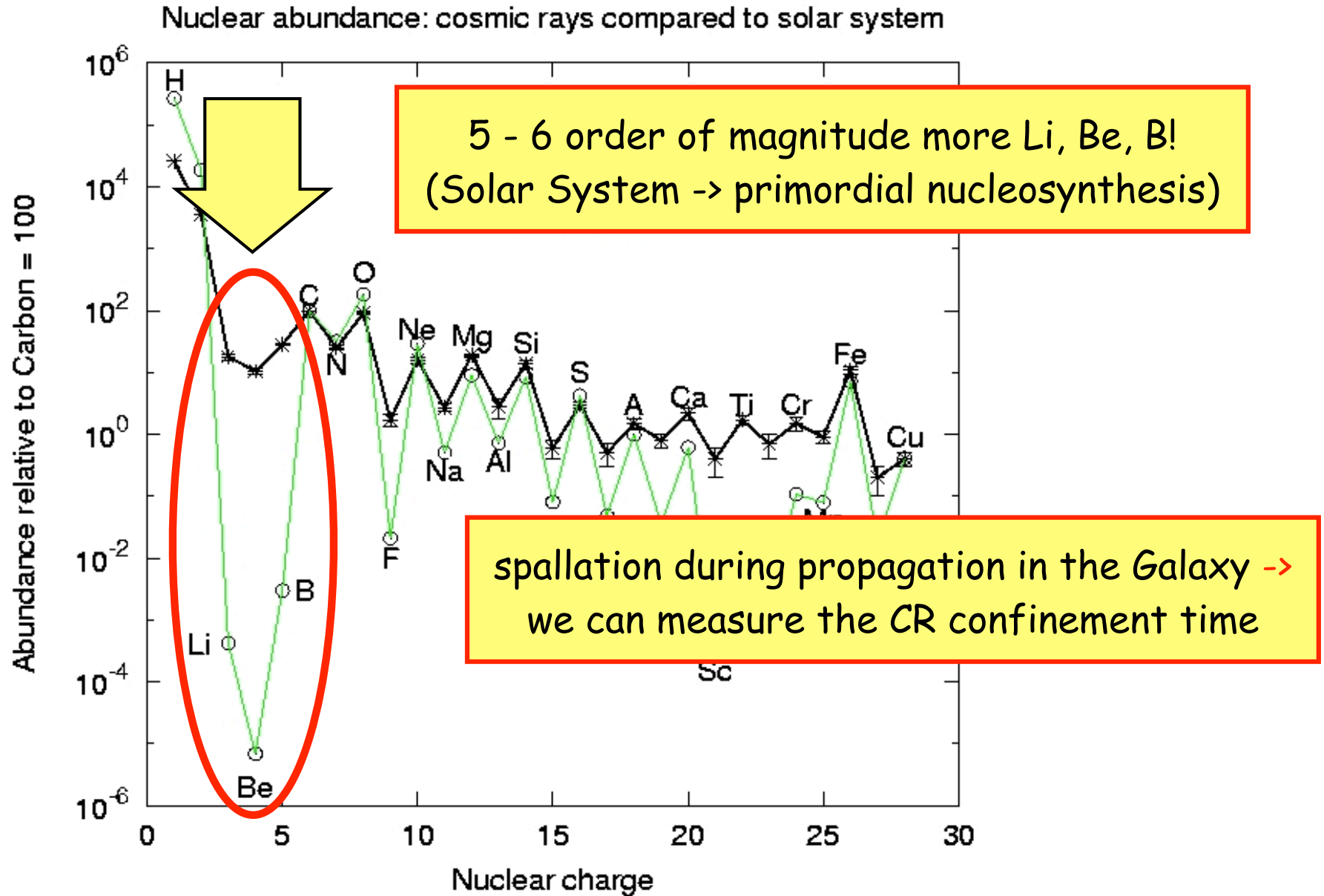




# Cosmic Ray composition

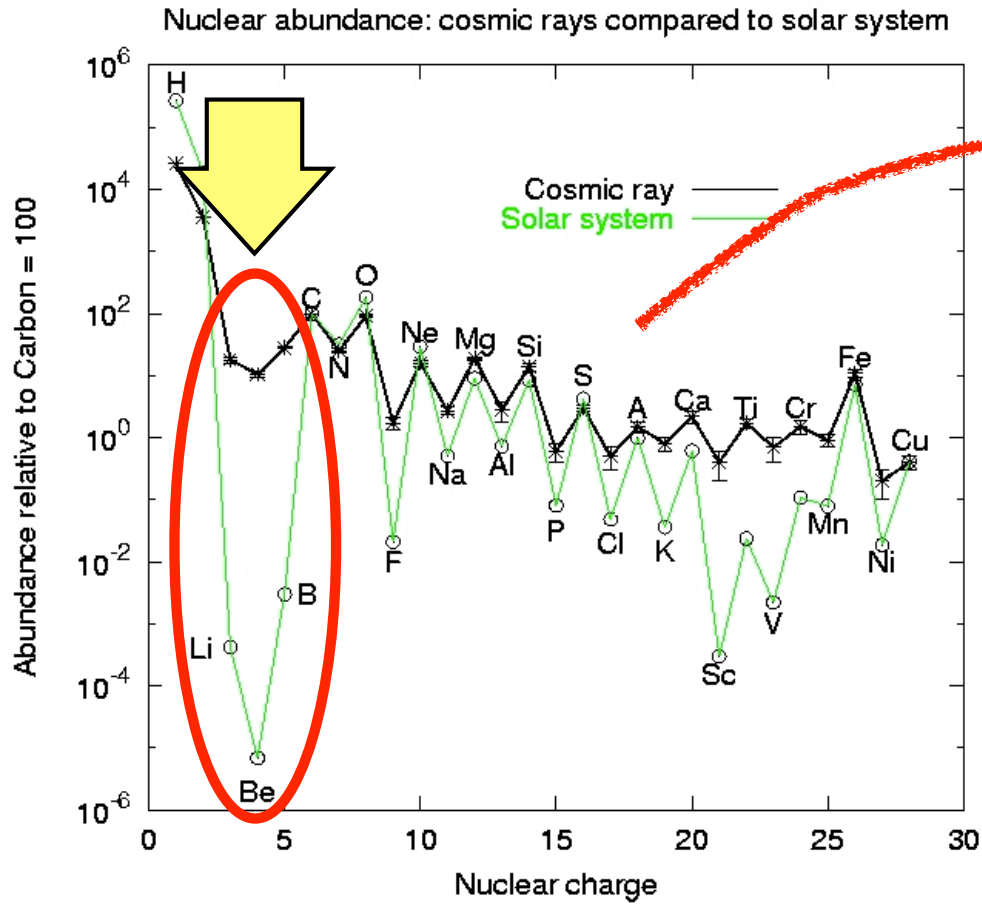


# Cosmic Ray composition



# A remarkable "coincidence"

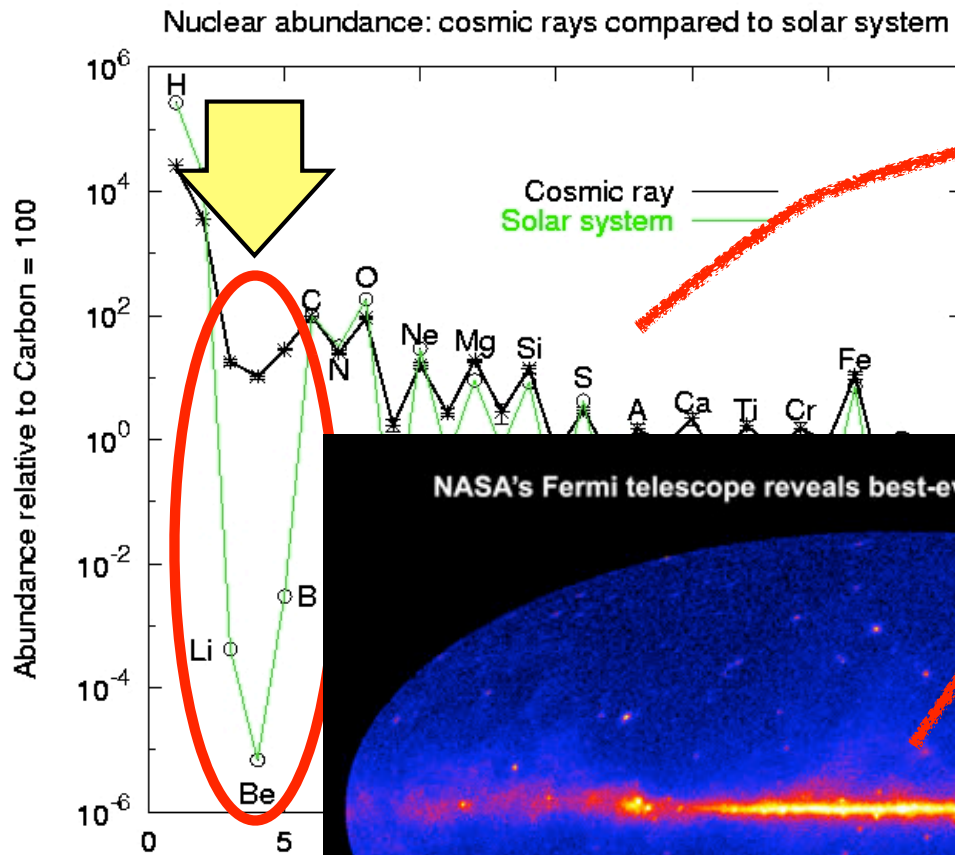
~ Baade & Zwicky, 1934



CR escape time

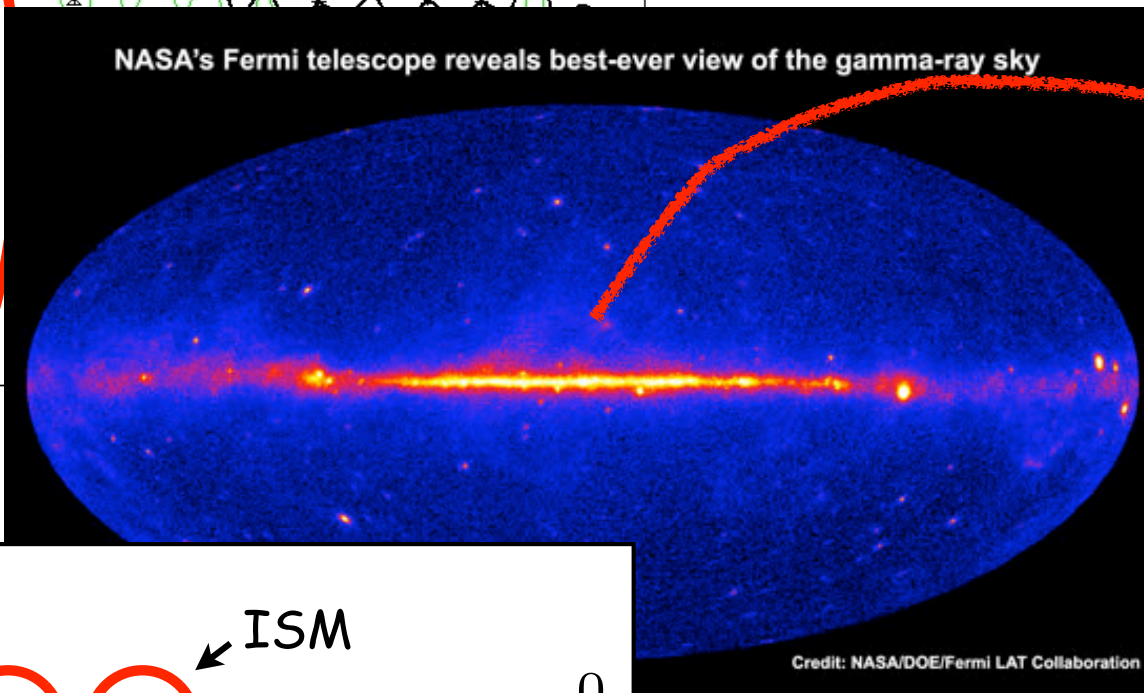
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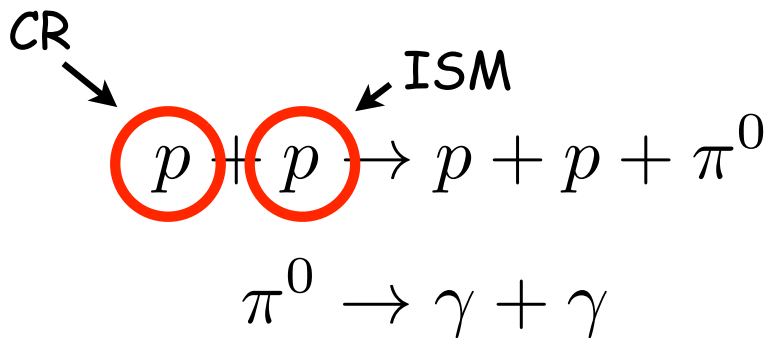


CR escape time

-> power of CR sources  $3 \times 10^{40}$  erg/s

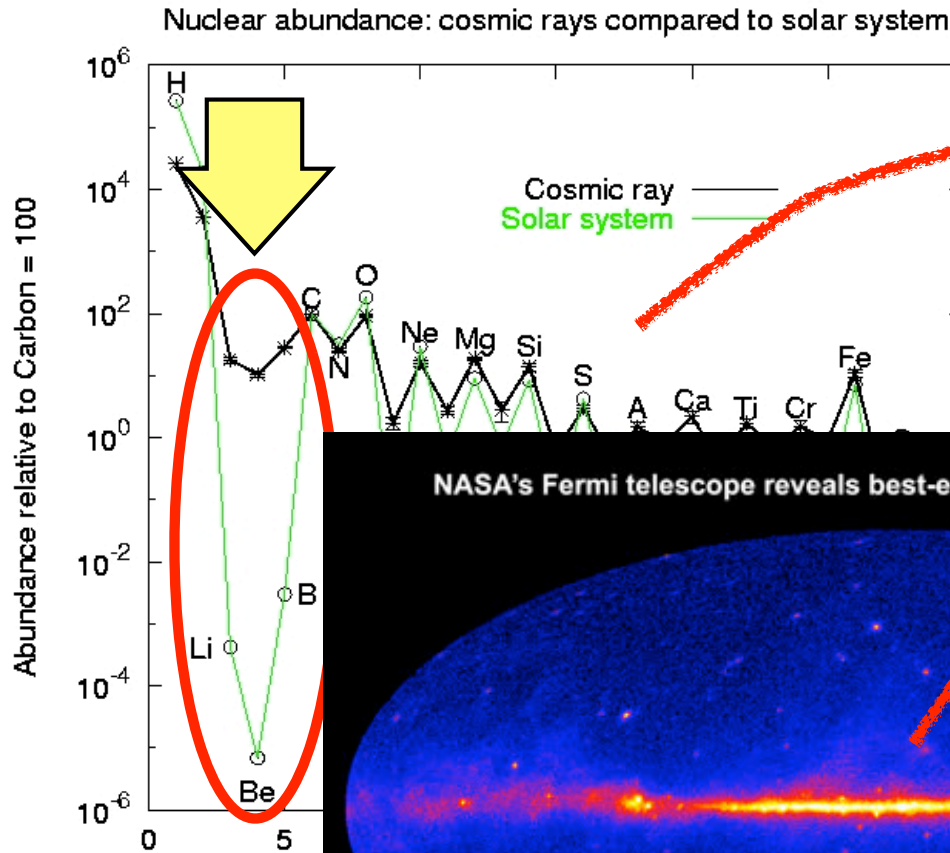


CR total energy



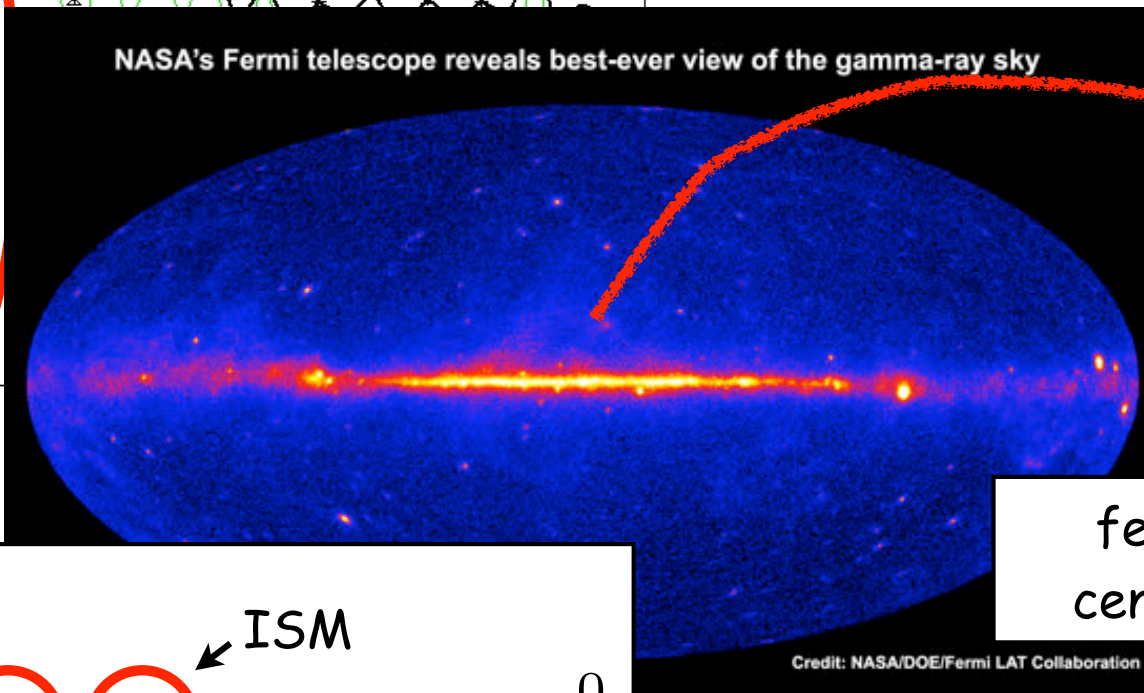
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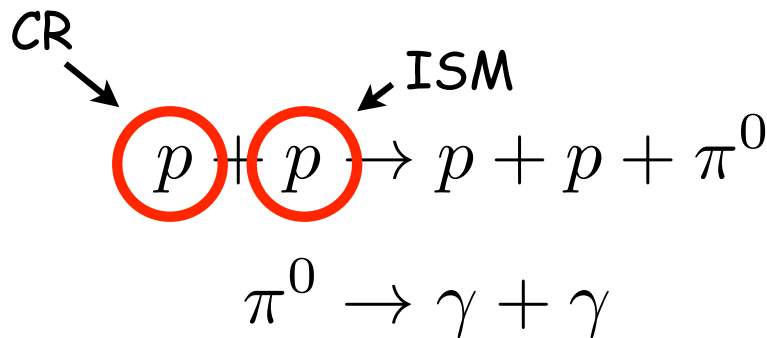
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CR total energy

