Studies of the D0-brane matrix models at low temperatures

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Dublin Institute for Advanced Studies: 13/10/22

Based on: 2110.01312, 2205.06098 & 2210.04881

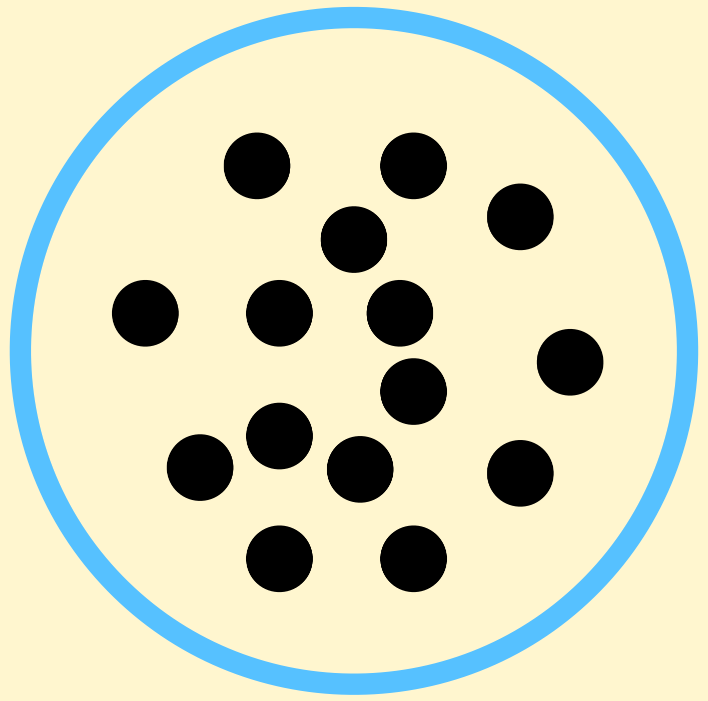
With: Bergner, Bodendorfer, Hanada, Rinaldi, Schäfer, Vranas, Watanabe (MCSMC)

Why low temperatures?

- Analytic gravity duals
- Matrix models → non-commutative → quantum effects
- Quantum traces of gravity?

Plan of the talk

- Definition of the models
- Holography
- Relation with gravity
- Confinement in D0-matrix model
- Simulations, tests at low temperatures
- Comparison with eternal energy of the black zero brane
- Role of gauge constraint?



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D0-matrix model (BFSS)

$$S = \frac{1}{2g_{YM}^2}\int dt Tr\left\{ (D_t X_M)^2 + [X_M, X_N]^2 + i\bar{\psi}^\alpha D_t \psi^\beta + \bar{\psi}^\alpha \gamma^M_{\alpha\beta}[X_M, \psi^\beta] \right\}$$

$X_{N\times N}$: $N \times N$ bosonic hermitian matrices with $M = 1,\cdots,9$

$\psi_{N\times N}$: $N \times N$ fermionic hermitian matrices with $\alpha = 1,\cdots,16$

D_t : $D_t\mathcal{O} = \partial_t\mathcal{O} - i[A_t, \mathcal{O}]$

$\lambda = g_{YM}^2 N = [\text{energy}]^3$

- Dimensional reduction of 4D $\mathcal{N} = 4$ / 10D $\mathcal{N} = 1$
- Matrix regularisation of 11D supermembrane De Wit-Hoppe-Nicolai, 1988
- Matrix model of M-theory (BFSS) Banks-Fischler-Shenker-Susskind, 1996
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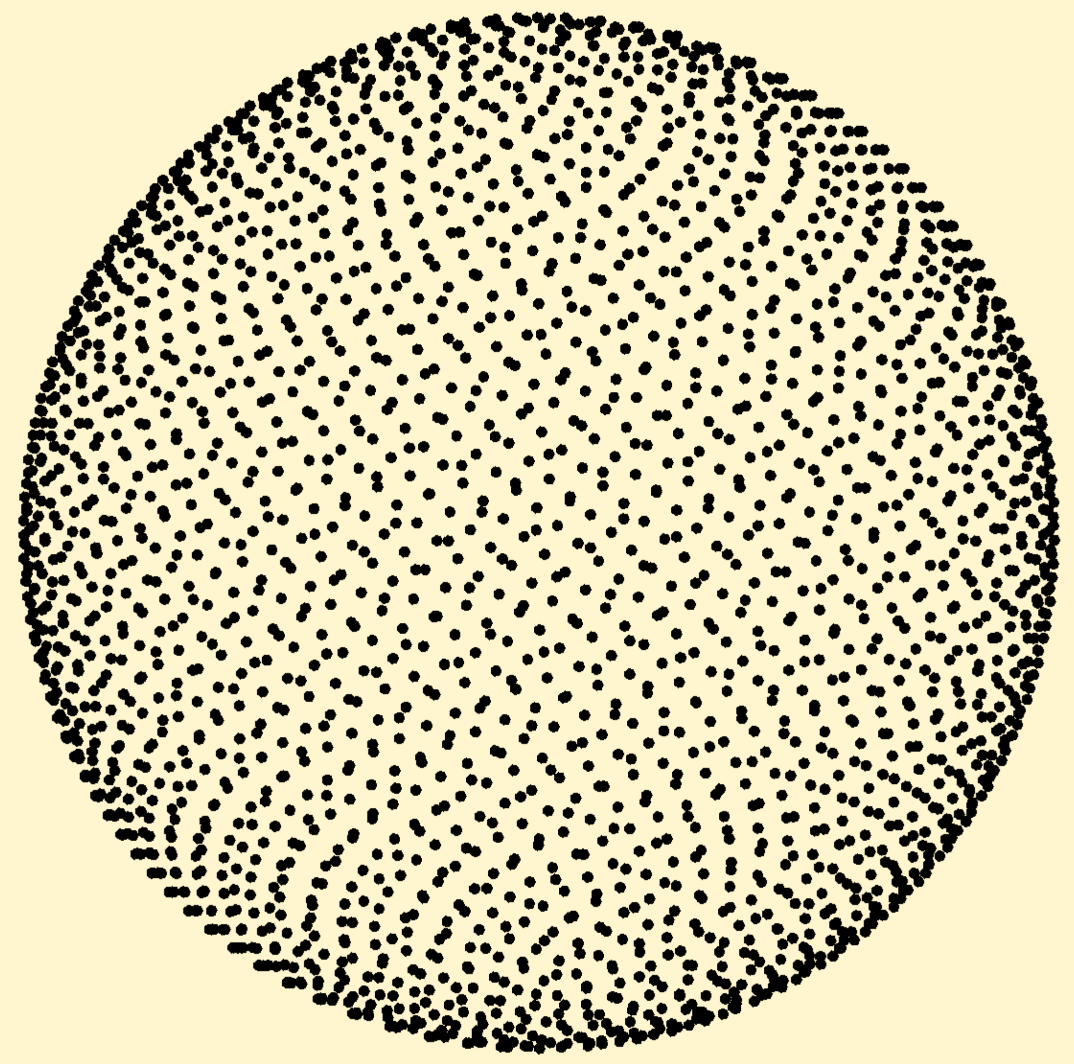