few supernovae per century in the Galaxy

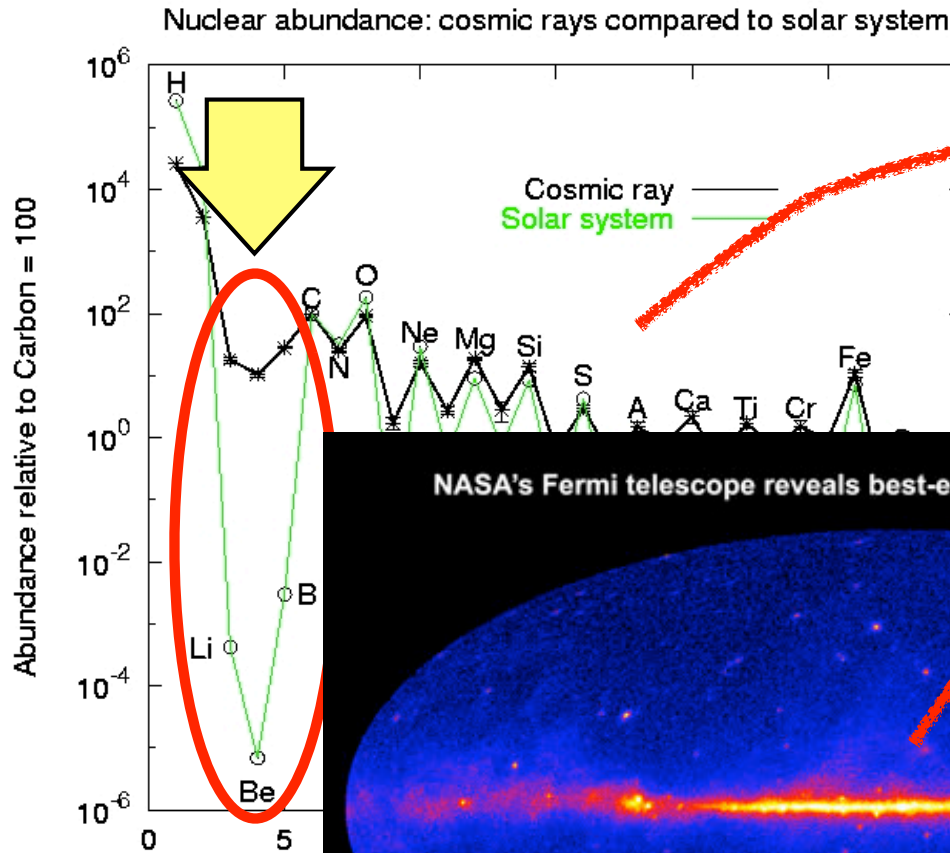


-> power of SuperNovae  $3 \times 10^{41}$  erg/s



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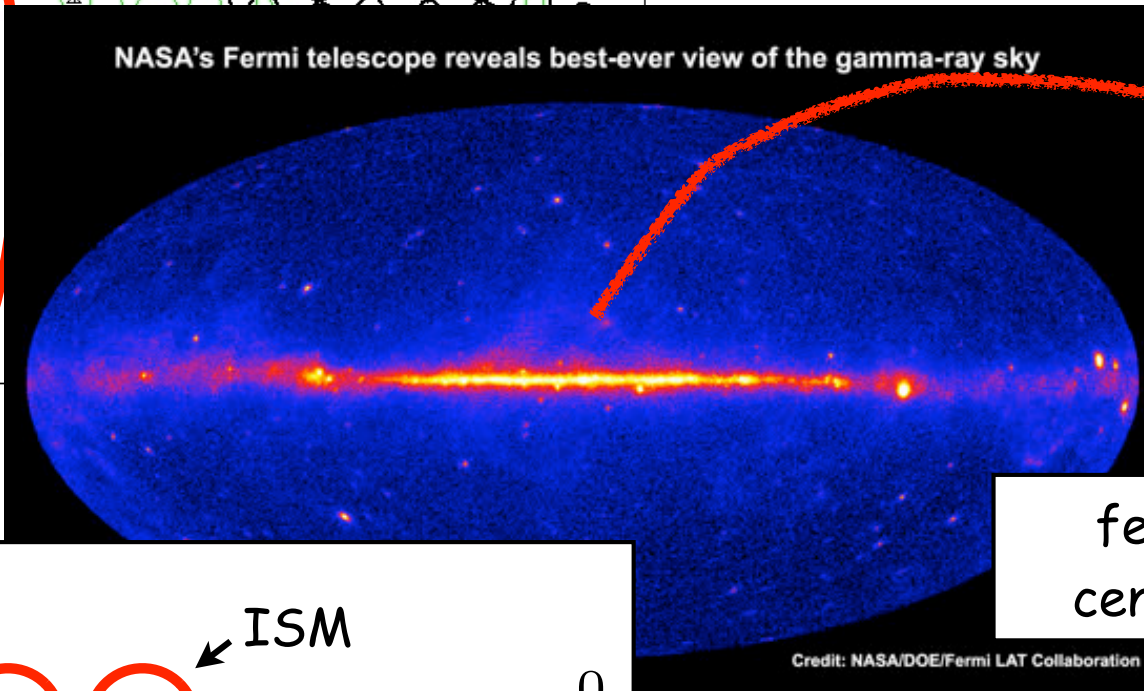
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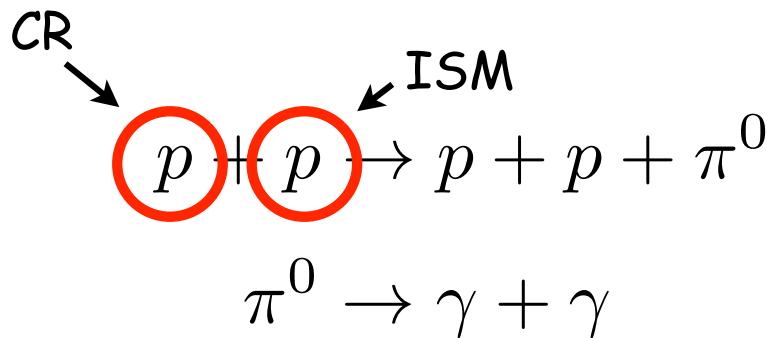
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NASA's Fermi telescope reveals best-ever view of the gamma-ray sky



CR total energy

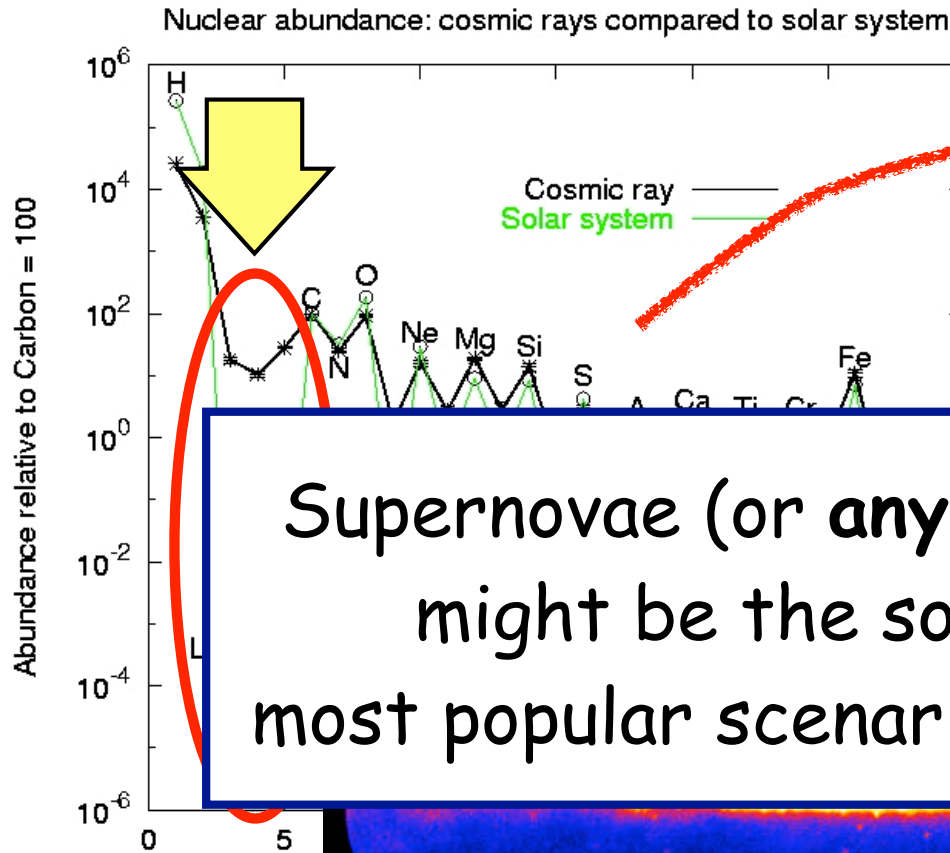
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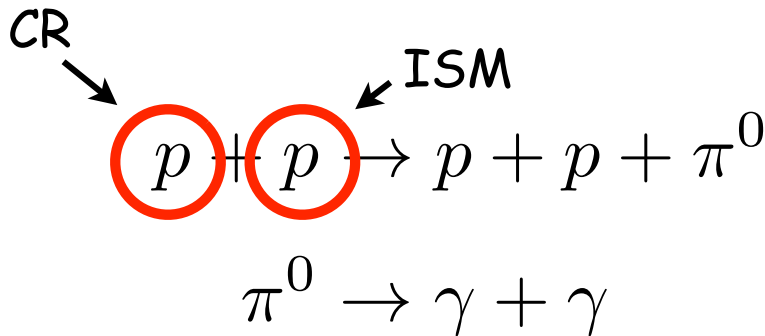
CR escape time

-> power of CR sources  $3 \times 10^{40}$  erg/s

Supernovae (or anything connected to them) might be the sources of cosmic rays: most popular scenario -> **supernova remnants**

energy

few supernovae per century in the Galaxy



Credit: NASA/DOE/Fermi LAT Collaboration

-> power of SuperNovae  $3 \times 10^{41}$  erg/s

# Cosmic Ray composition: spallation

**Spallation:** production of light elements as fragmentation products of the interaction of high energy particles with cold matter.

The anomaly is explained if ( $\sim GeV$ ) CRs transverse  $\lambda \approx 5 \text{ g/cm}^2$



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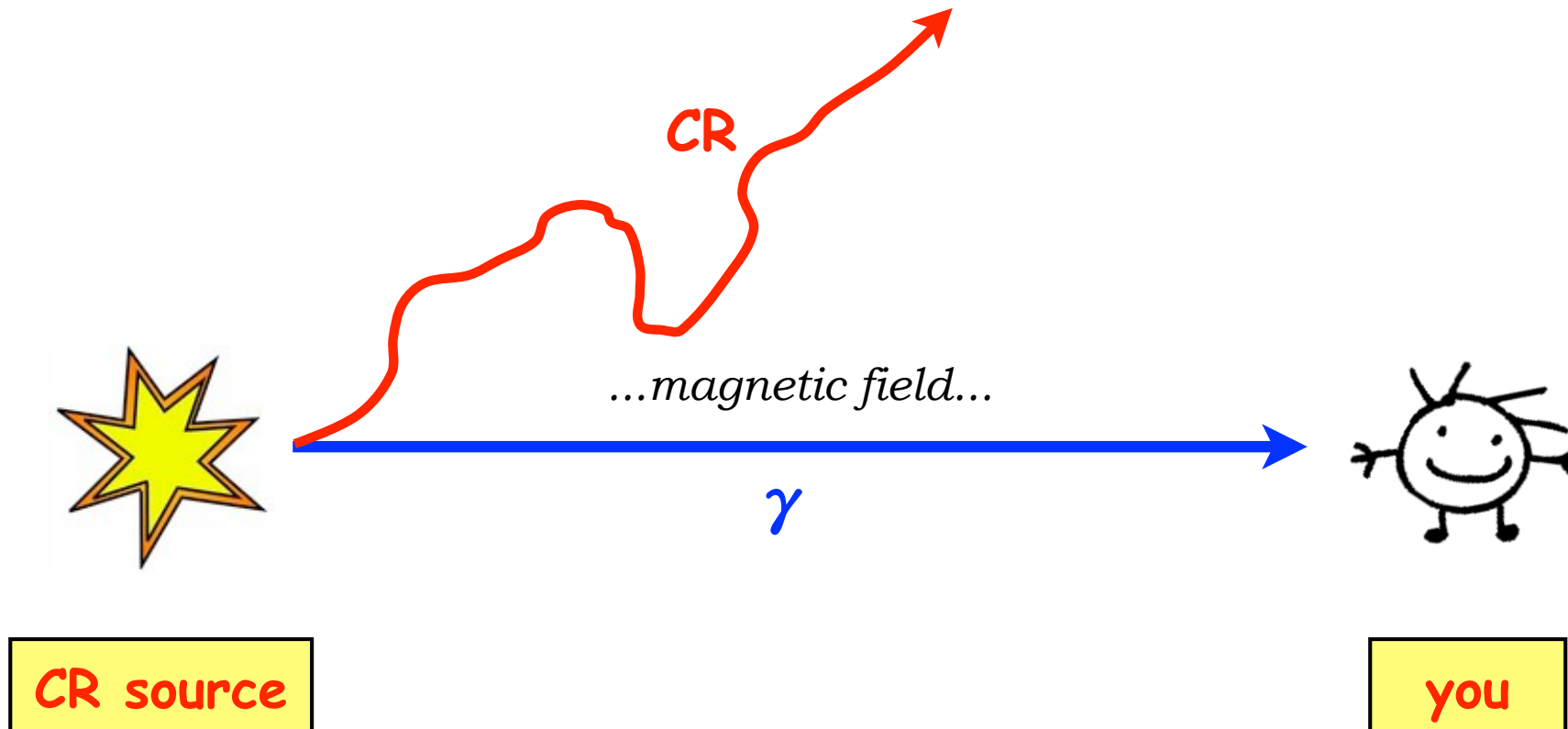
The anomaly is explained if ( $\sim \text{GeV}$ ) CRs transverse  $\lambda \approx 5 \text{ g/cm}^2$

Assuming propagation in the galactic disk:  $l_s = \frac{\lambda}{\rho_{ISM}} \approx 1 \text{ Mpc}$

 much larger than  
the size of the  
disk!!!

CRs don't go straight but are confined in the disk  
-> diffusive behavior -> isotropy!

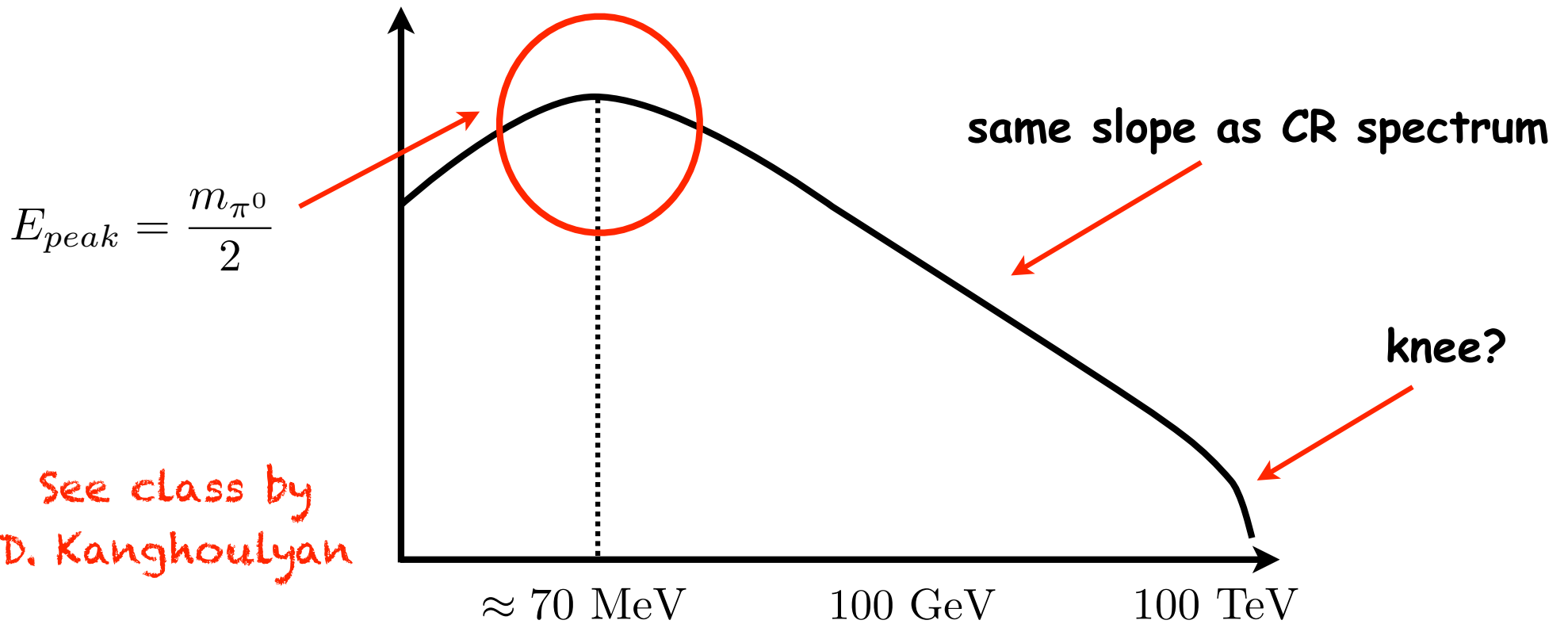
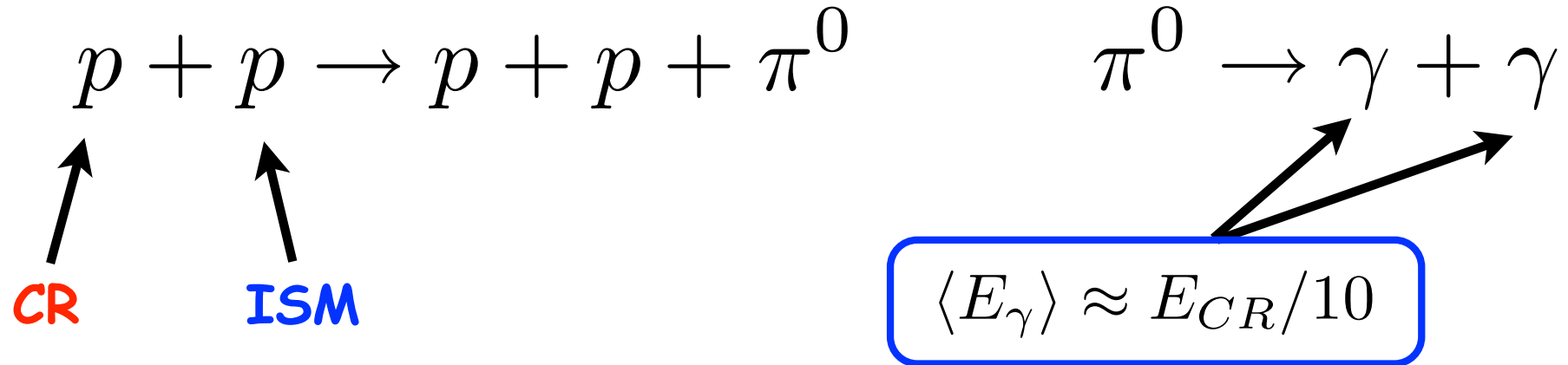
# Why is it so difficult?



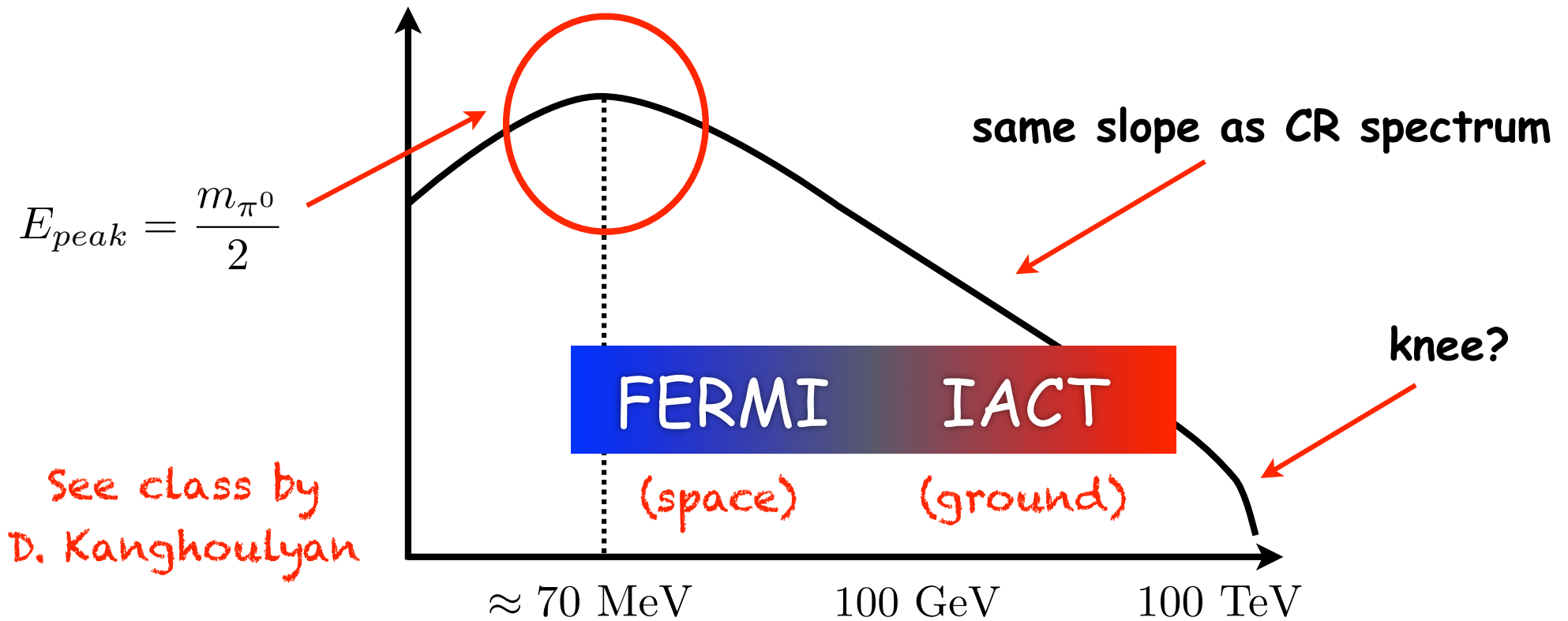
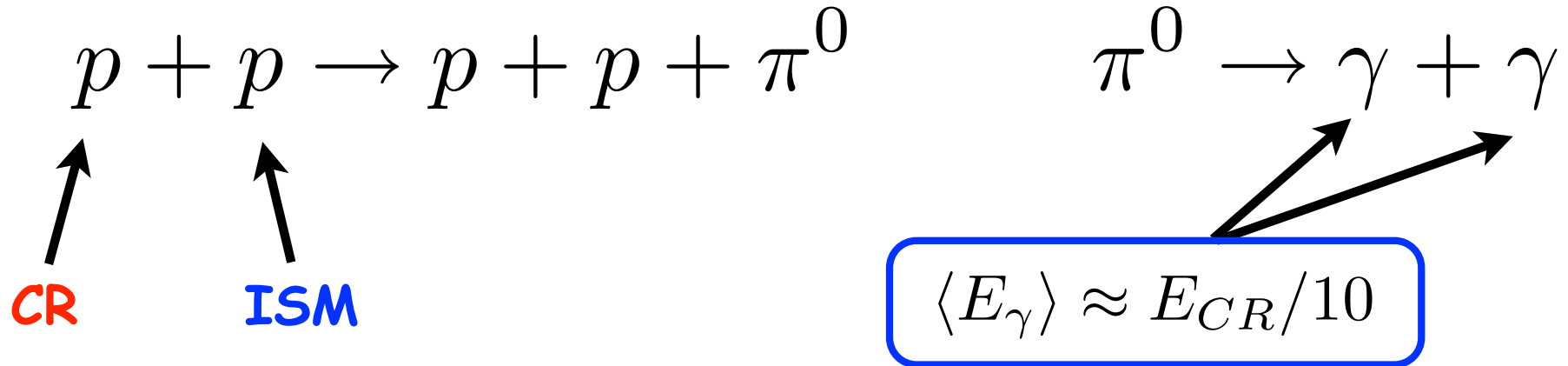
We cannot do CR Astronomy.

Need for indirect identification of CR sources.

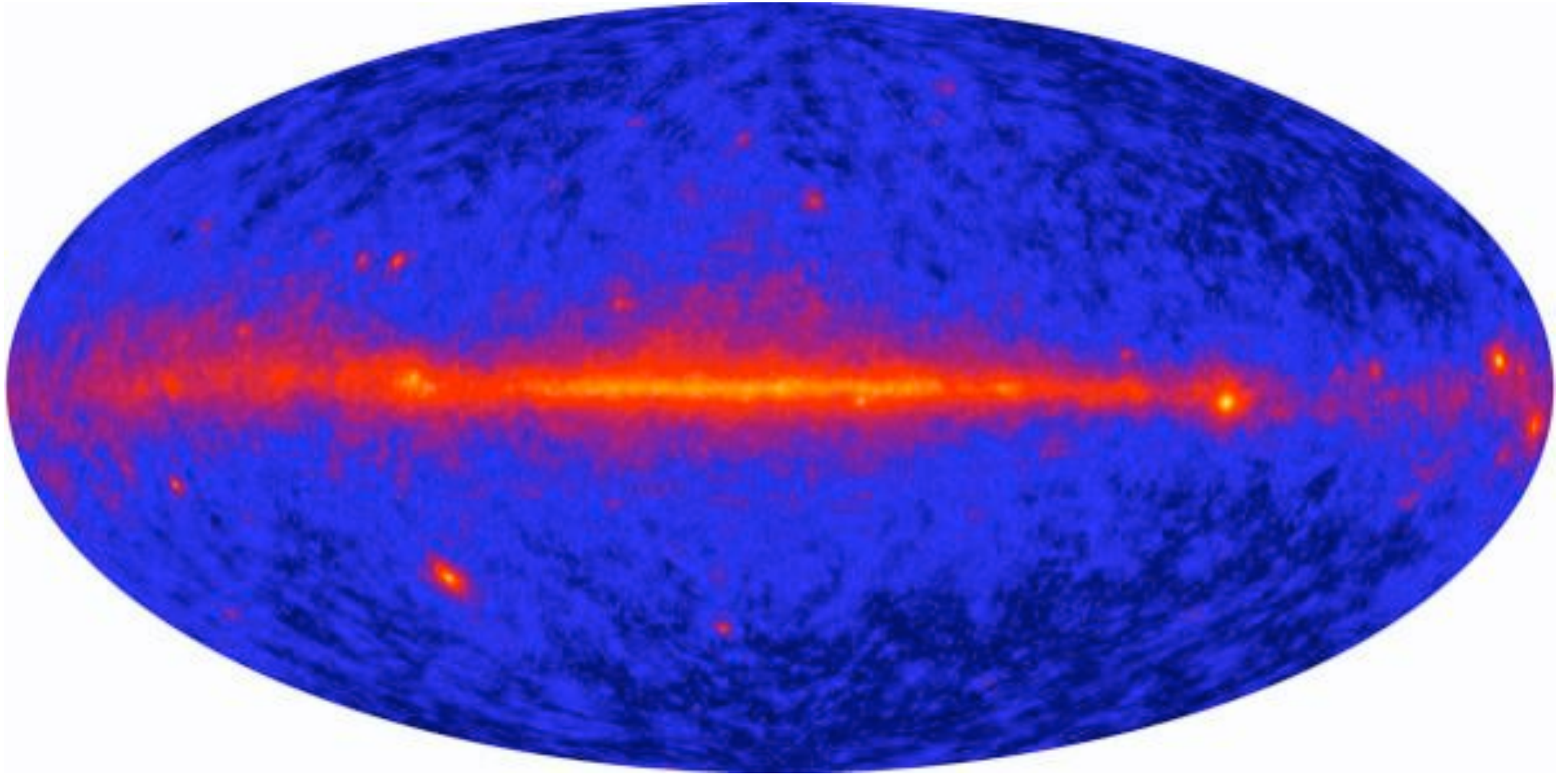
# Gamma-ray astronomy



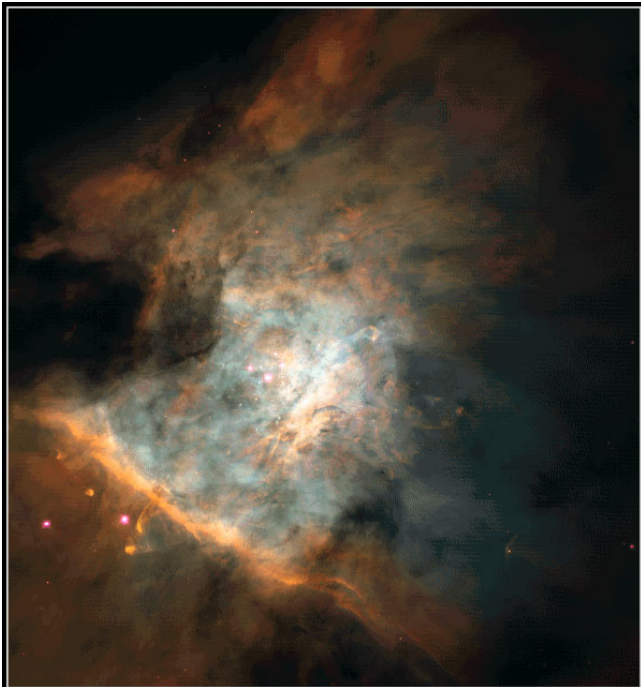
# Gamma-ray astronomy



# The sky @ $E > 100$ MeV (FERMI)



# Why molecular clouds?



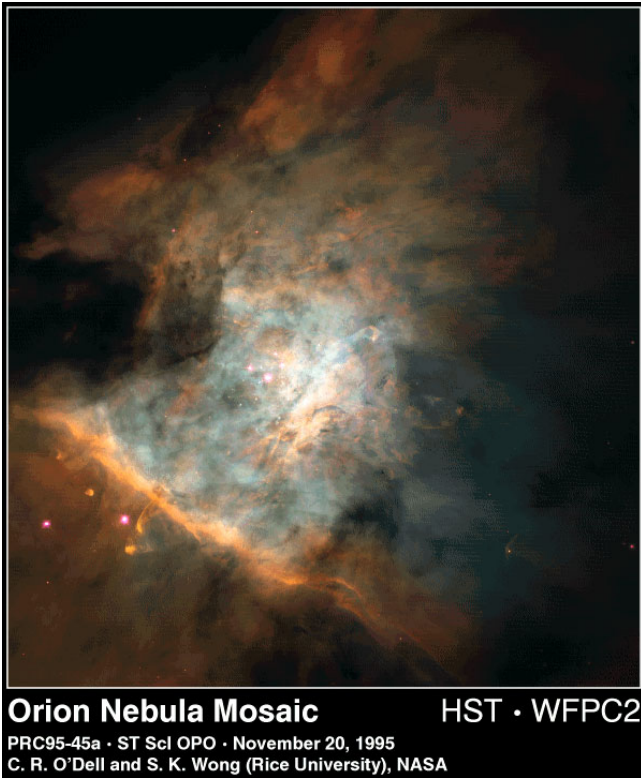
Orion Nebula Mosaic HST · WFPC2  
PRC95-45a · ST ScI OPO · November 20, 1995  
C. R. O'Dell and S. K. Wong (Rice University), NASA

**Molecular Clouds -> sites of star formation**

dense ->  $n \sim 100 \text{ cm}^{-3}$

massive -> Mass up to  $10^6 M_{\odot}$

# Why molecular clouds?



Molecular Clouds -> sites of star formation

dense ->  $n \sim 100 \text{ cm}^{-3}$

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Cloud mass

$$L_{\gamma} \approx \sigma c \int dV n_{CR} n_{ISM} = \sigma c n_{CR} \int dV n_{CR} n_{ISM} \propto M_{cl}$$

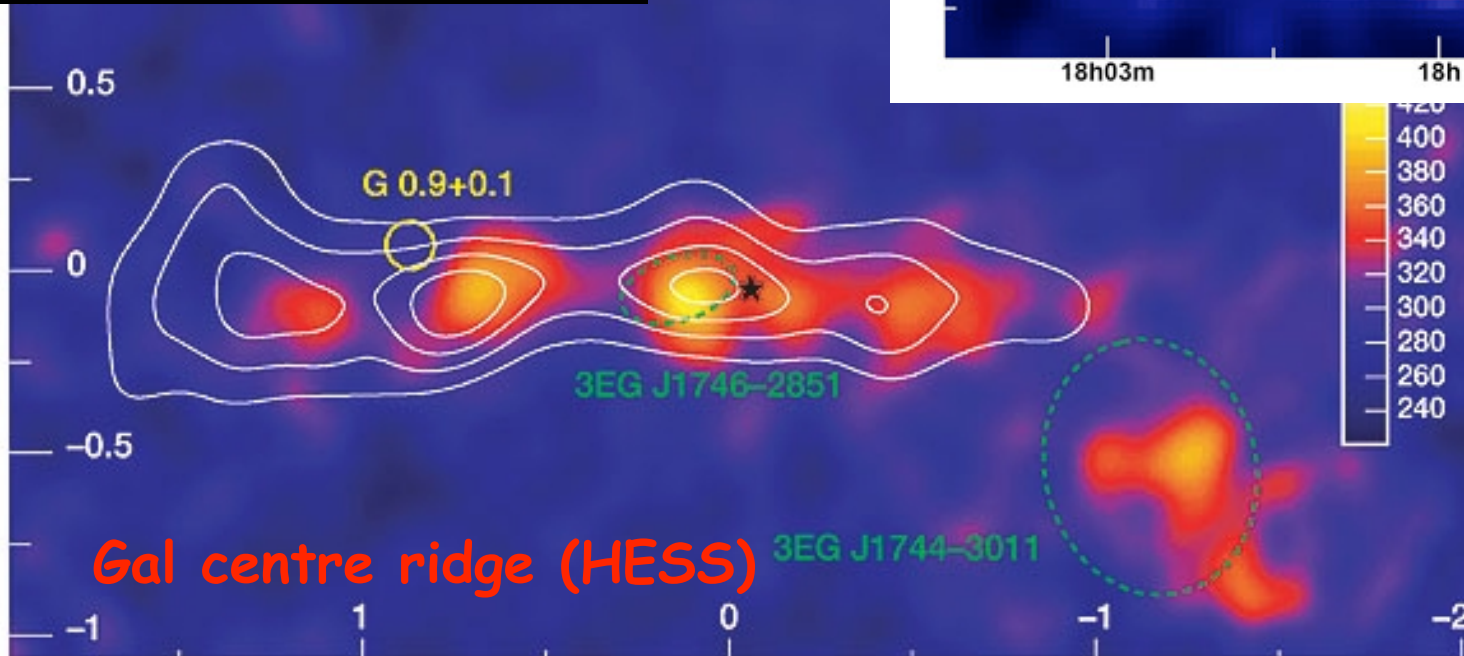
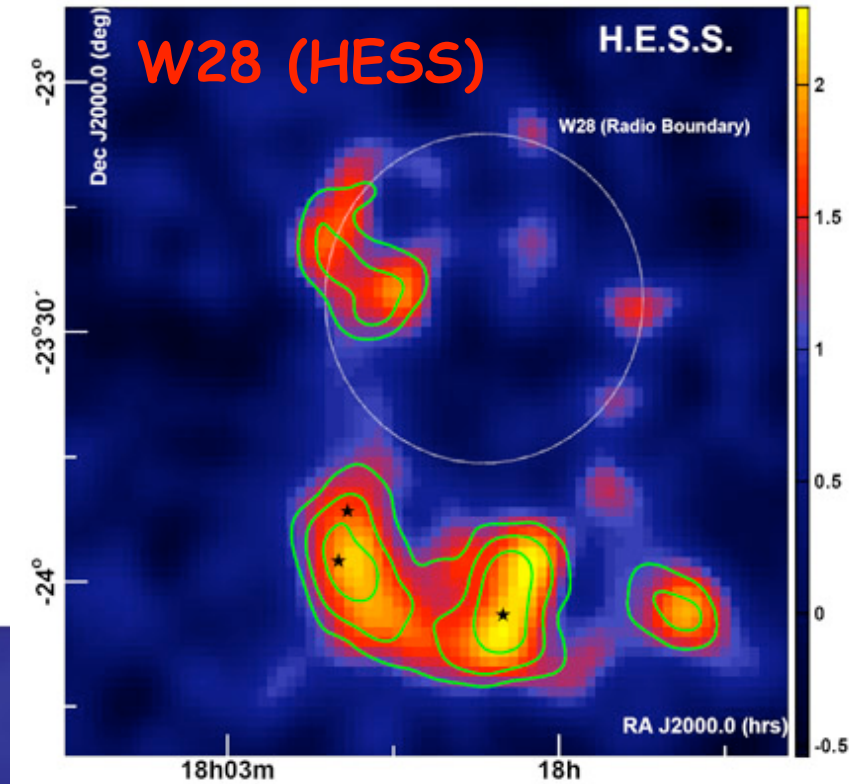
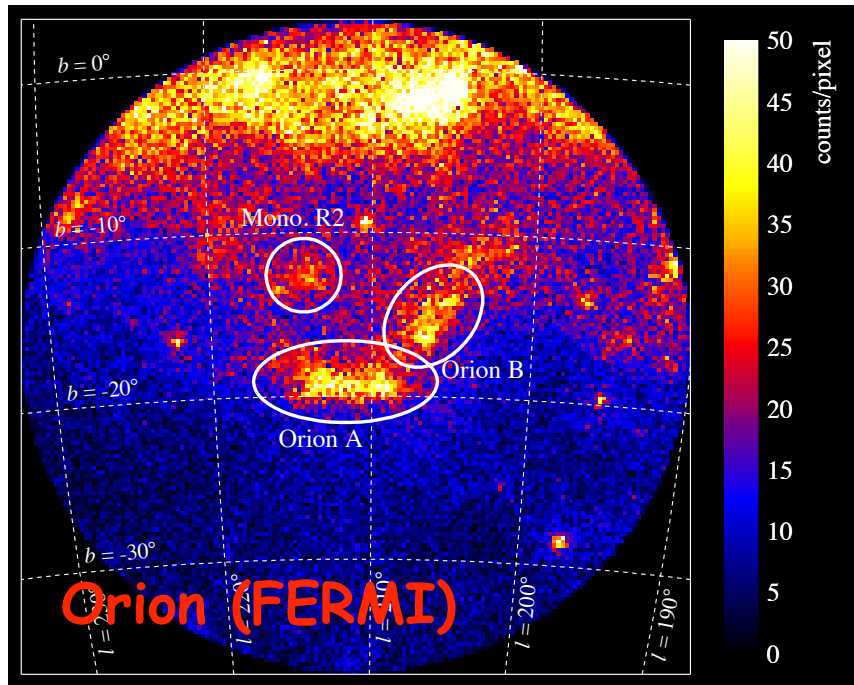
The equation shows the relationship between luminosity and cloud mass. A red 'X' is placed over the  $n_{CR}$  term inside the integral, and a red arrow points from the  $n_{CR}$  term outside the integral to the  $n_{CR}$  term inside the integral, indicating that the luminosity is proportional to the cloud mass.