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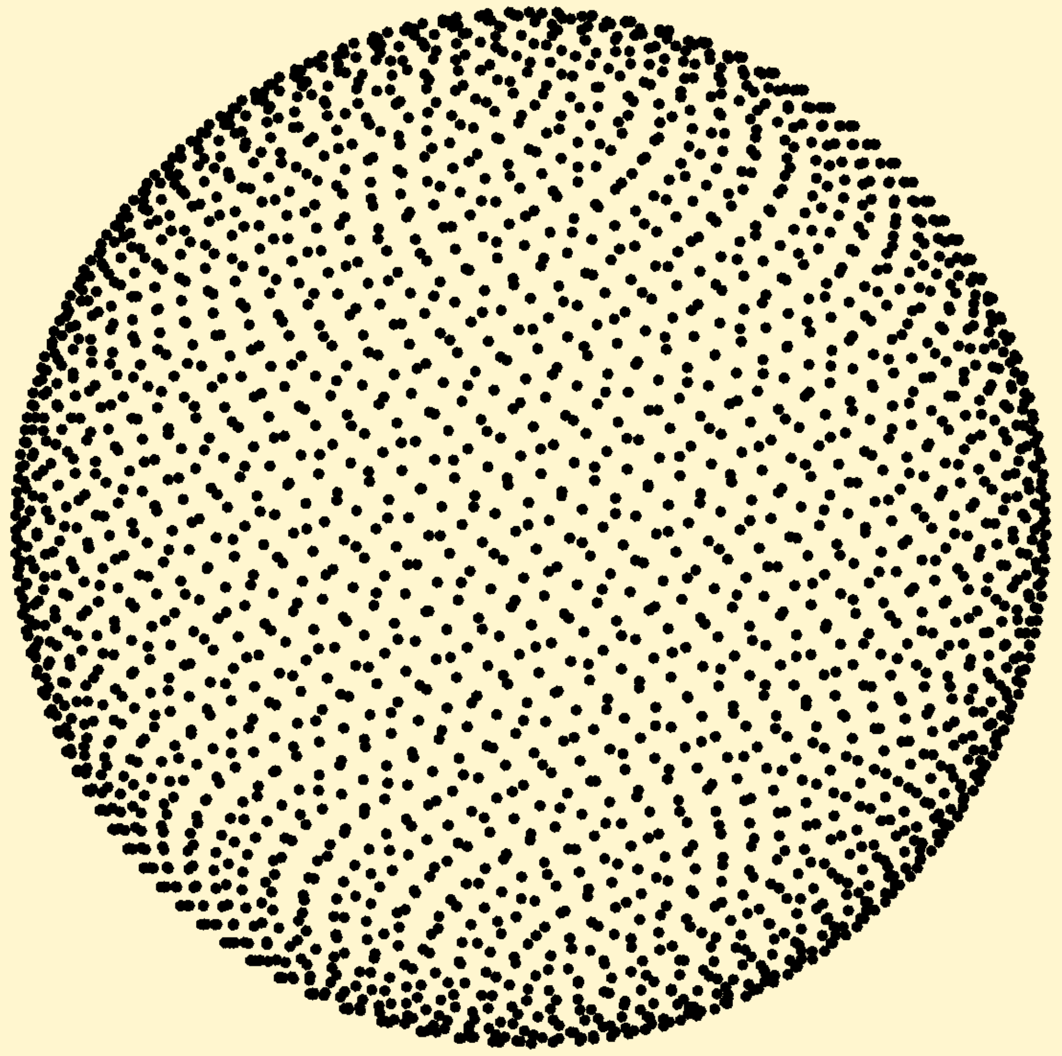
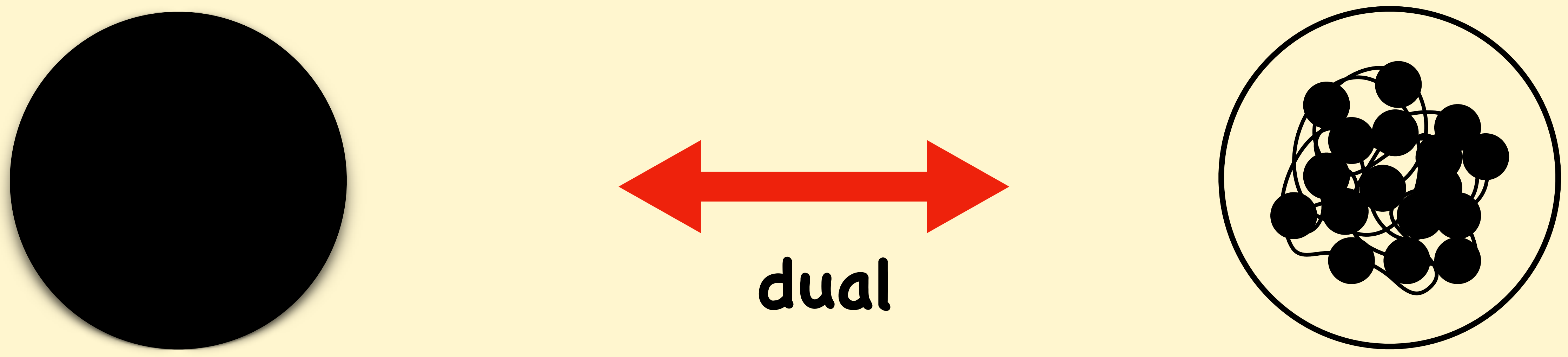


Gauge/gravity duality in string theory



Black p-brane
in IIA/IIB string

(p+1)-d U(N) SYM
(Dp-branes + strings)

In this talk p=0

The curious case of p=0

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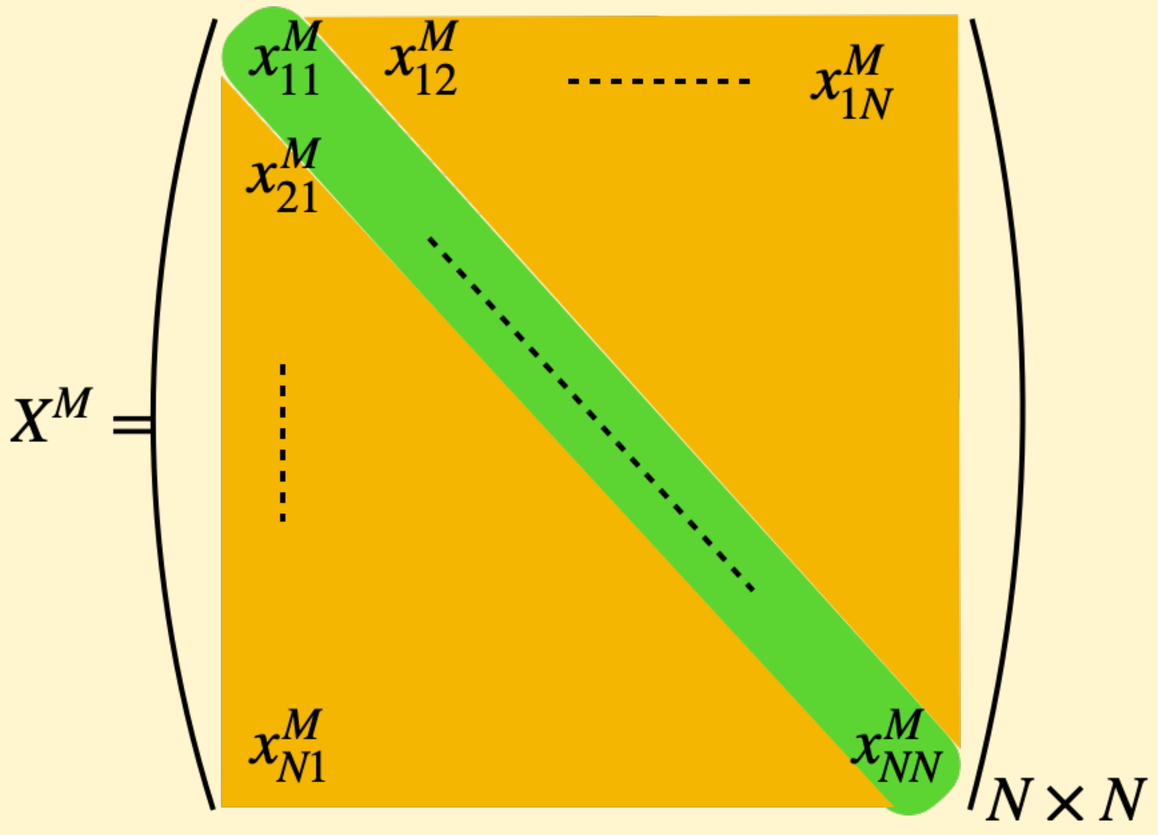
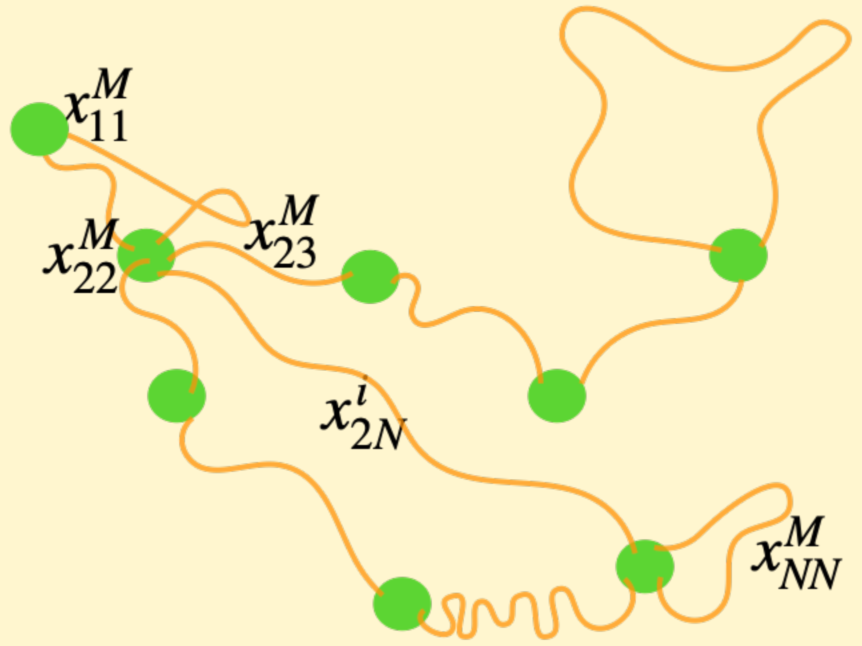