...because they are massive



# Molecular clouds are gamma ray sources

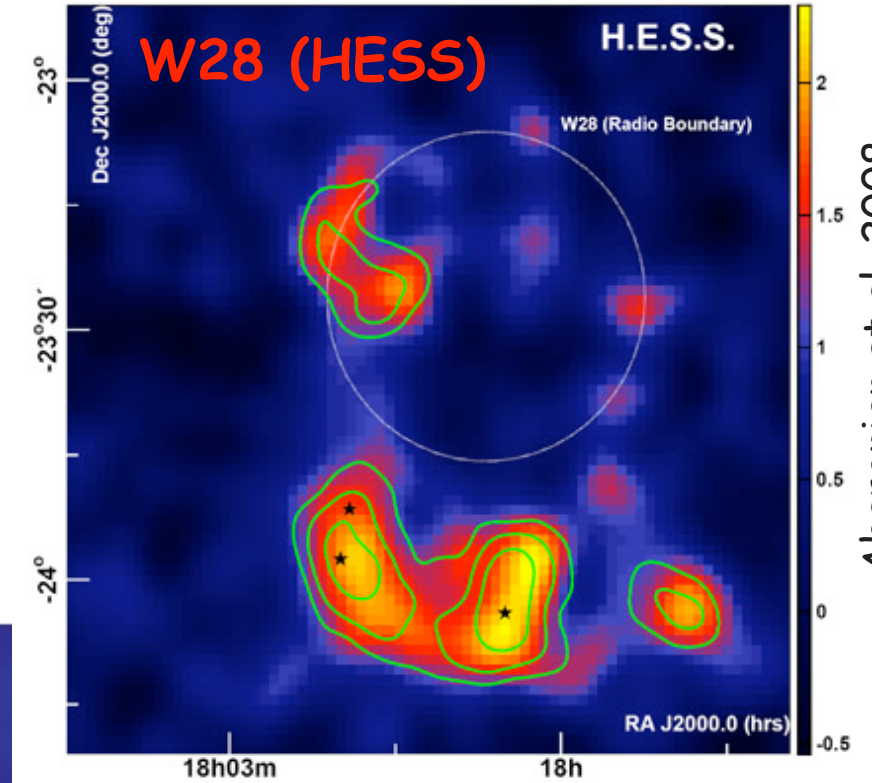
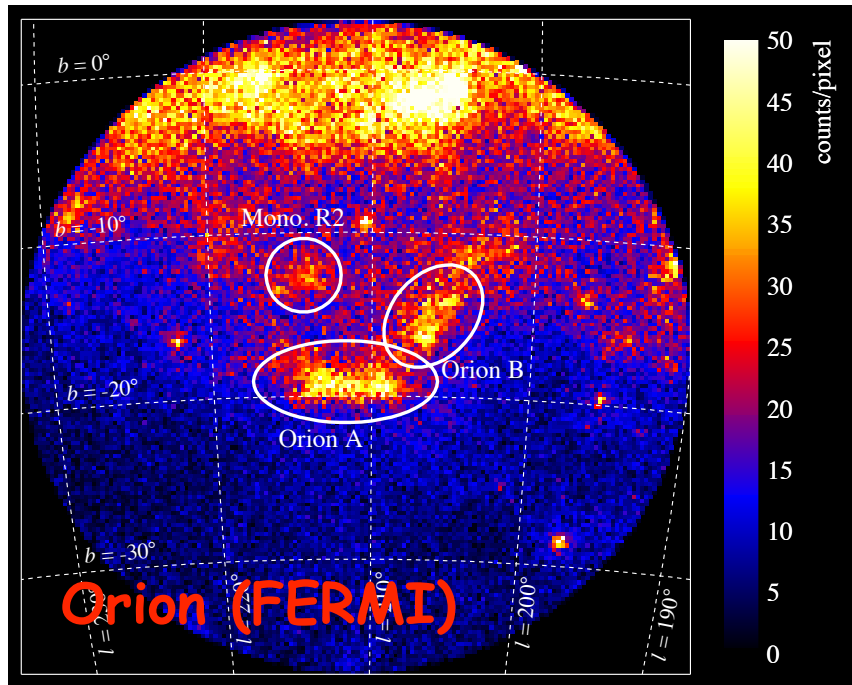
Okumura et al, 2009





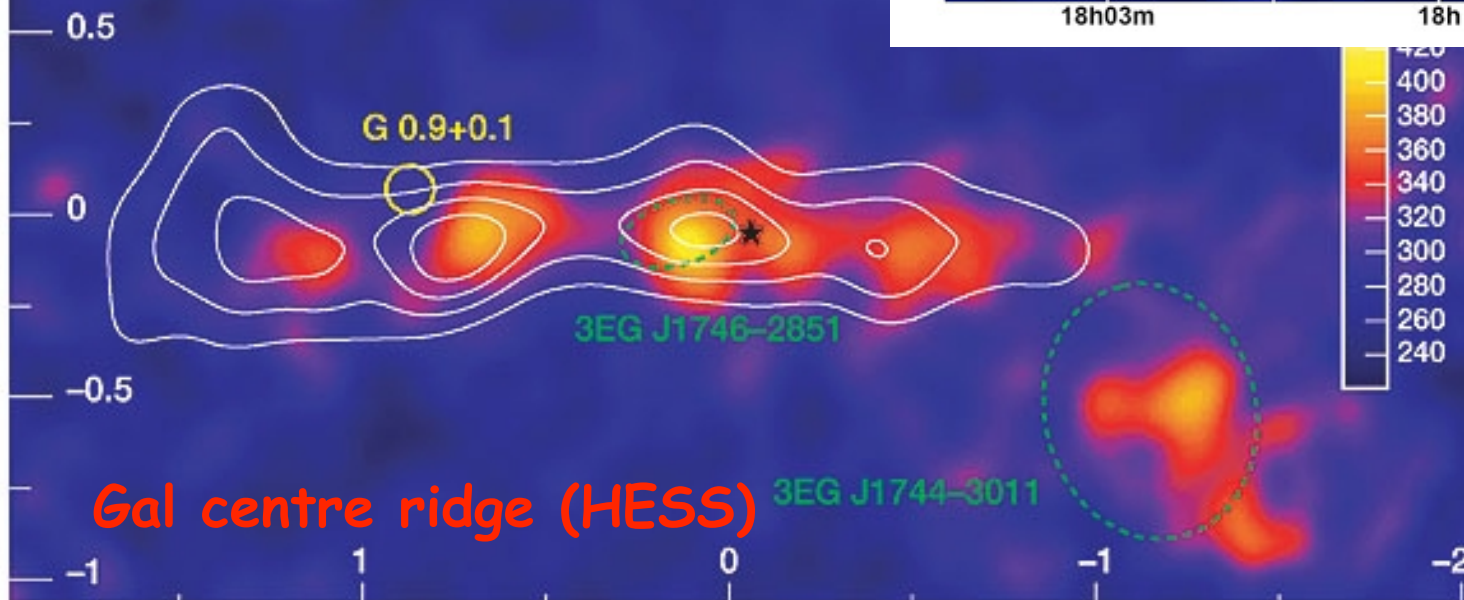
# Molecular clouds are gamma ray sources

Okumura et al, 2009



Aharonian et al, 2008

Aharonian et al, 2006



**(1) Molecular clouds  
as cosmic ray barometers**

# Molecular Clouds as CR barometers

(Issa & Wolfendale, 1981 ; Aharonian, 1991)

**Zero-th order approximation:** the CR spectrum everywhere in the Galaxy is identical to the spectrum we observe at Earth

$$\Rightarrow F_{\gamma} = A \left( \frac{M_{cl}}{d^2} \right)$$

known constant

# Molecular Clouds as CR barometers

$$F_{\gamma} = A \left( \frac{M_{cl}}{d^2} \right)$$

Conversely, if we know  $M_{cl}$  and  $d$  (from CO measurements) we can derive  $A$  and estimate both the normalization and spectrum of CRs at the cloud -> **Molecular Clouds are CR Barometers**

## Two caveats:

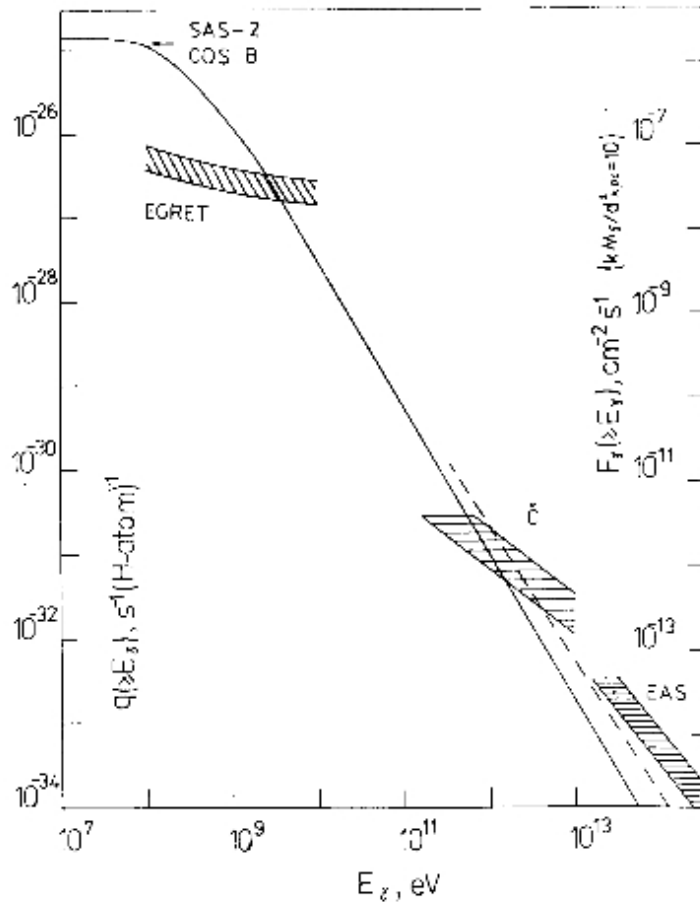
- error in the determination of the mass (CO -> H<sub>2</sub> conversion)
- effective penetration of CR into the cloud (if not see Gabici et al. 2007)

# Molecular Clouds as CR barometers: GeVs

detectable with EGRET if:

$$\frac{M_5}{d_{kpc}^2} > 10$$

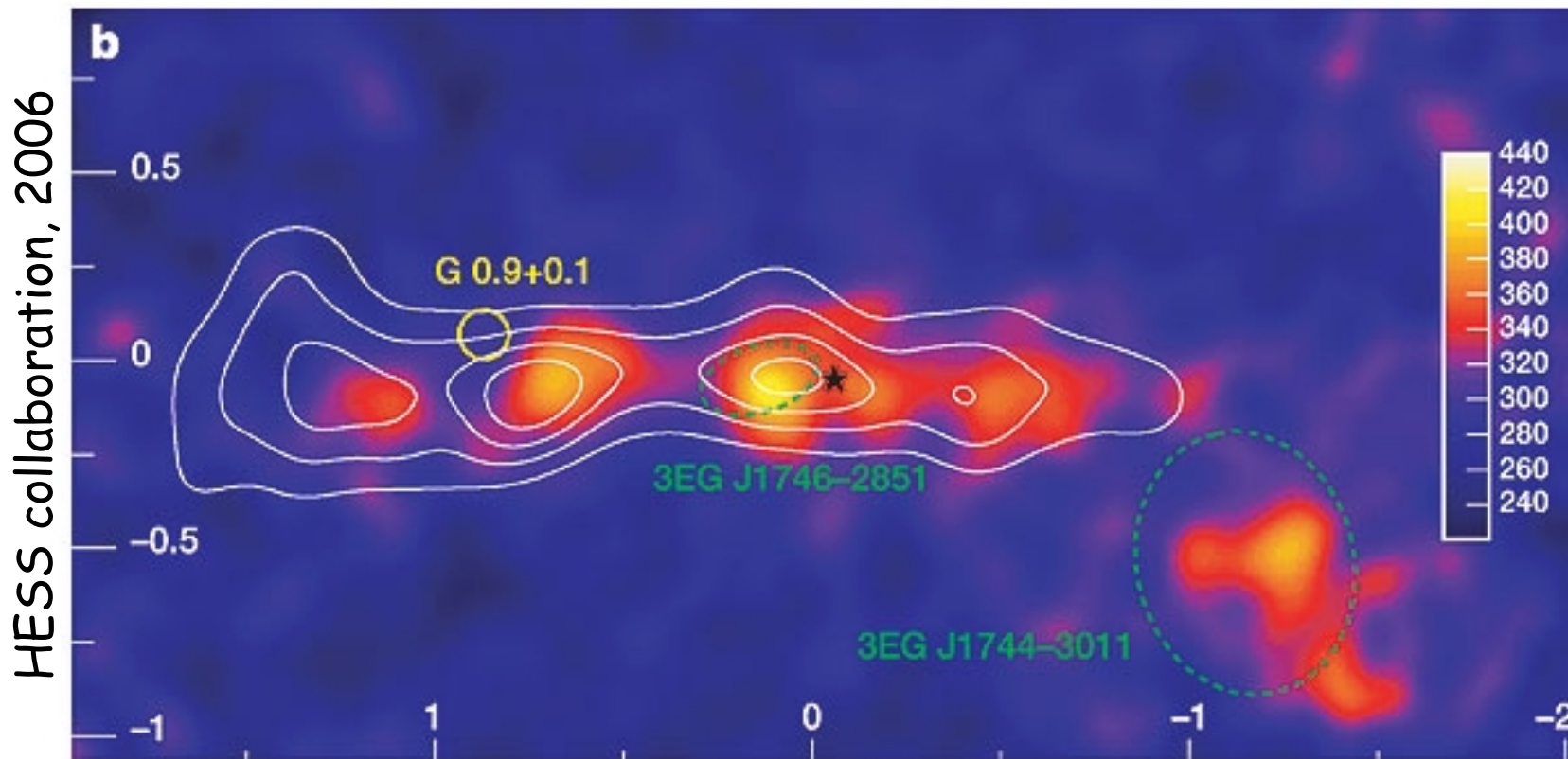
Akharonian, 1991



- ☀ only a few (Orion, Monoceros) -> Digel et al.2001
- ☀ we need **FERMI**
- ☀ gamma ray spectrum -> steep power law from GeV to TeV energies

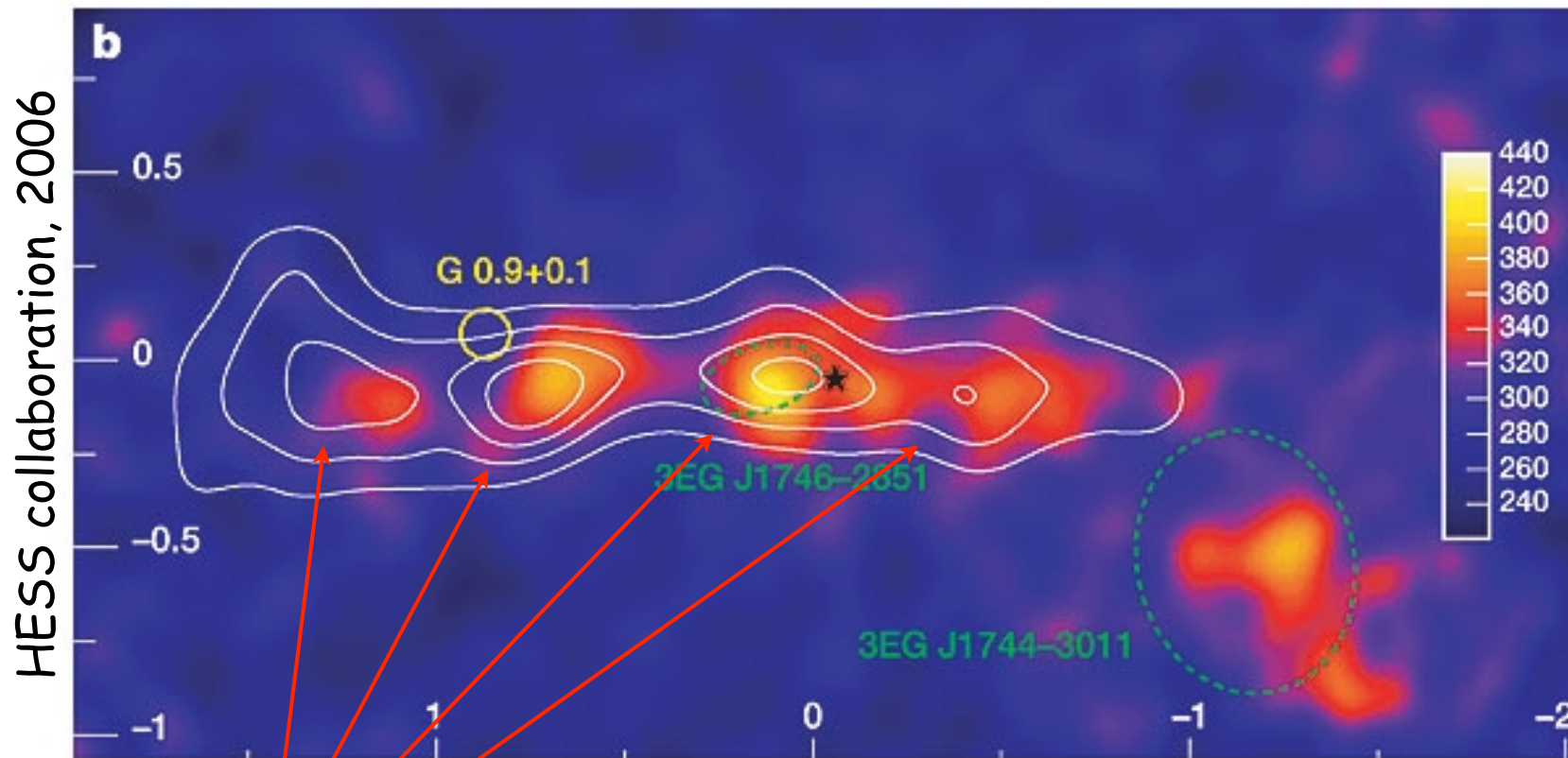
# Molecular Clouds as CR barometers: TeVs

The galactic centre ridge as seen by HESS



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The galactic centre ridge as seen by HESS

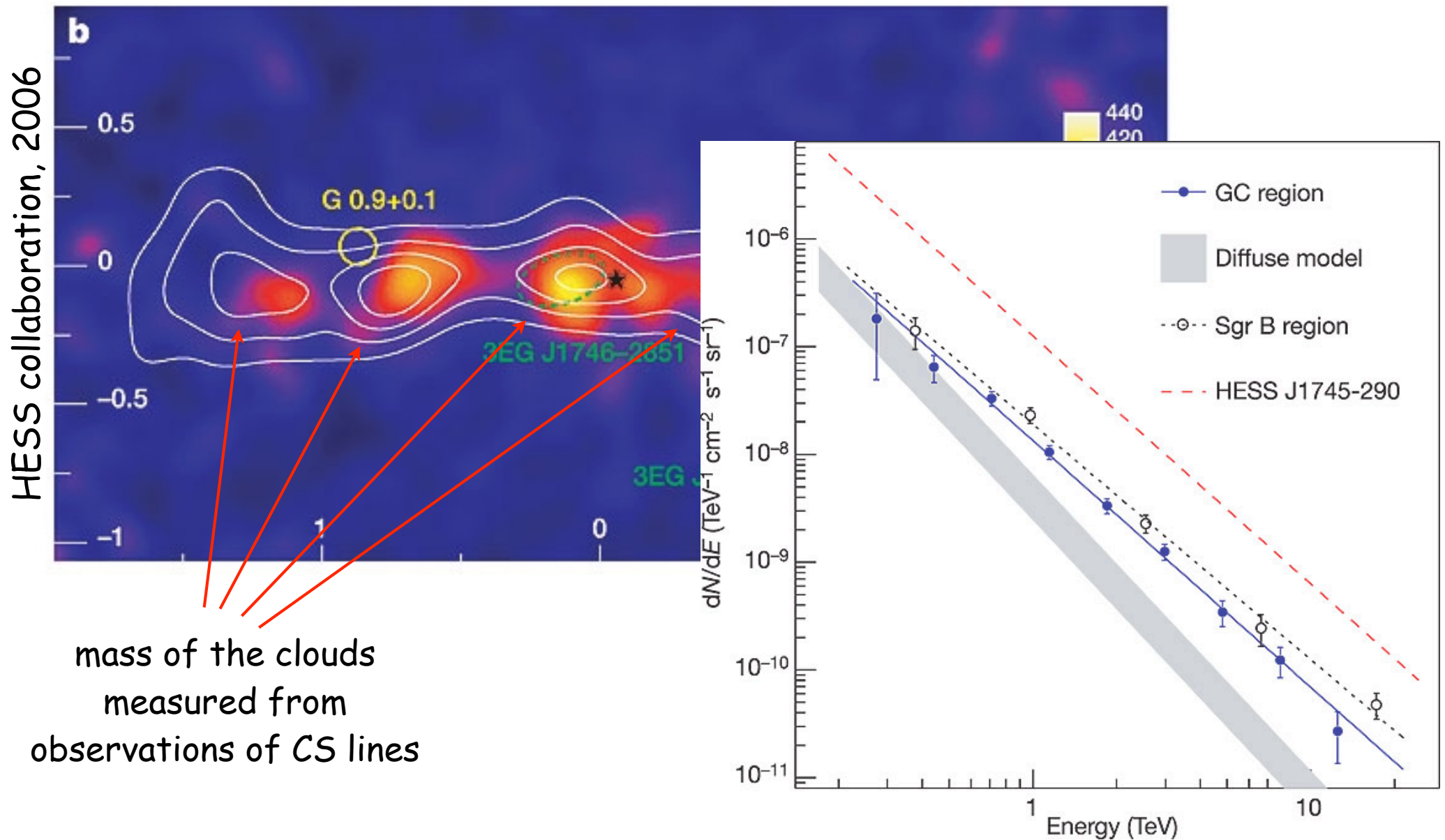


mass of the clouds  
measured from  
observations of CS lines



# Molecular Clouds as CR barometers: TeVs

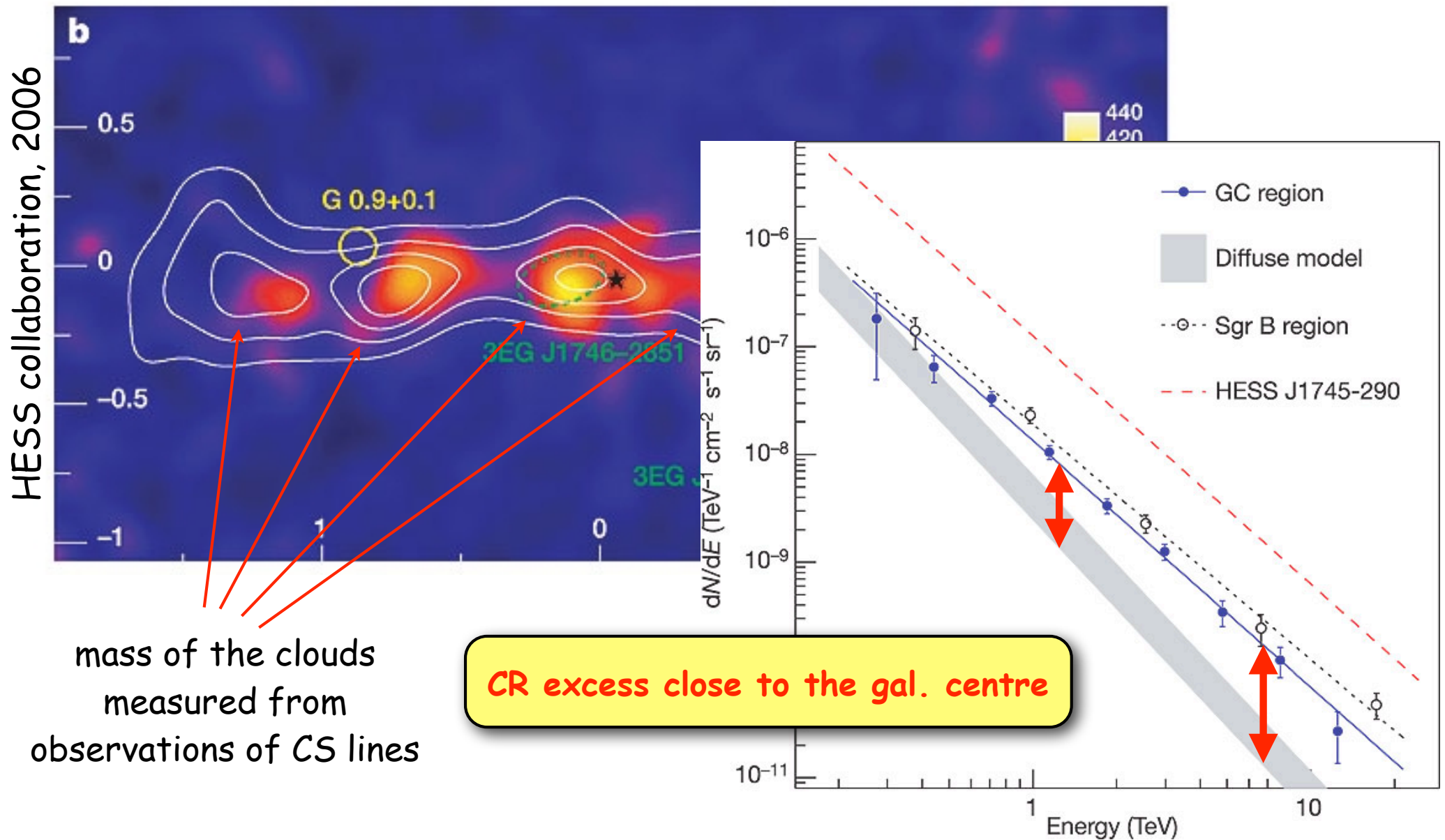
The galactic centre ridge as seen by HESS





# Molecular Clouds as CR barometers: TeVs

The galactic centre ridge as seen by HESS



# Detectability at TeV energies: the role of CTA

Gamma-ray flux from the cloud @1TeV

$$2 \times 10^{-13} \delta \frac{M_5}{d_{kpc}^2} \text{ TeV/cm}^2/\text{s} > 10^{-14} \left( \frac{\epsilon_{CTA}}{0.1} \right) \left( \frac{\theta}{0.1^\circ} \right) \text{ TeV/cm}^2/\text{s}$$

flux from a  
passive cloud

enhancement with  
respect to passive  
cloud

mass and distance  
of the cloud

# Detectability at TeV energies: the role of CTA

Sensitivity of CTA @1TeV

$$2 \times 10^{-13} \delta \frac{M_5}{d_{kpc}^2} \text{TeV/cm}^2/\text{s} > 10^{-14} \left( \frac{\epsilon_{CTA}}{0.1} \right) \left( \frac{\theta}{0.1^\circ} \right) \text{TeV/cm}^2/\text{s}$$

HESS sensitivity  
divided by 10

how much CTA is  
better than HESS

angular resolution

# Detectability at TeV energies: the role of CTA

Simplifying assumption:

$$2 \times 10^{-13} \delta \frac{M_5}{d_{kpc}^2} \text{ TeV/cm}^2/\text{s} > 10^{-14} \left( \frac{\epsilon_{CTA}}{0.1} \right) \left( \frac{\theta}{0.1^\circ} \right) \text{ TeV/cm}^2/\text{s}$$

all the clouds have the same density ( $\sim 100 \text{ cm}^{-3}$ ):  $\theta \approx 1^\circ \frac{M_5^{1/3}}{d_{kpc}}$

# Detectability at TeV energies: the role of CTA

Detectability condition:

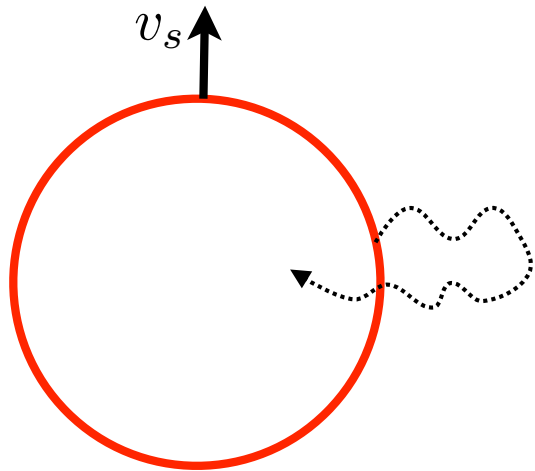
$$d_{kpc} < 2 \delta M_5^{2/3}$$

- **HESS** cannot detect passive clouds
- **CTA** will be able to detect local passive clouds ( $\sim$  kpc distance scale)
- **CTA (HESS)** will probe the Cosmic Ray pressure in regions of the Galaxy

where  $\delta \gg 1$  ( $\delta \gg 10$ )

**(2) Illuminated molecular clouds  
(MC/SNR associations)**

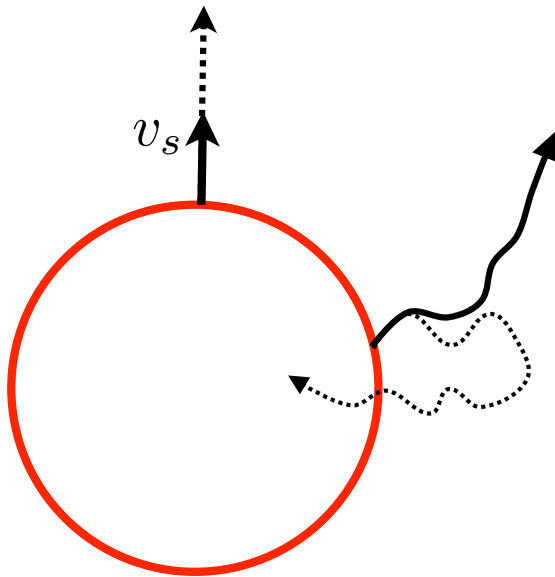
# How to use Molecular Clouds...



SNRs accelerate CRs

# How to use Molecular Clouds...

CRs "somehow" escape the SNR

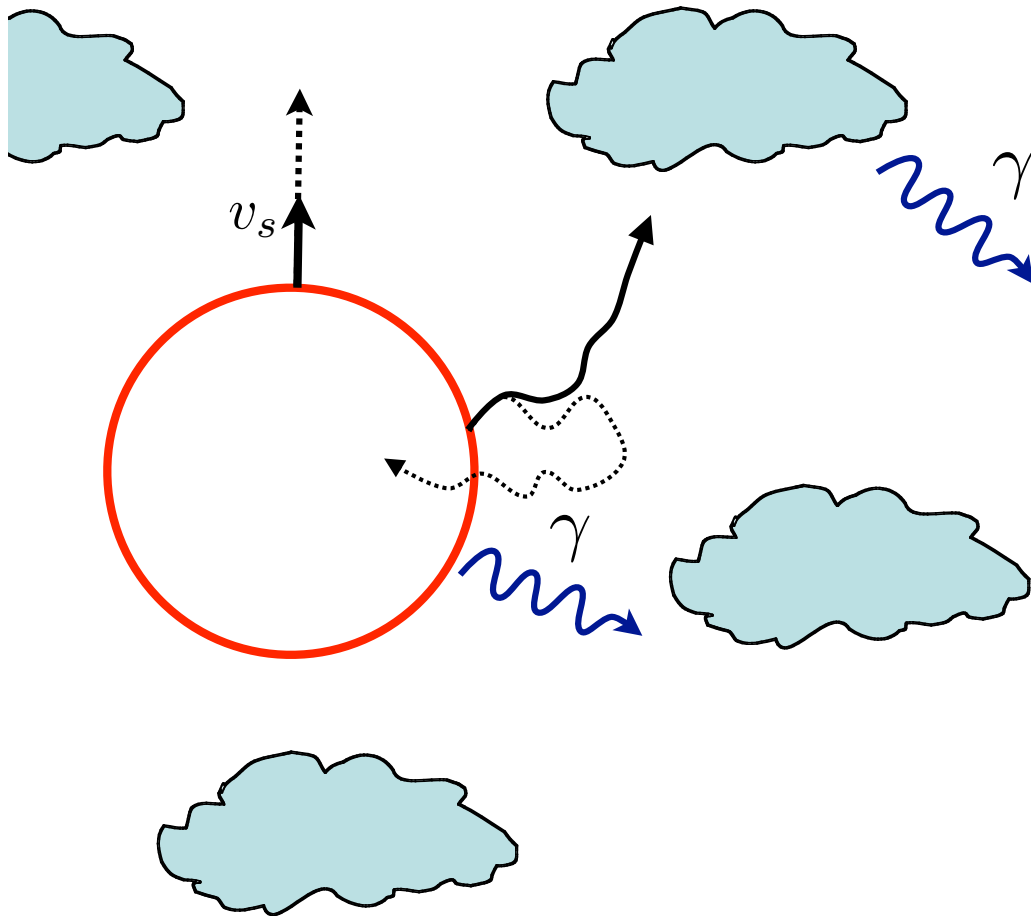


SNRs accelerate CRs



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CRs "somehow" escape the SNR



MCs enhance the gamma ray emission

SNRs accelerate CRs

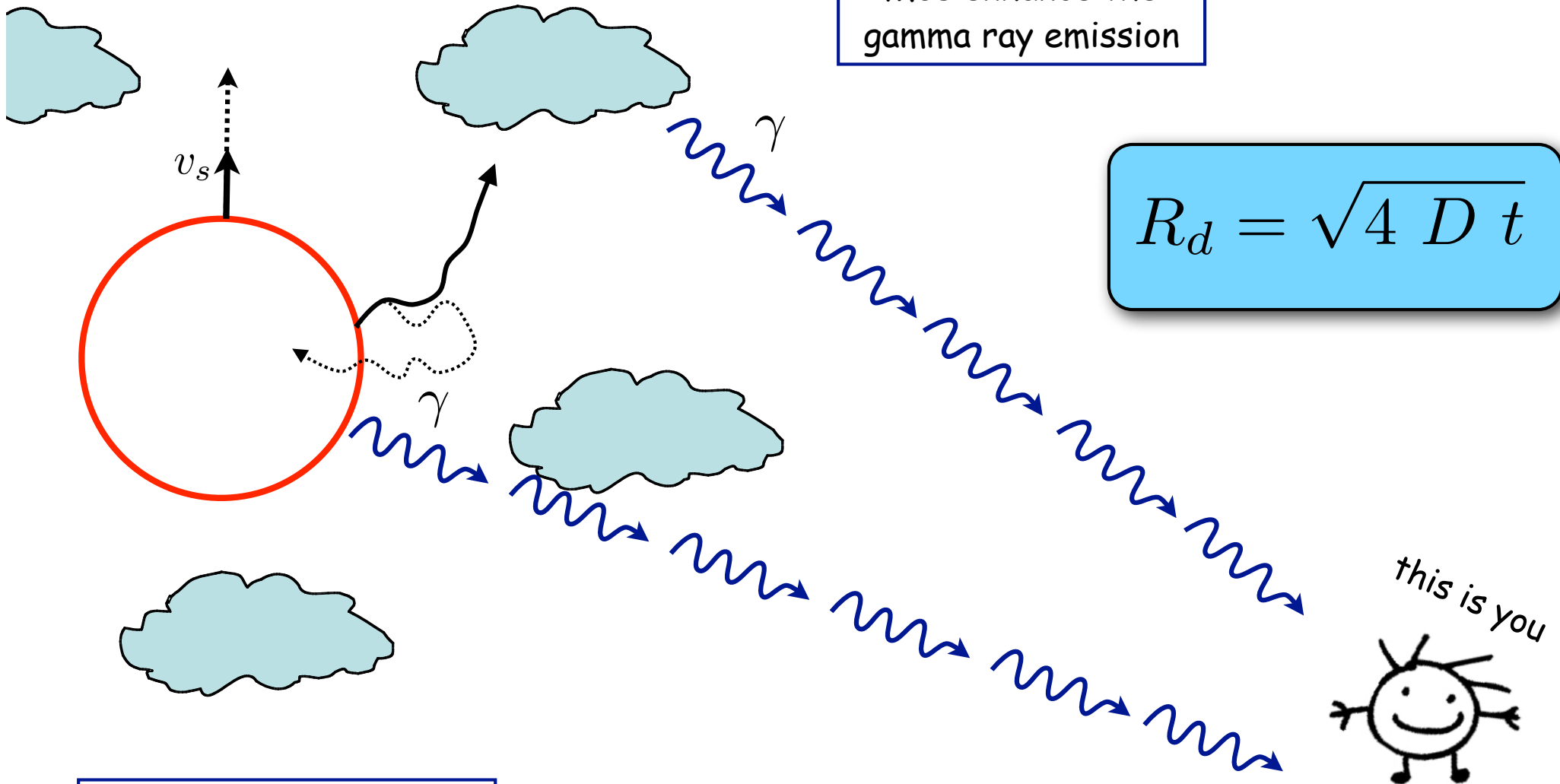
# How to use Molecular Clouds...

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$$R_d = \sqrt{4 D t}$$

SNRs accelerate CRs



# How to use Molecular Clouds...

CRs "somehow" escape the SNR

MCs enhance the gamma ray emission

$$R_d = \sqrt{4 D t}$$

**We can try to:**

- constrain diffusion
- identify the sources of CRs

SNRs accelerate CRs

this is you



# The diffusion of CRs

Spallation measurements tell us that cosmic rays follow tortuous paths before escaping the Galaxy. **Why?**

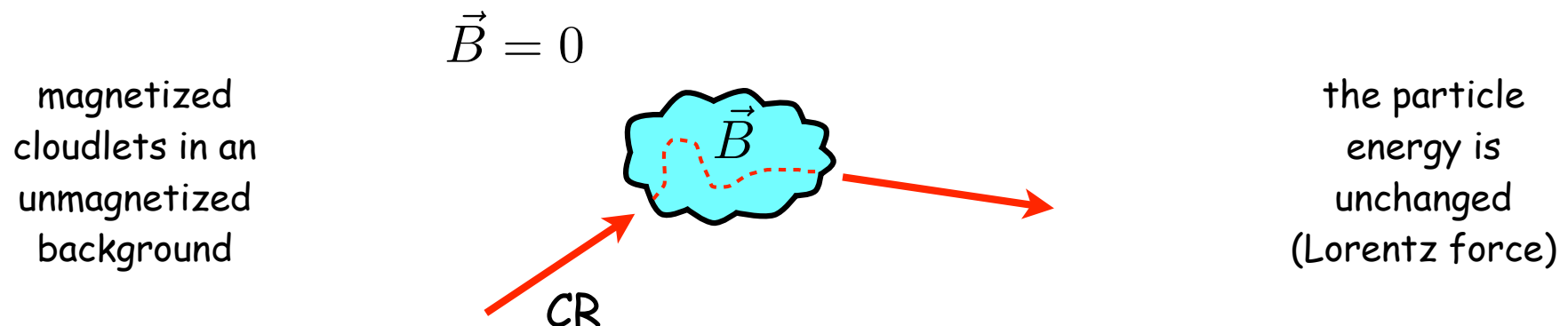
The galactic magnetic field or, better, **irregularities in the Galactic magnetic field** are responsible for the diffusive propagation of cosmic rays.