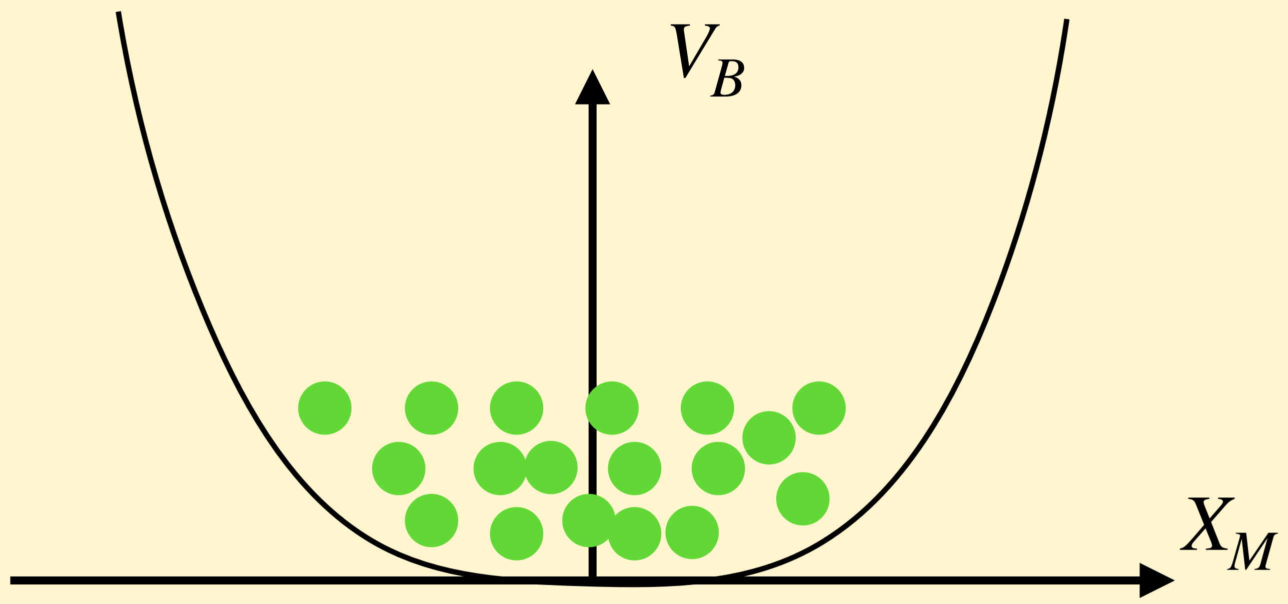


Witten 1995

$\mathbb{R}^9$

$M = 1, \cdots, 9$

$$V_B = [X_M, X_N]^2 \sim X_M^4$$



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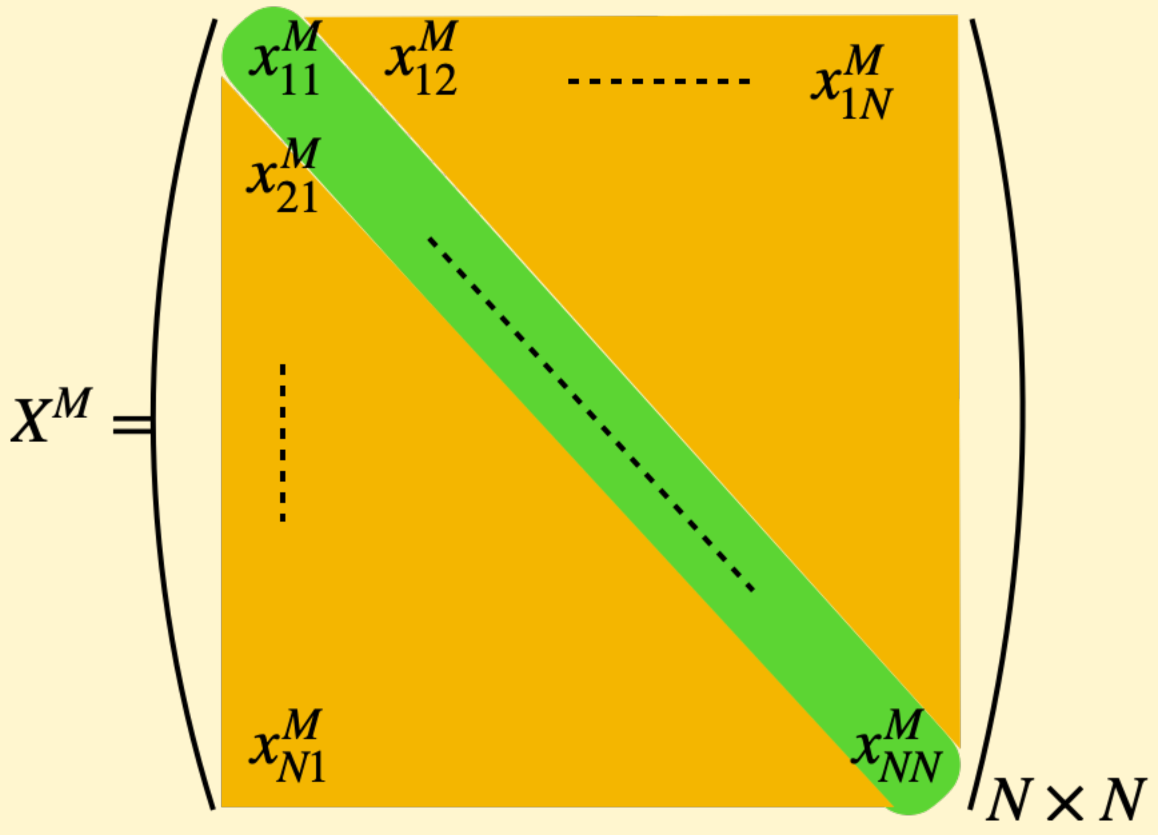

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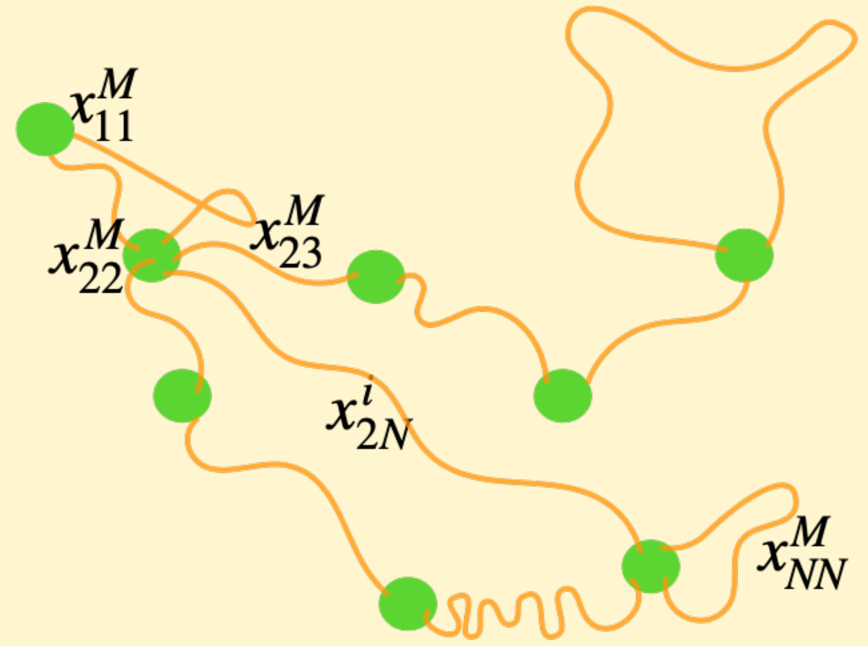