# The diffusion of CRs

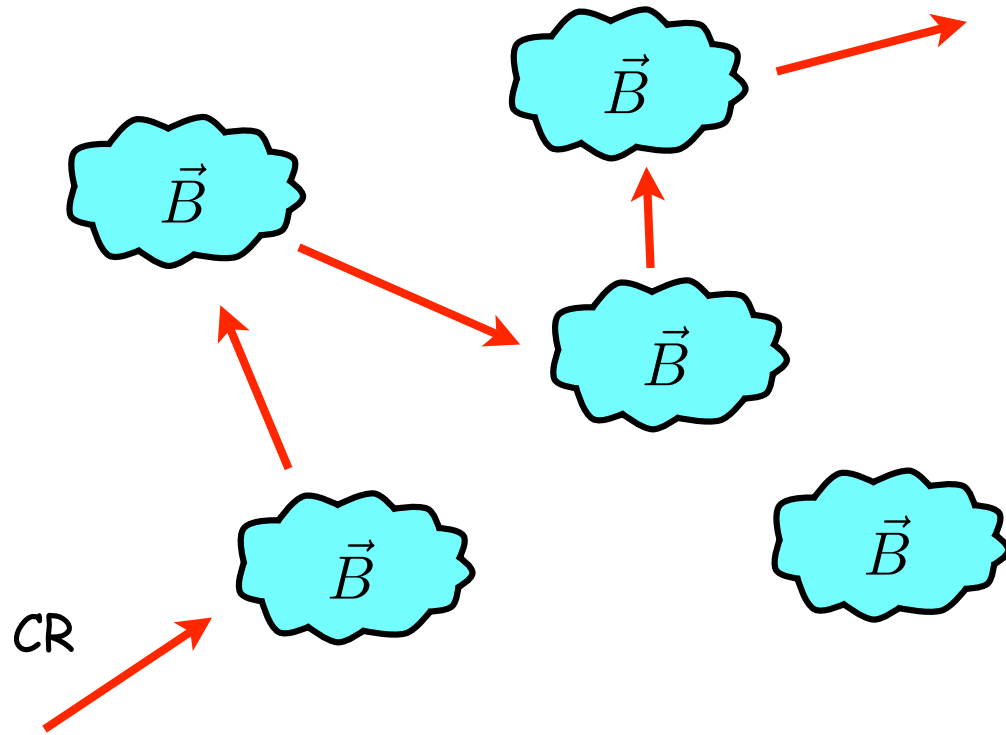
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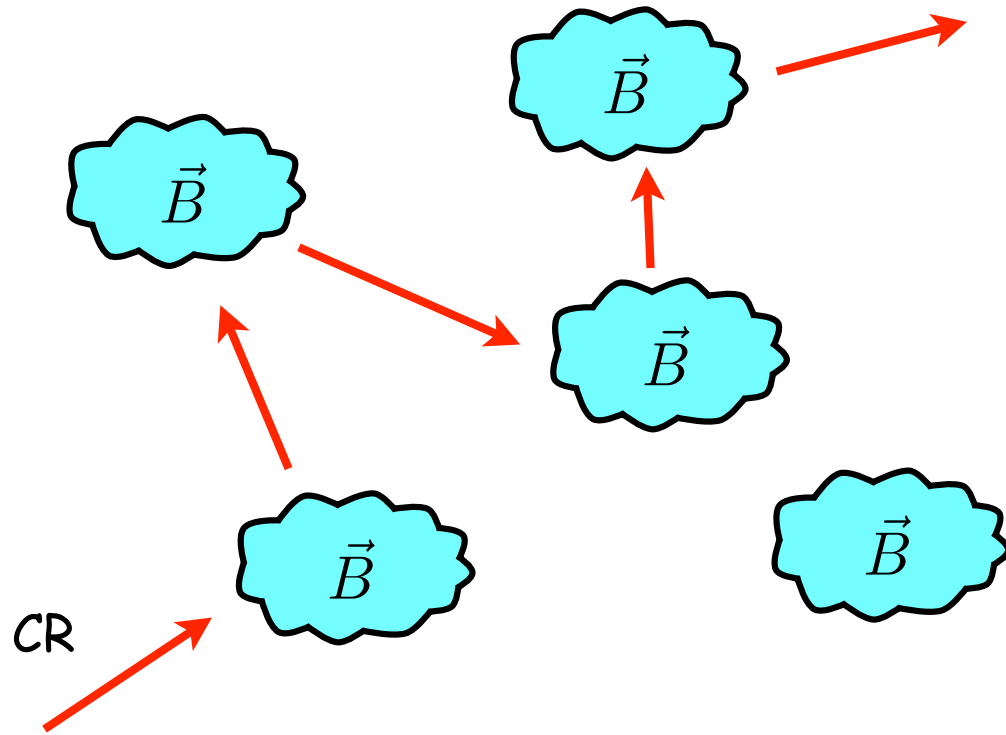
(Oversimplified picture)



# The diffusion of CRs



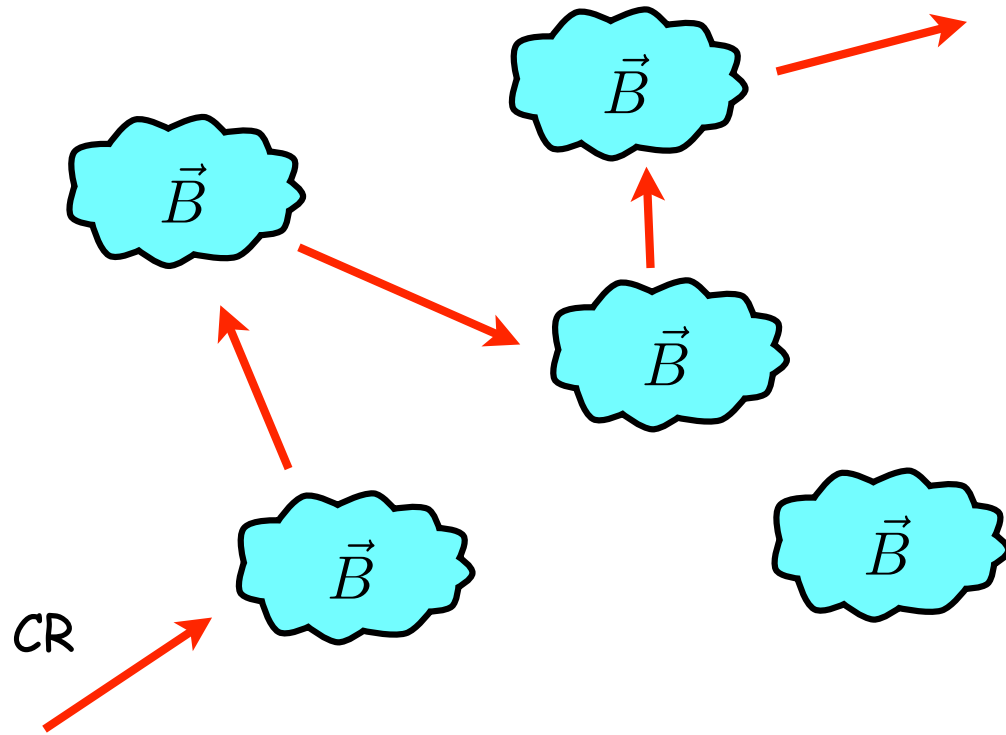
# The diffusion of CRs



$\lambda$   $\rightarrow$  mean free path



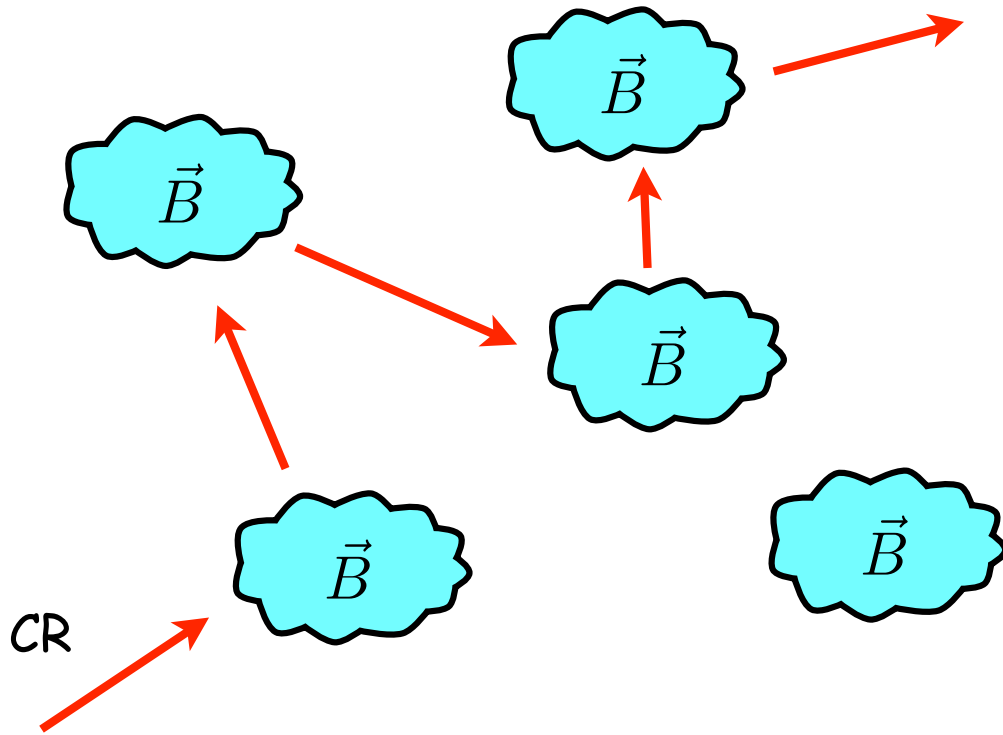
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$$\tau_c = \frac{\lambda}{c} \rightarrow \text{collision time}$$

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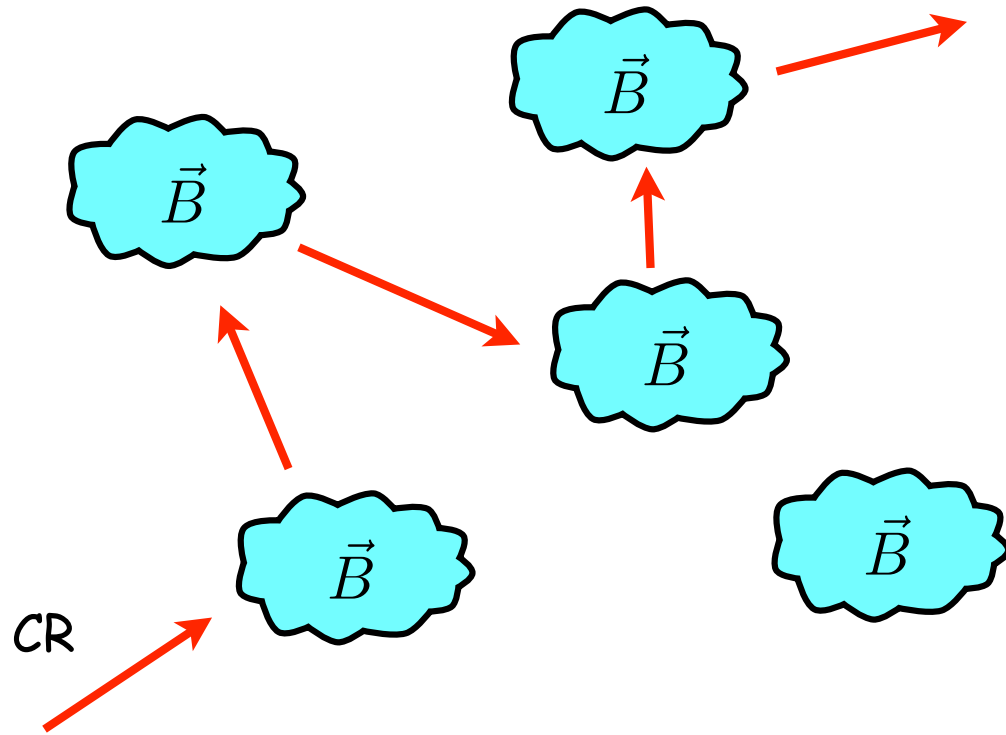


$\lambda$  -> mean free path

$$\tau_c = \frac{\lambda}{c} \rightarrow \text{collision time}$$

$$N = \frac{t}{\tau_c} \rightarrow \text{\# collisions after time } t$$

# The diffusion of CRs



$\lambda$  -> mean free path

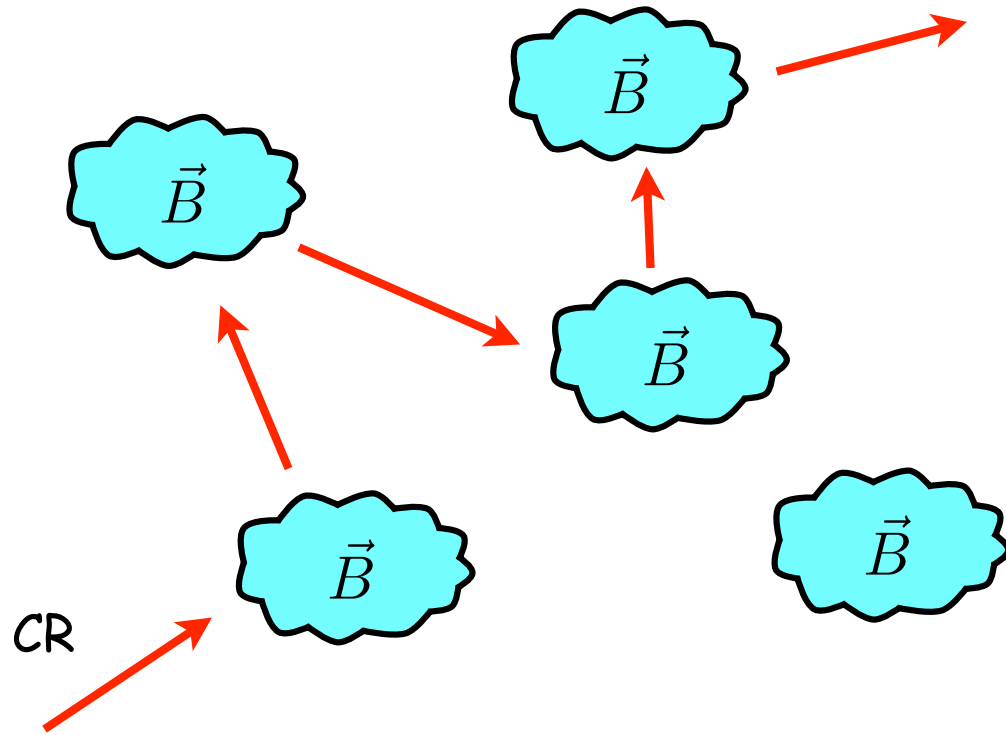
$\tau_c = \frac{\lambda}{c}$  -> collision time

$N = \frac{t}{\tau_c}$  -> # collisions after time t

diffusion length ->  $l_d = \lambda \sqrt{N}$

random walk

# The diffusion of CRs



$\lambda$  → mean free path

$\tau_c = \frac{\lambda}{c}$  → collision time

$N = \frac{t}{\tau_c}$  → # collisions after time  $t$

diffusion length →

$$l_d = \lambda \sqrt{N} = \lambda \sqrt{\frac{t}{\tau_c}} = \lambda \sqrt{\frac{t c}{\lambda}} = \sqrt{\lambda c t}$$

random walk

this product determines the diffusion properties of the particle

# The diffusion of CRs

It is convenient to define the quantity  $D = \lambda c$  called **diffusion coefficient**

diffusive propagation ->  $l_d = \sqrt{D t} \propto \sqrt{t}$

straight line propagation ->  $l_{sl} = c t \propto t$

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Spallation measurements allow us to measure the average diffusion coefficient in the Galaxy

$$l_{disk} = \sqrt{D t_{disk}} \longrightarrow D = \frac{l_{disk}^2}{t_{disk}} = 10^{28} \text{ cm}^2/\text{s}$$

$\nearrow$  ~300 pc       $\nearrow$  3 Myr (from spallation)       $\nearrow$  @ 10 GeV

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$\sim 300 \text{ pc}$  (pointing to  $l_{disk}$ )      3 Myr (from spallation) (pointing to  $t_{disk}$ )      @ 10 GeV (pointing to  $10^{28}$ )

**Energy dependent ->  $D \propto E^{0.6}$**



# Implications of the SNR hypothesis

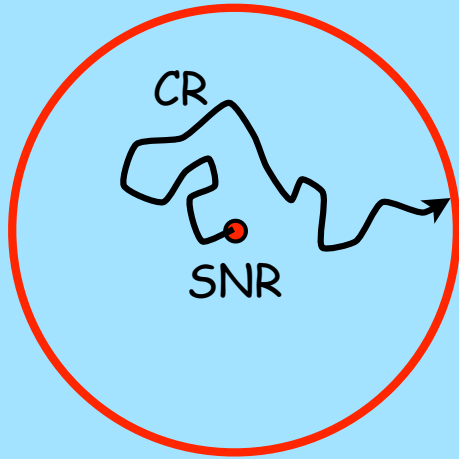
CR sea  $\rightarrow 1 \text{ eV/cm}^3$

●  
SNR

$$E_{CR}^{SNR} = 10^{50} \text{ erg}$$

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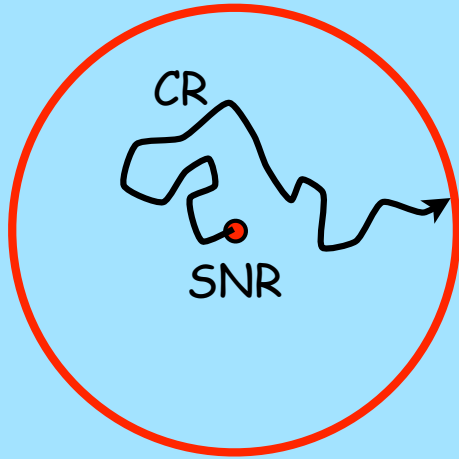
volume affected by CRs from the SNR

$$\frac{E_{CR}^{SNR}}{\left(\frac{4\pi}{3} R_{CR}^3\right)} = 1 \text{ eV/cm}^3$$

$\Rightarrow R_{CR} \approx 100 \text{ pc}$

# Implications of the SNR hypothesis

CR sea  $\rightarrow 1 \text{ eV/cm}^3$



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volume affected by CRs from the SNR

$$\frac{E_{CR}^{SNR}}{\left(\frac{4\pi}{3} R_{CR}^3\right)} = 1 \text{ eV/cm}^3$$

$\Rightarrow R_{CR} \approx 100 \text{ pc}$

such a volume is affected for a time:

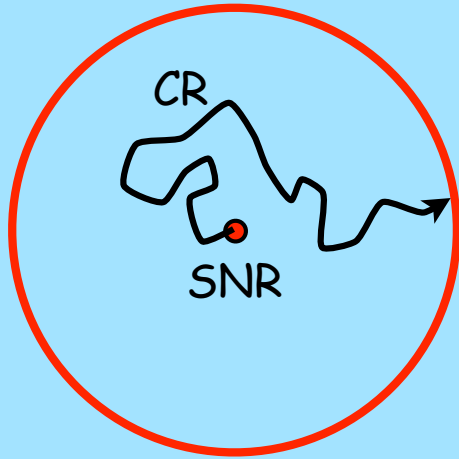
$$D = 10^{28} \left(\frac{E}{10 \text{ GeV}}\right)^{0.6} \text{ cm}^2/\text{s} \quad \Rightarrow \quad D(1 \text{ TeV}) \approx 2 \times 10^{29} \text{ cm}^2/\text{s}$$

$$t \approx \frac{R_{CR}^2}{D} \approx 10^4 \text{ yr}$$

# Implications of the SNR hypothesis

Very rough!