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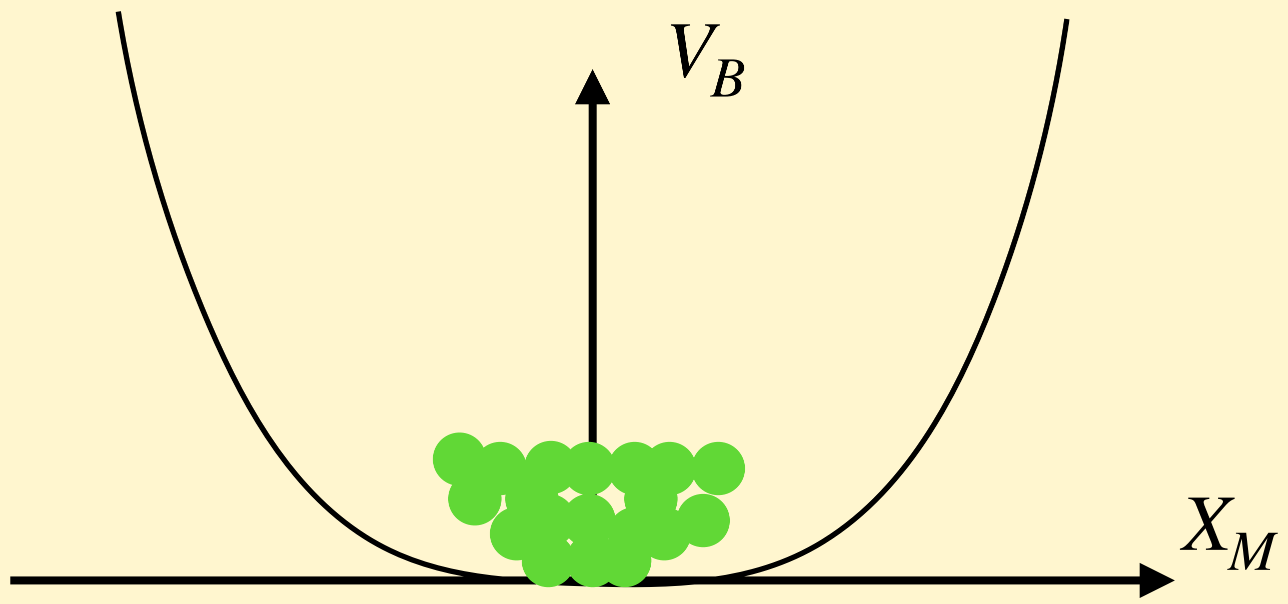

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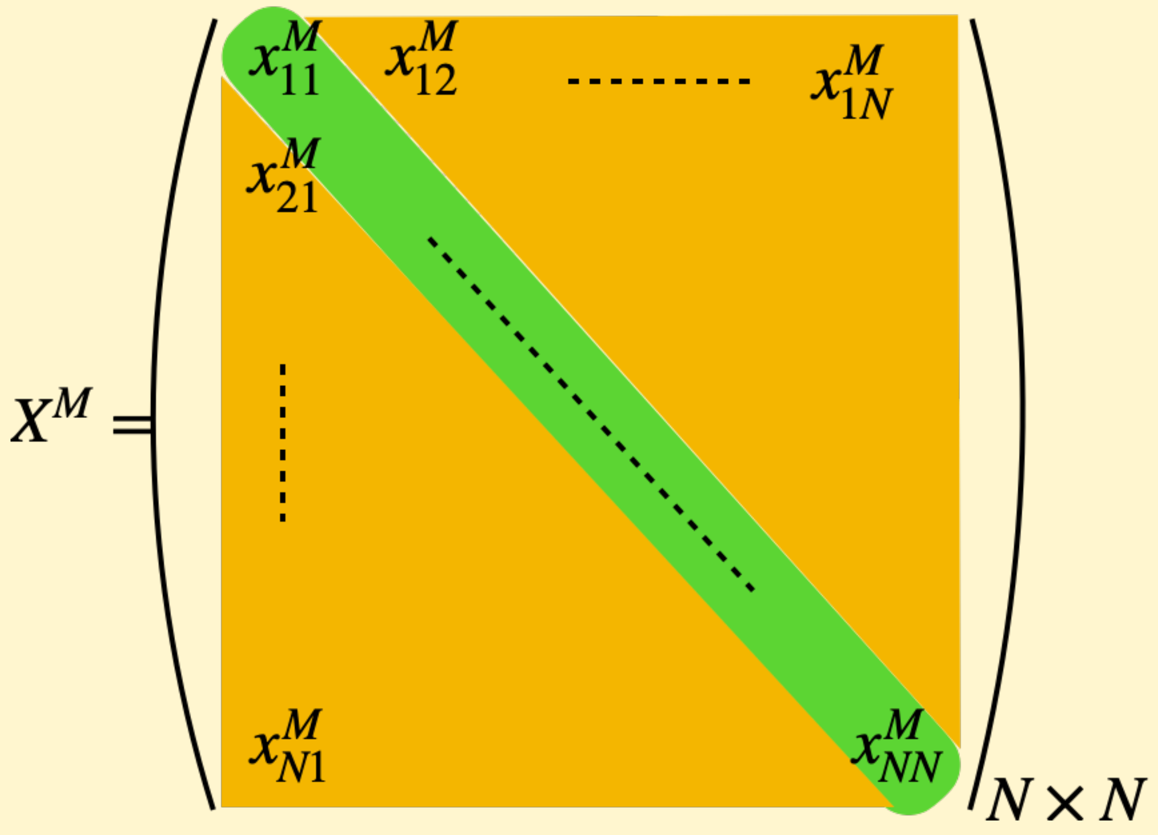

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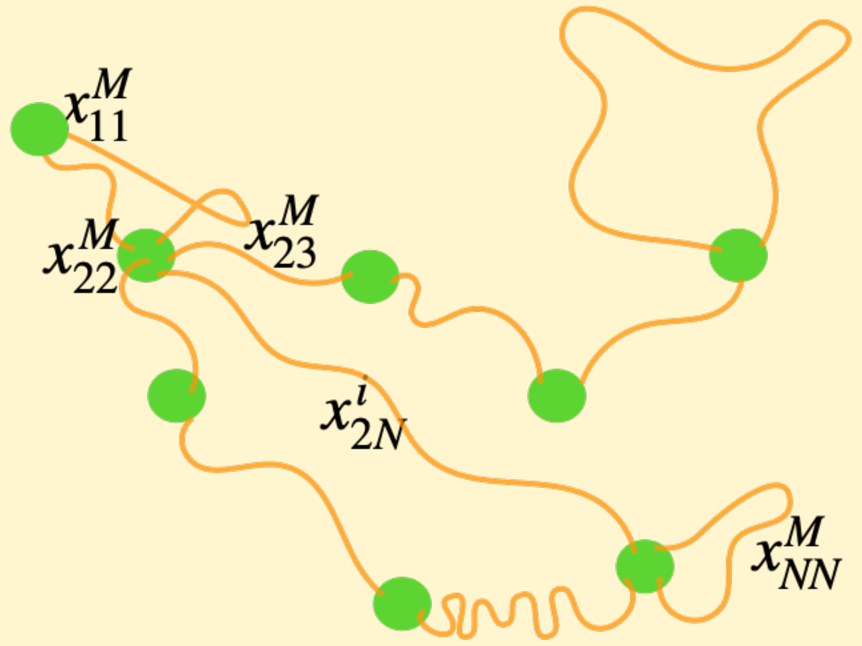