CR sea  $\rightarrow 1 \text{ eV/cm}^3$



$$E_{CR}^{SNR} = 10^{50} \text{ erg}$$

volume affected by CRs from the SNR

$$\frac{E_{CR}^{SNR}}{\left(\frac{4\pi}{3} R_{CR}^3\right)} = 1 \text{ eV/cm}^3$$

$$\Rightarrow R_{CR} \approx 100 \text{ pc}$$

such a volume is affected for a time:

$$D = 10^{28} \left(\frac{E}{10 \text{ GeV}}\right)^{0.6} \text{ cm}^2/\text{s} \Rightarrow D(1 \text{ TeV}) \approx 2 \times 10^{29} \text{ cm}^2/\text{s}$$

$$t \approx \frac{R_{CR}^2}{D} \approx 10^4 \text{ yr}$$

# Montmerle's SNOBs

adapted from Montmerle, 1979 ; Casse & Paul, 1980

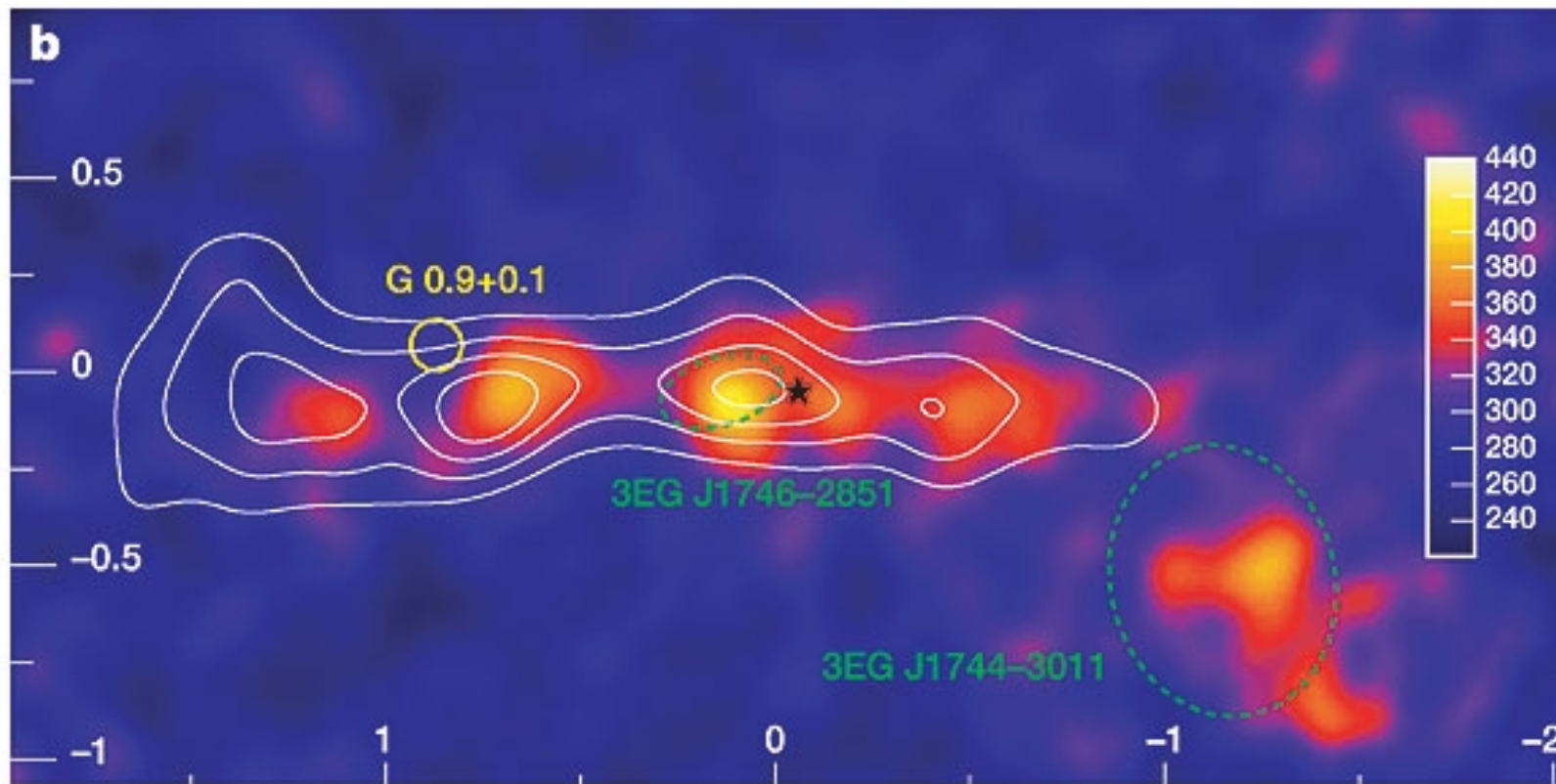
Super Novae-OB  
stars associations

- ✓ Massive (OB) stars form in dense regions -> molecular cloud complexes
- ✓ OB stars evolve rapidly and eventually explode forming SNRs
- ✓ SNR shocks accelerate COSMIC RAYS
- ✓ CRs escape from their sources and diffuse away in the DENSE circumstellar material -> molecular cloud complex
- ✓ ...and produce there gamma rays!

An association between cosmic ray sources and molecular cloud is expected

# Example: the galactic centre ridge

The galactic centre ridge as seen by HESS

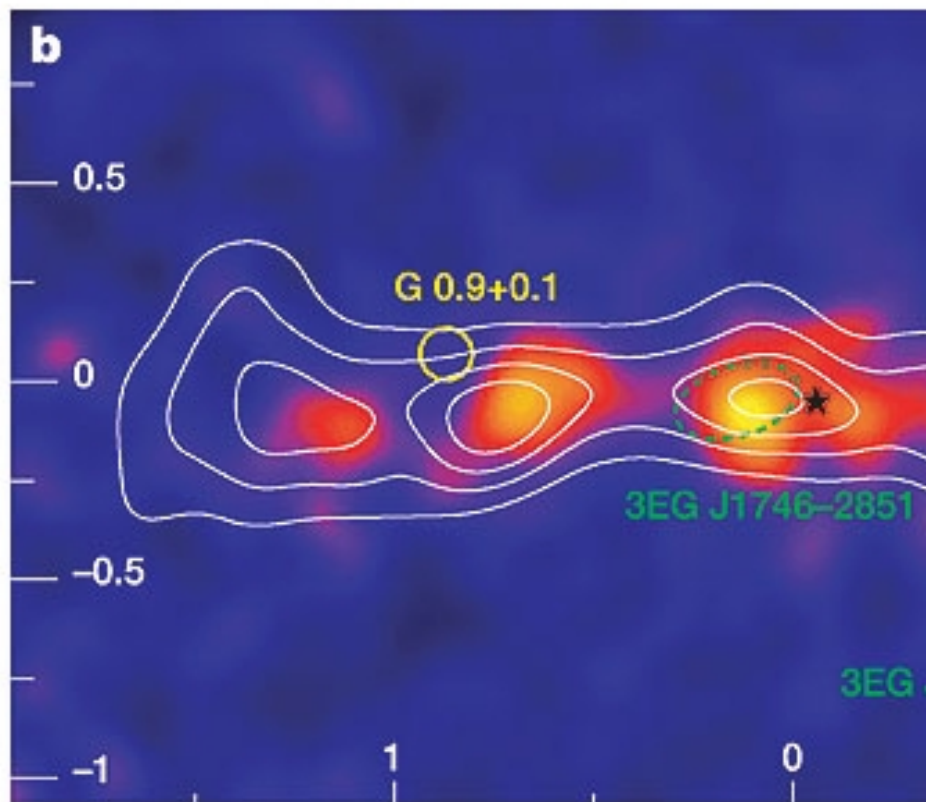


HESS collaboration, 2006

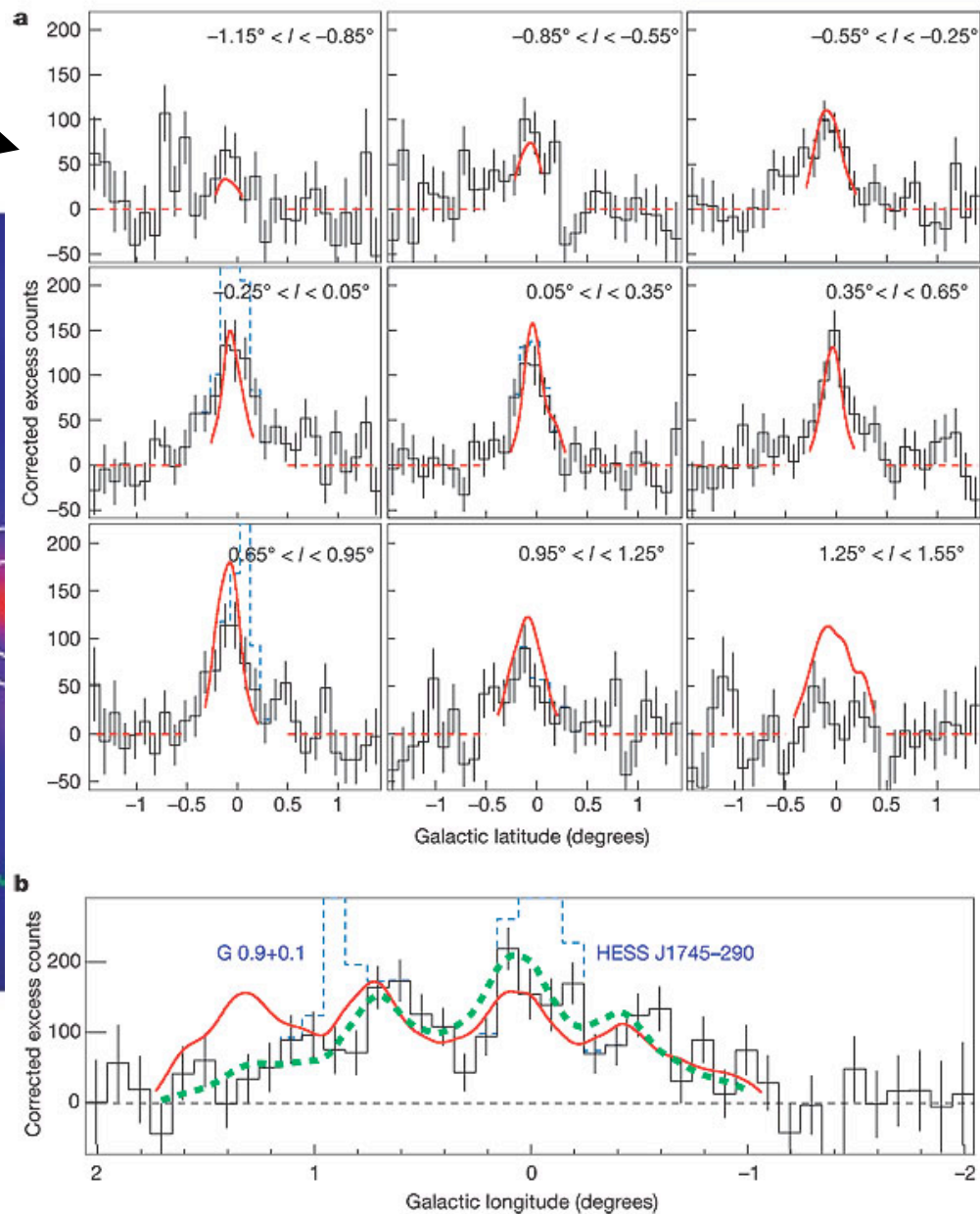
# Example: the galactic centre ridge

The galactic centre ridge as seen by HESS

good match between CS lines and TeV emission

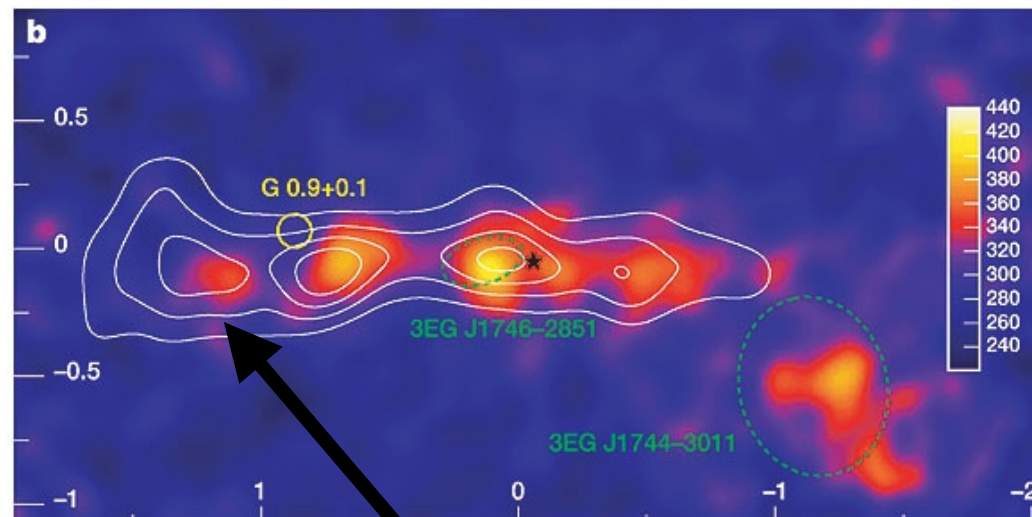
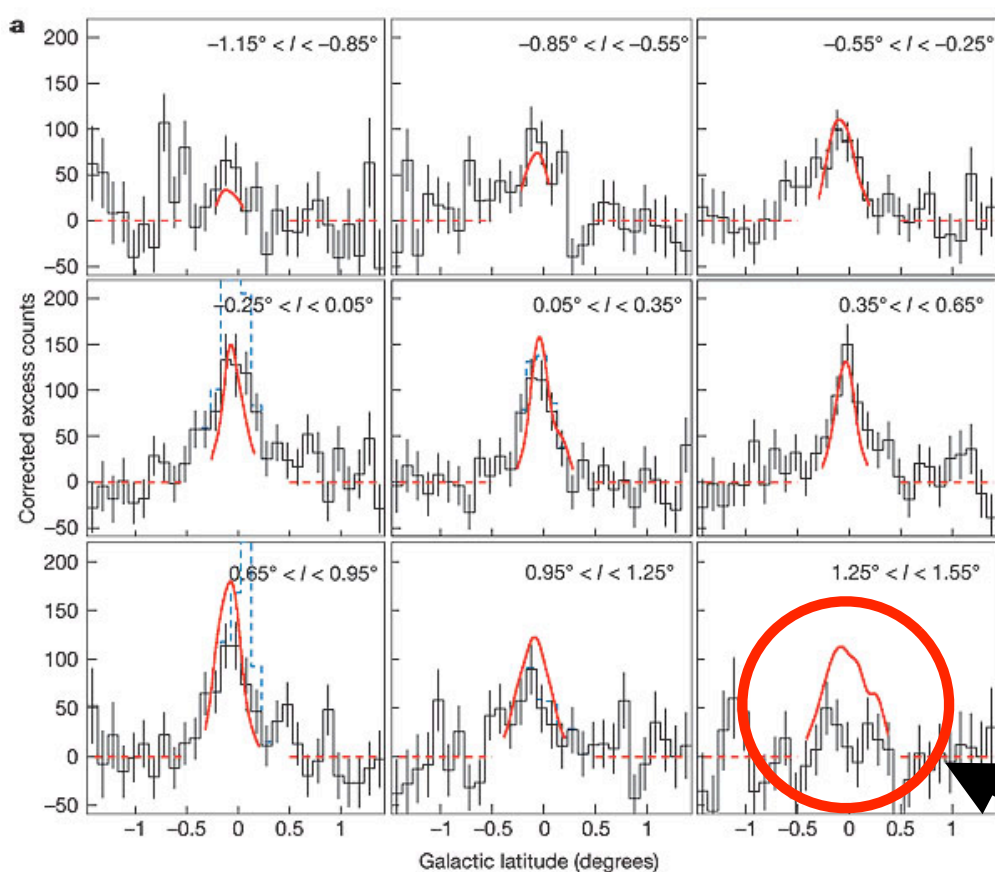


HESS collaboration, 2006

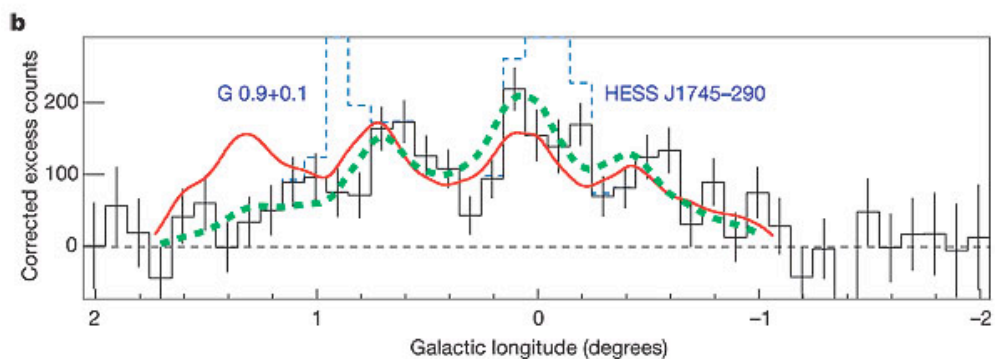




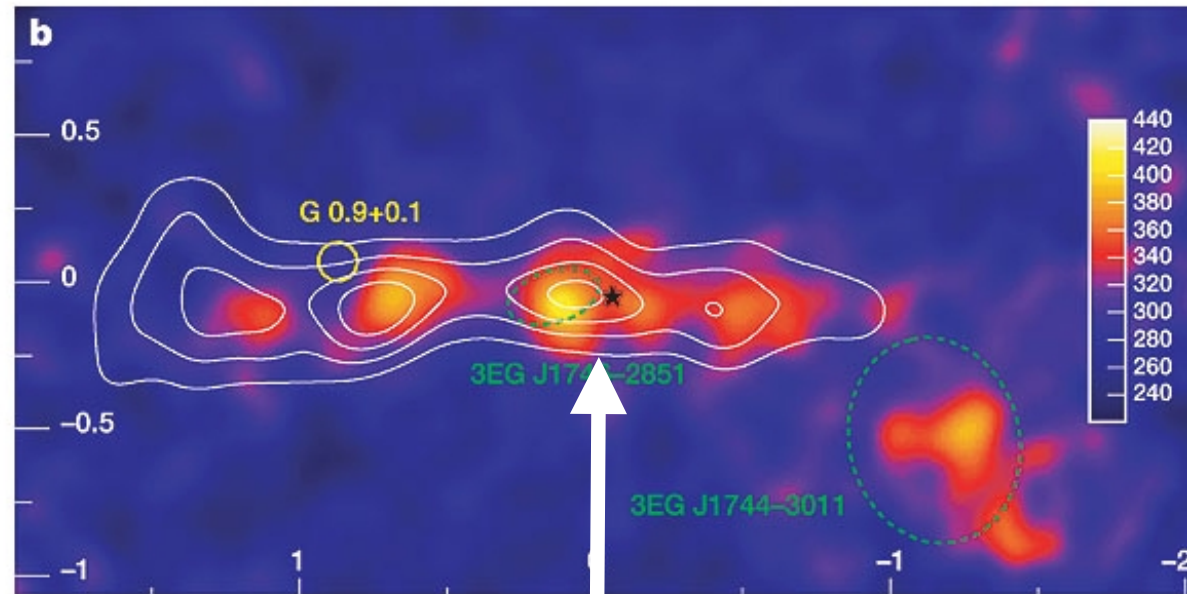
# Example: the galactic centre ridge



the correlation between gamma ray intensity and gas density is worse for the cloud which is the farthest away from the galactic centre