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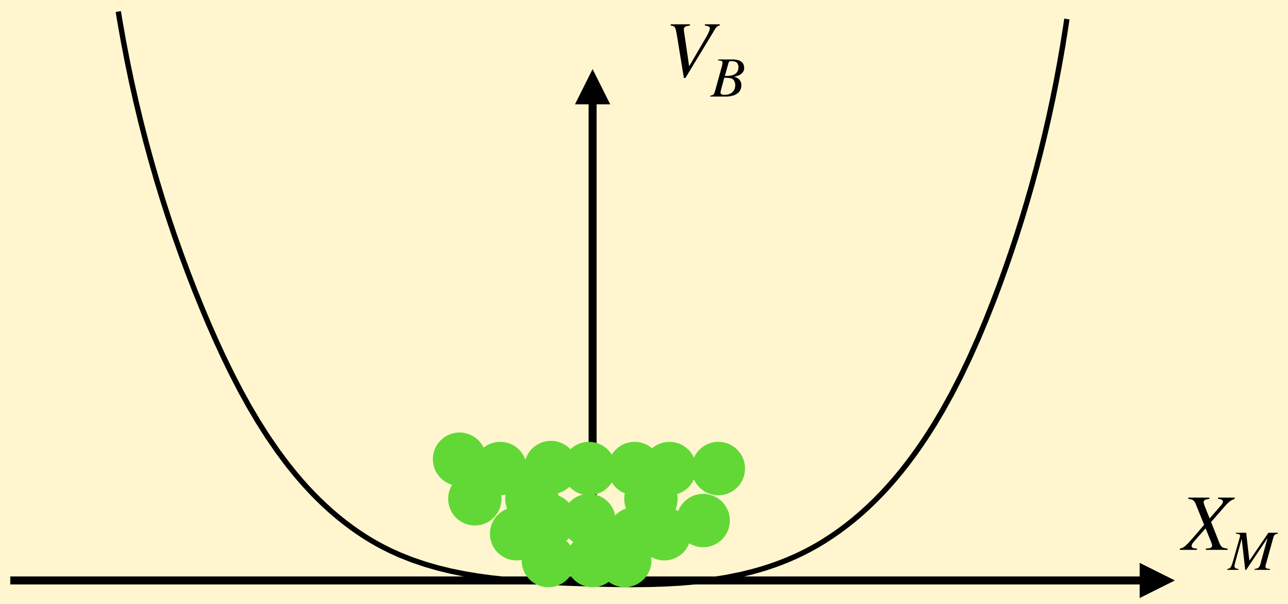


What does this correspond to in the gravity side?

Deformation of the D0-matrix model

The BMN model

Berestein, Maldacena, Nastase, 2002

$$S_{BMN} = S_b + S_f + \Delta S_b + \Delta S_f$$

BFSS

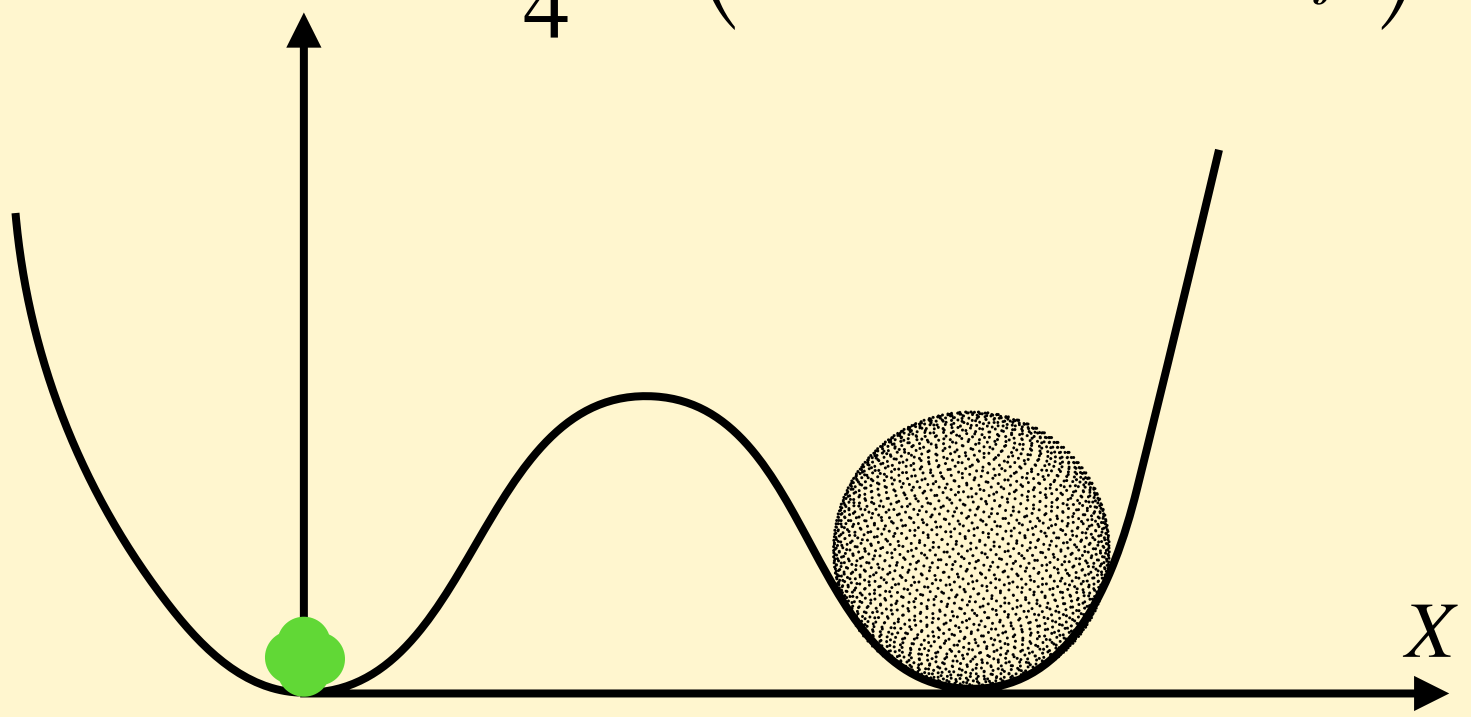
$$S_b = \frac{N}{\lambda}\int_0^\beta dt Tr\left\{\frac{1}{2}\sum_{I=1}^{9}(D_t X_I)^2 - \frac{1}{4}\sum_{I,J=1}^{9}[X_I, X_J]^2\right\},$$

$$S_f = \frac{N}{\lambda}\int_0^\beta dt Tr\left\{i\bar{\psi}\gamma^{10}D_t\psi - \sum_{I=1}^{9}\bar{\psi}\gamma^I[X_I,\psi]\right\},$$

$$\Delta S_b = \frac{N}{\lambda}\int_0^\beta dt Tr\left\{\frac{\mu^2}{2}\sum_{i=1}^{3}X_i^2 + \frac{\mu^2}{8}\sum_{a=4}^{9}X_a^2 + i\mu\sum_{i,j,k=1}^{3}\epsilon^{ijk}X_iX_jX_k\right\},$$

$$\Delta S_f = \frac{3i\mu N}{4\lambda}\int_0^\beta dt Tr\left(\bar{\psi}\gamma^{123}\psi\right),$$

$$V_{SO(3)} = \frac{1}{4}Tr\left(\mu\epsilon^{ijk}X_k + i[X_i, X_j]\right)^2$$



- Mass terms for bosons, fermions
- $SO(9) \rightarrow SO(3) \times SO(6)$
- $SU(2)$ vacua, i.e fuzzy spheres

The curious case of p=0

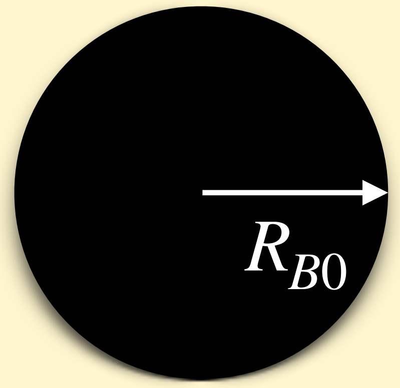
$$g_{\text{eff}} = \frac{\lambda}{E^3}$$

Strong coupling ⟷ Low energies

Black zero-brane in IIA SUGRA

$$\frac{ds^2}{\alpha'} = H(r)^{-\frac{1}{2}} f(r) dt^2 + H(r)^{\frac{1}{2}} \left(\frac{dr^2}{f(r)} + r^2 d\Omega_8^2 \right)$$

$$H(r) = \frac{240\pi^5 \lambda}{r^7},\ f(r) = 1 - \left(\frac{r_0}{r}\right)^7$$



$$E = 7.41 N^2 \lambda^{-3/5} T^{14/5},\ \ S = 11.52 N^2 \lambda^{-3/5} T^{9/5}$$

$$\frac{R_{B0}^2}{\alpha'} \sim g_{\text{eff}}^{\frac{1}{2}} \sim \sqrt{\frac{\lambda}{E^3}}$$

$$e^{\phi}\Big|_{horizon} \sim \frac{g_{\text{eff}}^{7/4}}{N}$$

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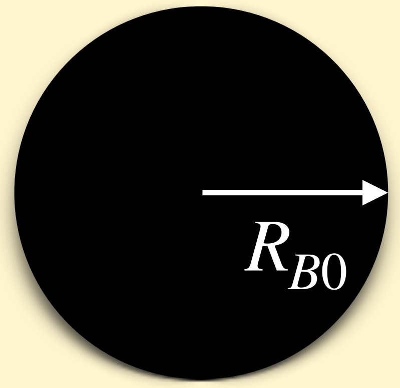