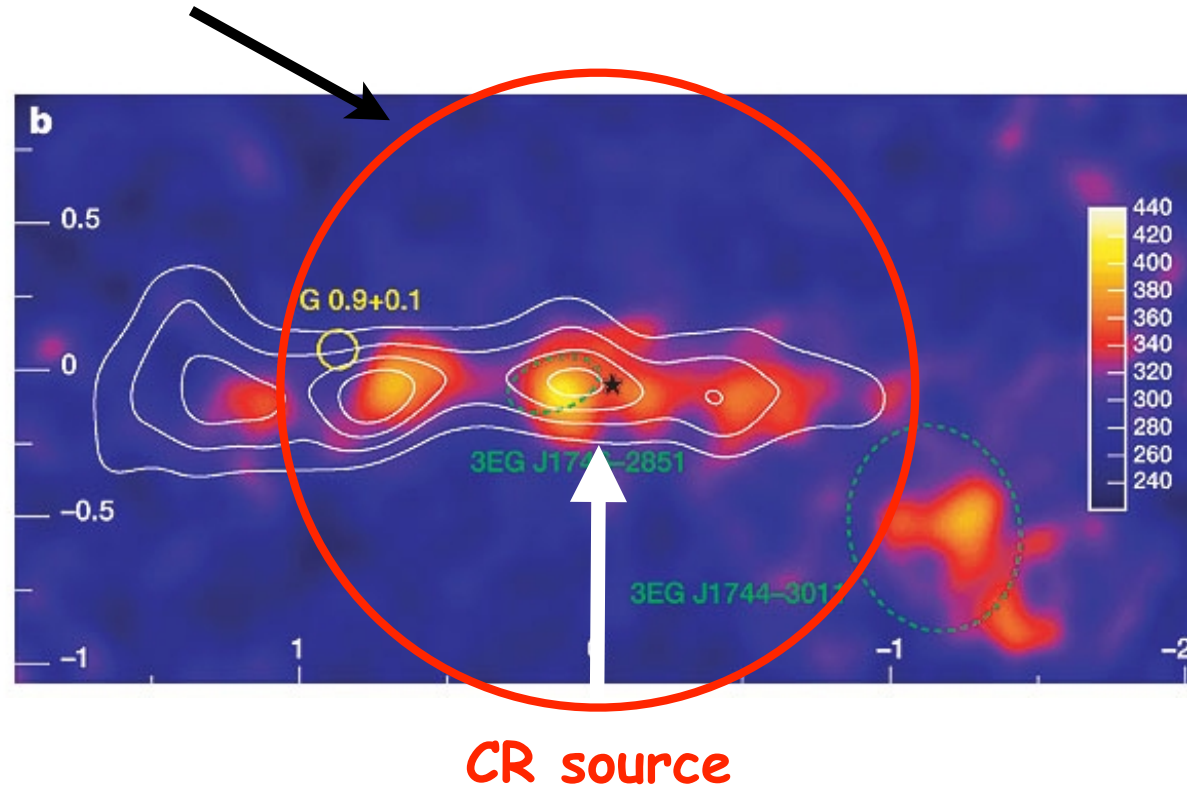
# Example: the galactic centre ridge



CR source

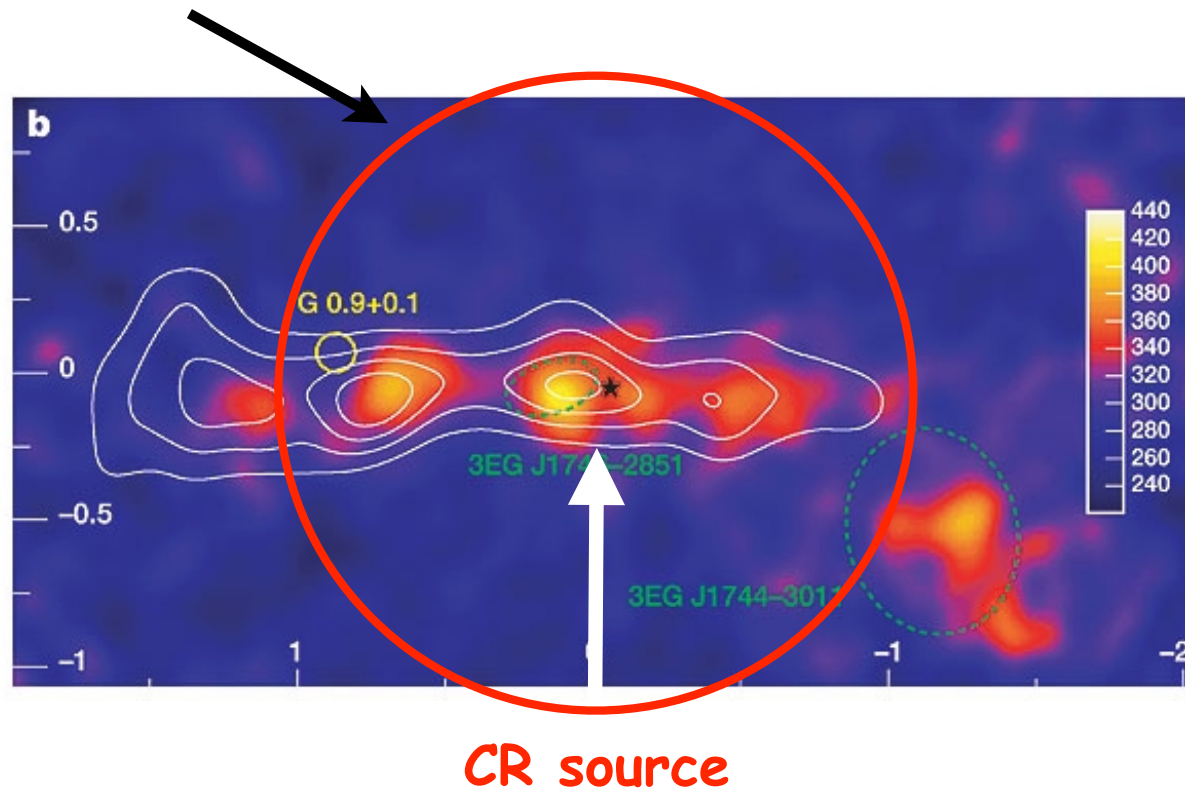
# Example: the galactic centre ridge

after a time  $t_{\text{diff}}$  CRs fill a volume like this



# Example: the galactic centre ridge

after a time  $t_{diff}$  CRs fill a volume like this



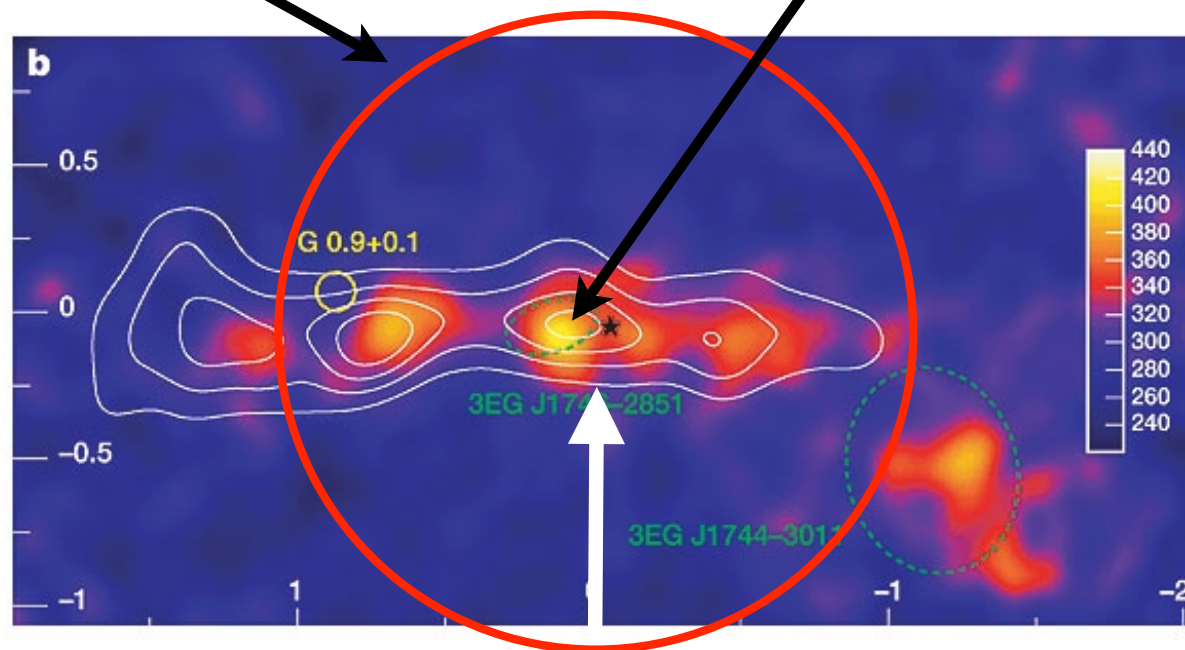
$$t_{diff} \approx \frac{l^2}{D} \quad \rightarrow$$

if we know the age of the source we can estimate the diffusion coefficient!

# Example: the galactic centre ridge

after a time  $t_{diff}$  CRs fill a volume like this

SNR SgrA East  $\rightarrow t \sim 10^4$  yr  
(though quite uncertain)



CR source

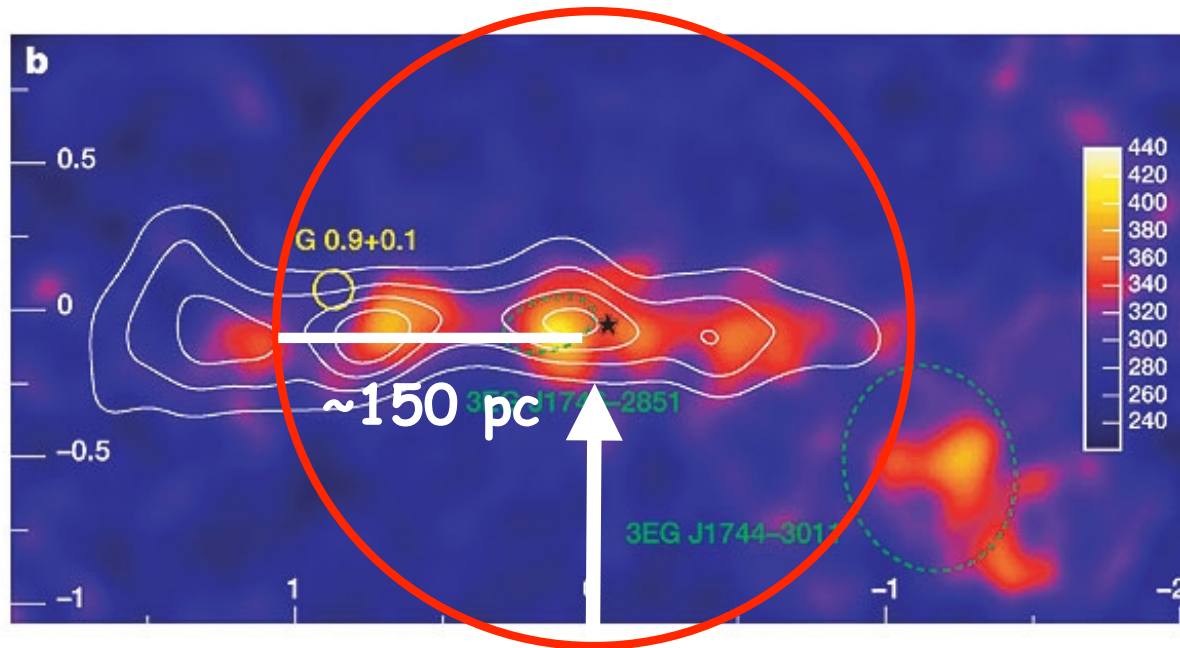
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if we know the age of the source we can estimate the diffusion coefficient!



# Example: the galactic centre ridge

$$t_{diff} \approx \frac{l^2}{D}$$



CR source ( $\sim 10^4$  yr)

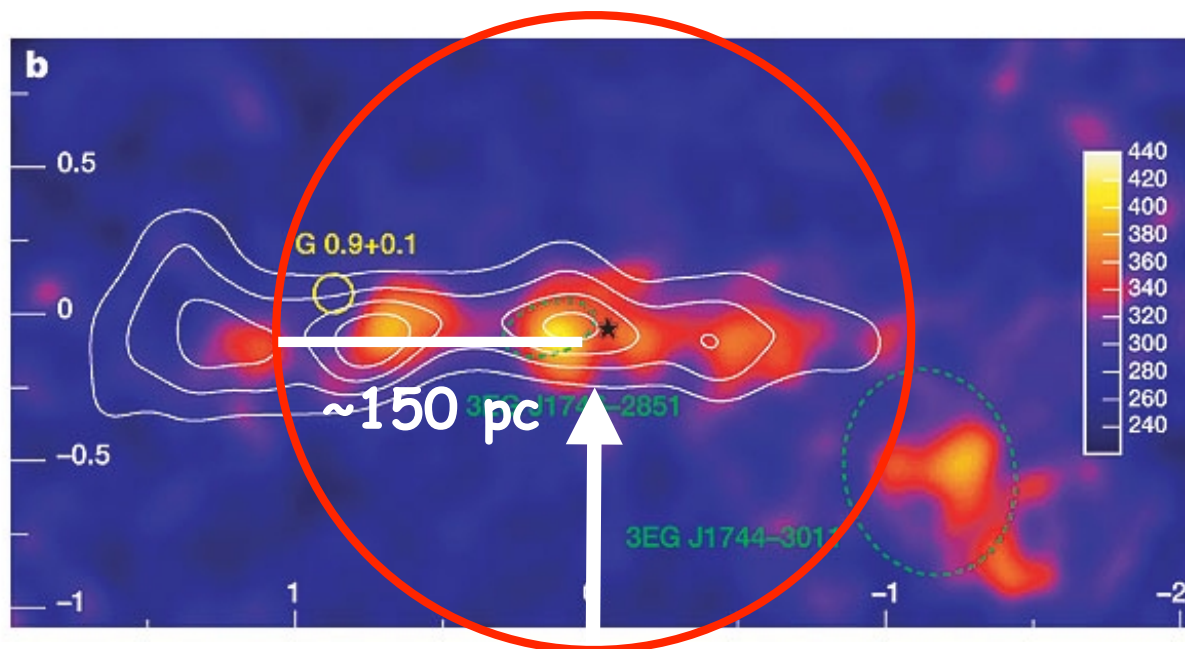
# Example: the galactic centre ridge

$$t_{diff} \approx \frac{l^2}{D}$$



$$D \lesssim 7 \times 10^{29} \text{ cm}^2 / \text{s}$$

not too different from the  
average diffusion  
coefficient in the Galaxy



CR source ( $\sim 10^4$  yr)



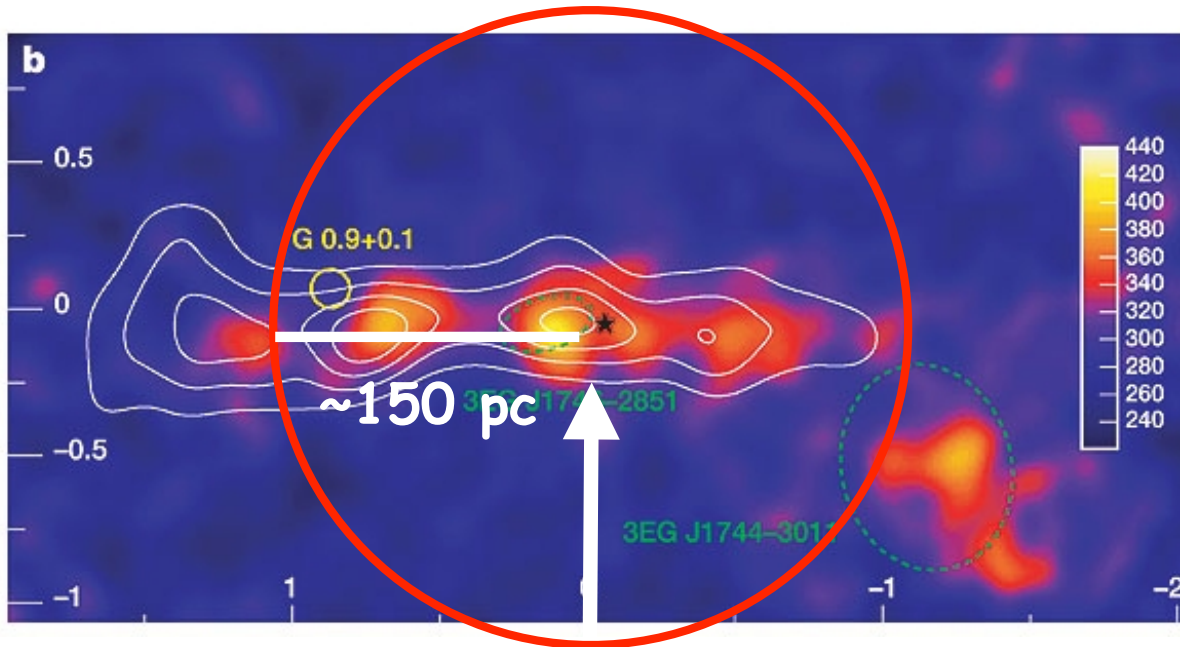
# Example: the galactic centre ridge

$$t_{diff} \approx \frac{l^2}{D}$$



$$D \lesssim 7 \times 10^{29} \text{ cm}^2/\text{s}$$

not too different from the average diffusion coefficient in the Galaxy

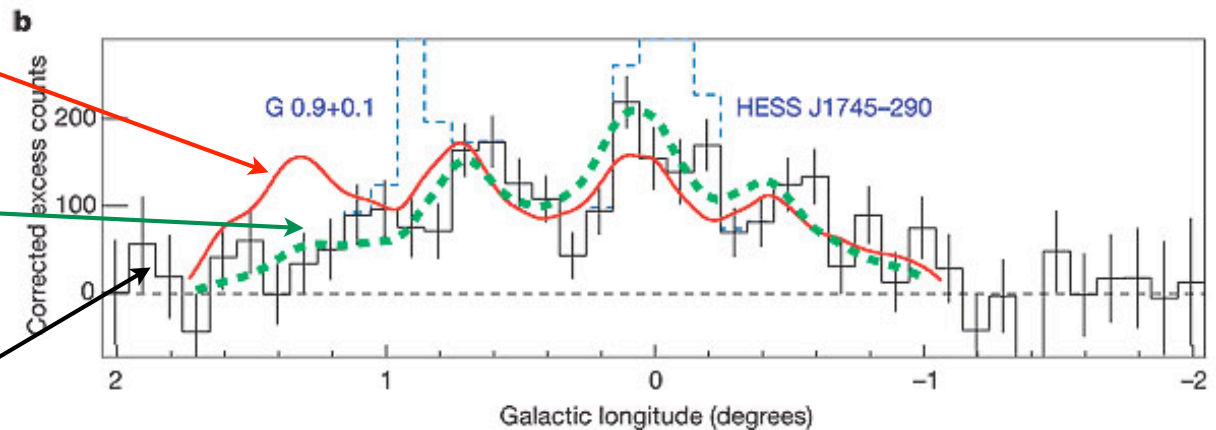


CR source ( $\sim 10^4$  yr)

density

predicted gamma rays  
(diffusion from a source in the Galactic centre)

observed gamma rays



# Take-home messages

- ☑ **Molecular clouds** are massive -> **lots of gammas** expected
- ☑ **MCs** -> **massive stars** -> **SNRs** -> **CR acceleration&escape** -> **gamma rays!**
- ☑ can be used to **probe** the **CR intensity** throughout the Galaxy
- ☑ and/or to identify the **sources of CRs** and **constrain the diffusion coefficient**