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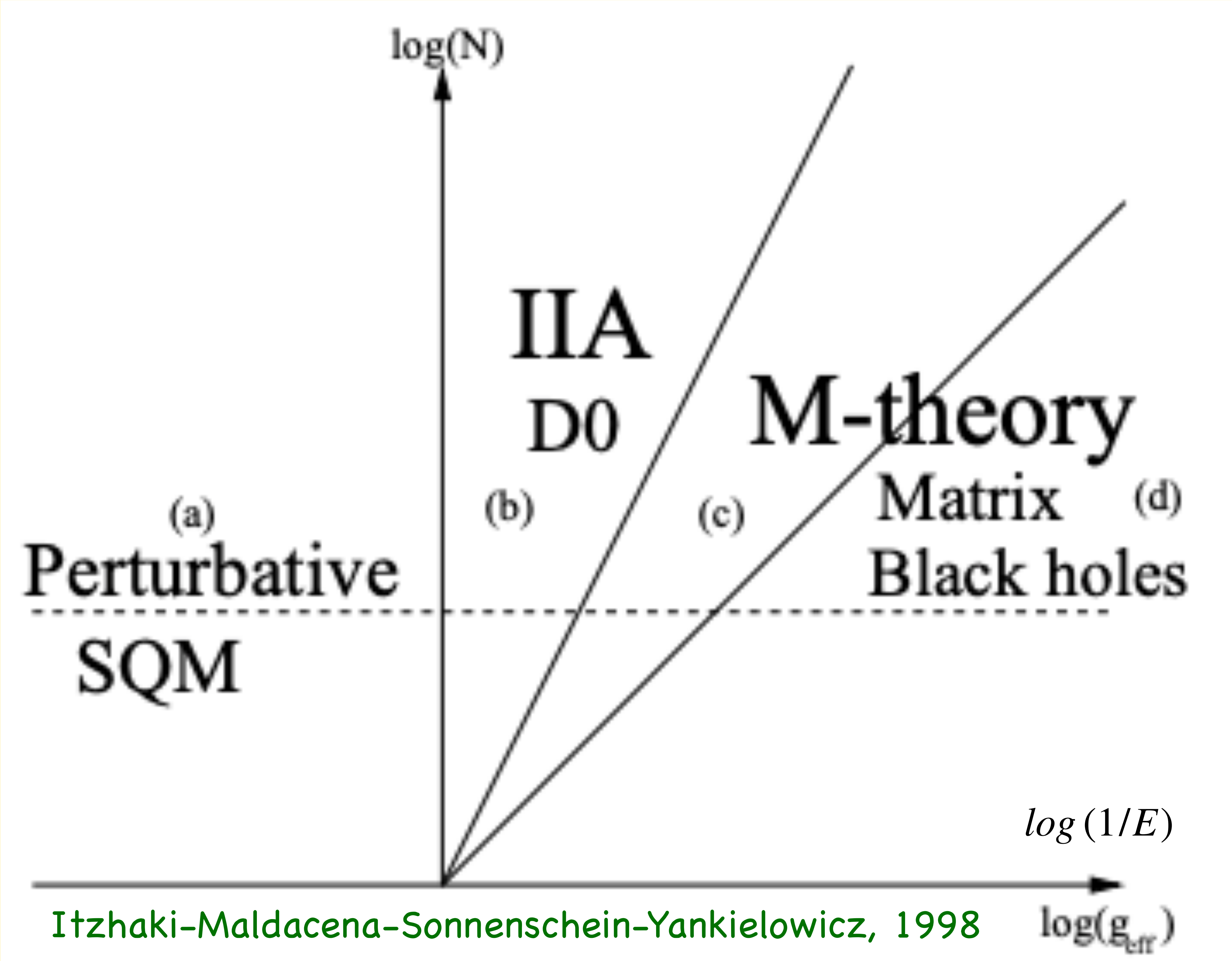

Itzhaki-Maldacena-Sonnenschein-Yankielowicz, 1998

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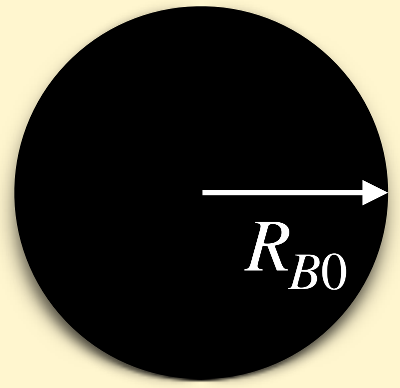

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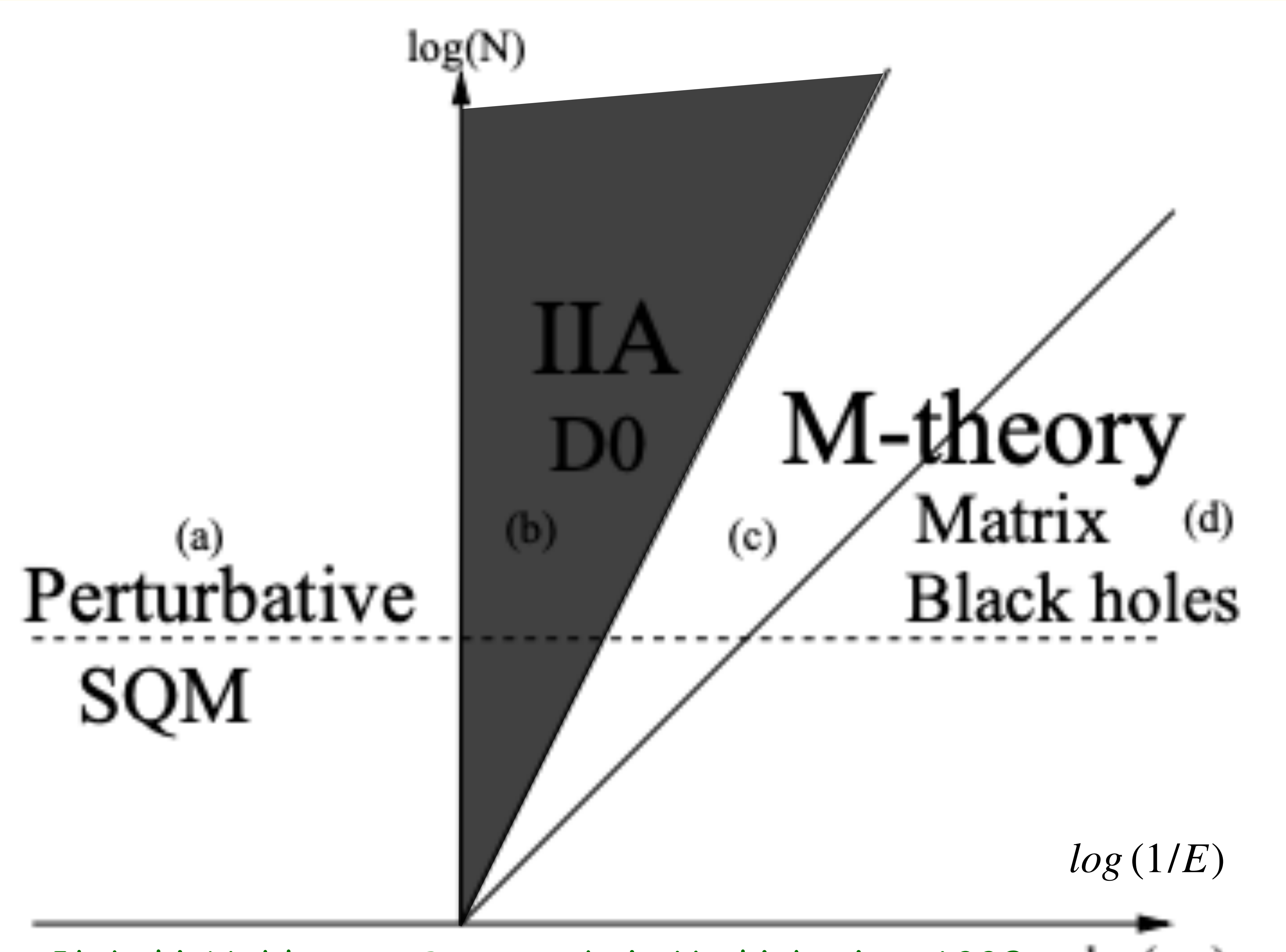

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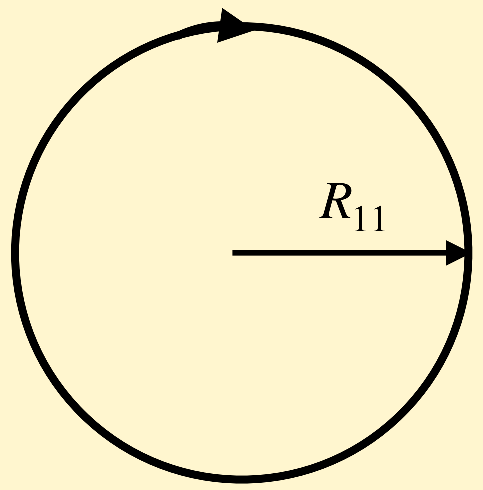

Type IIA string theory is defined as M-theory compactified on a circle $\mathbf{S}^1$

$R_{11} \sim 2\pi g_s l_s = 2\pi l_s e^{\phi} \simeq l_s \frac{g_{\rm eff}^{7/4}}{N}, \quad g_{\rm eff} = \frac{\lambda}{E^3}$ Strong coupling/low energies corresponds to the M-theory region

To probe M-theory region $E \ll 1$ $(E = 7.41 N^2 \lambda^{-3/5} T^{14/5})$ ⟶ Low temperatures

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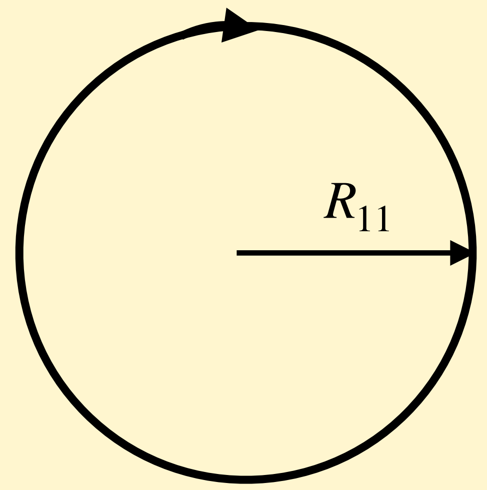

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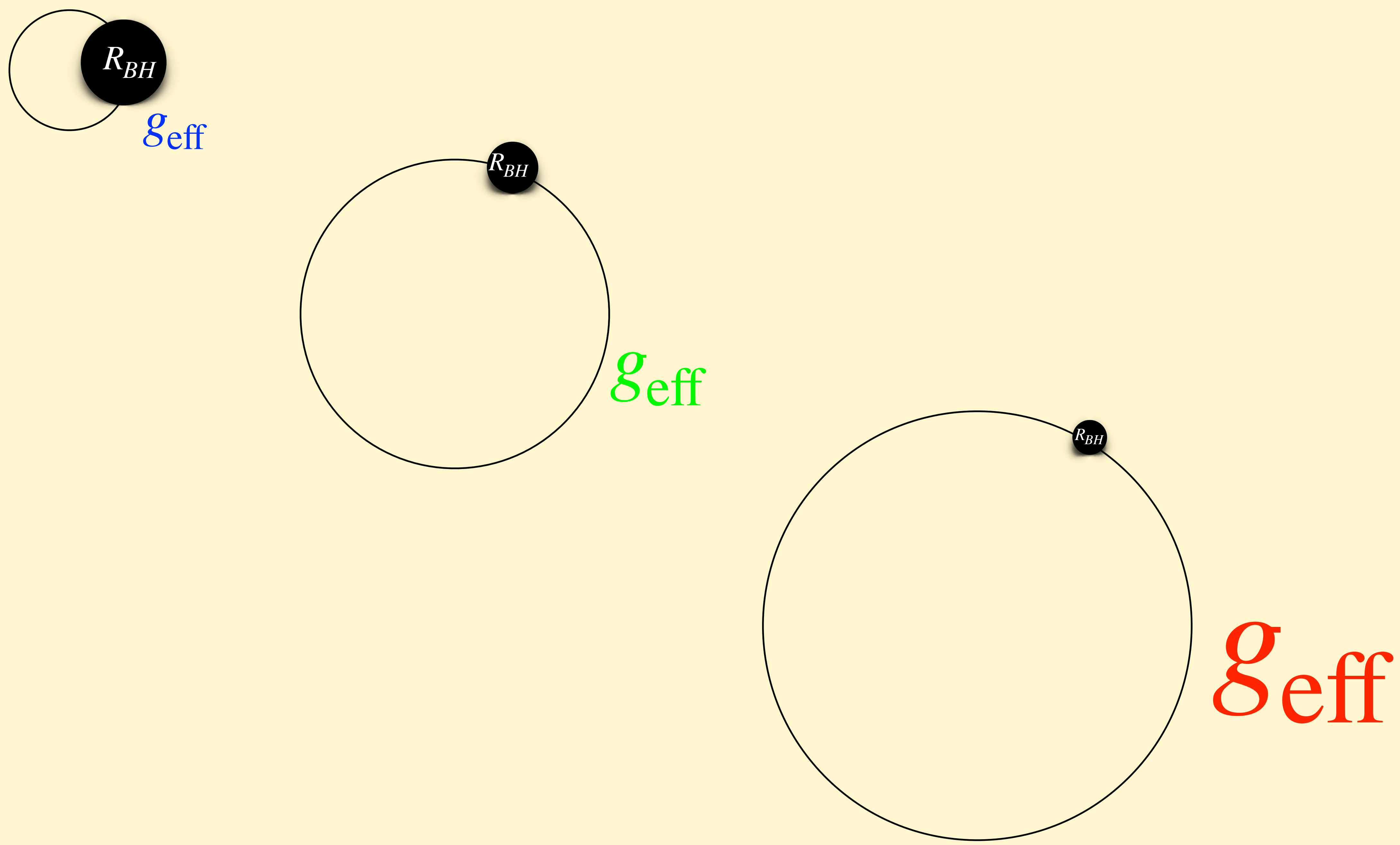
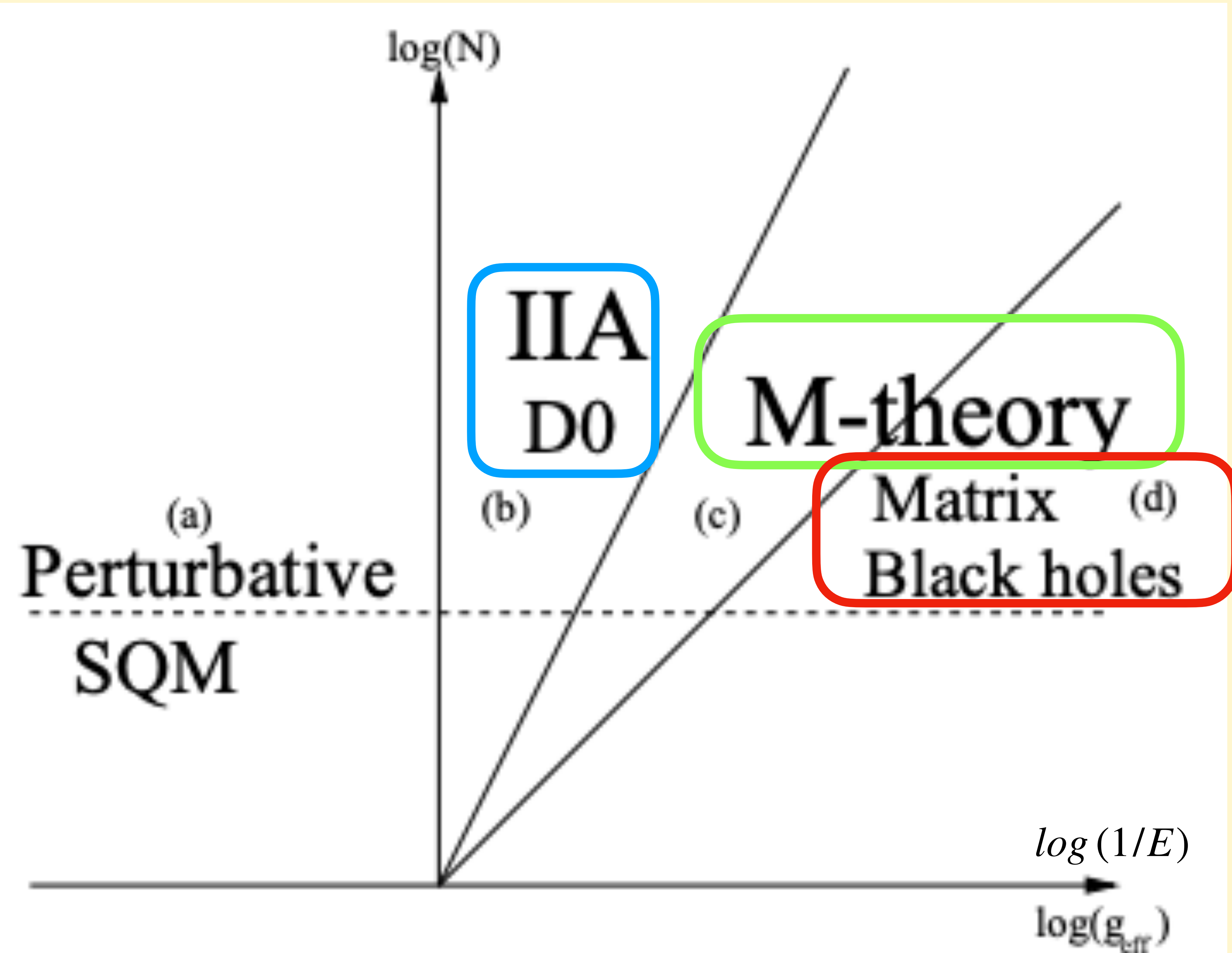
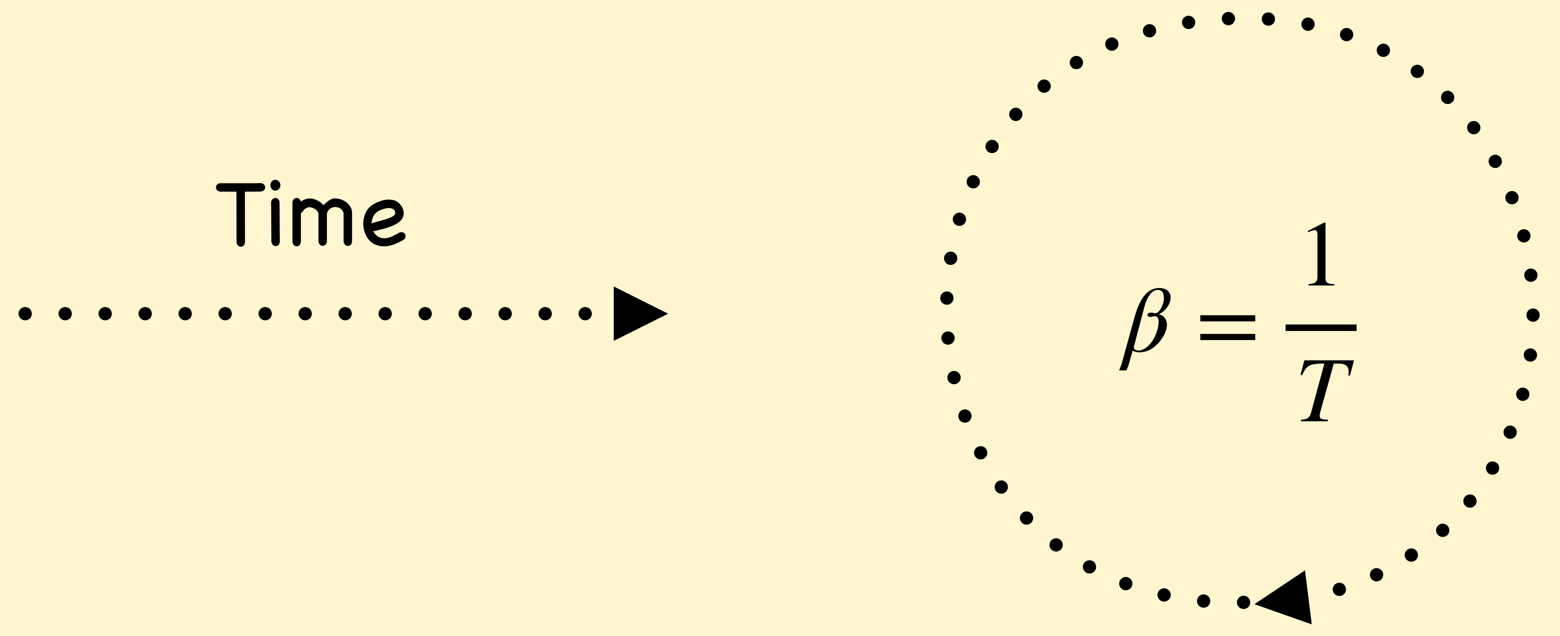




Switch to simulations

Model on the lattice

- We can do Monte Carlo simulations
- Borrow techniques from lattice QCD
- (0+1)-d matrix quantum mechanics



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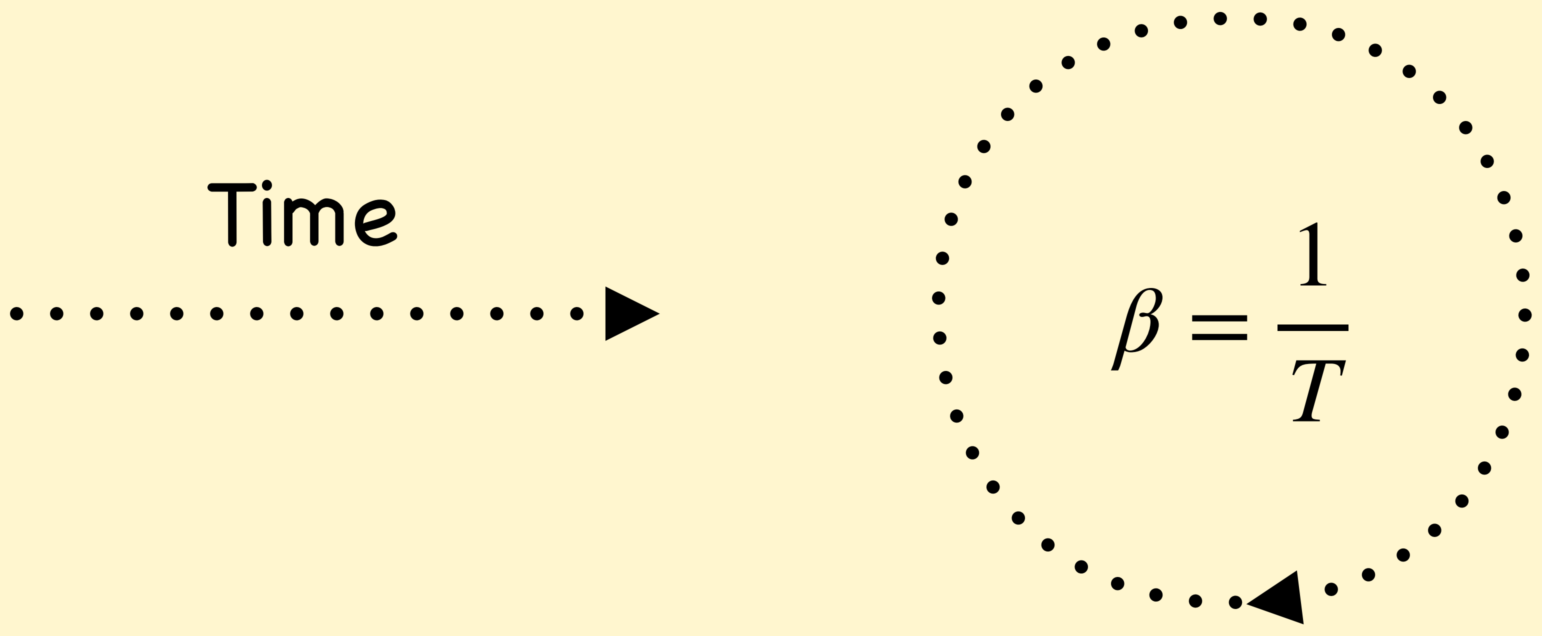

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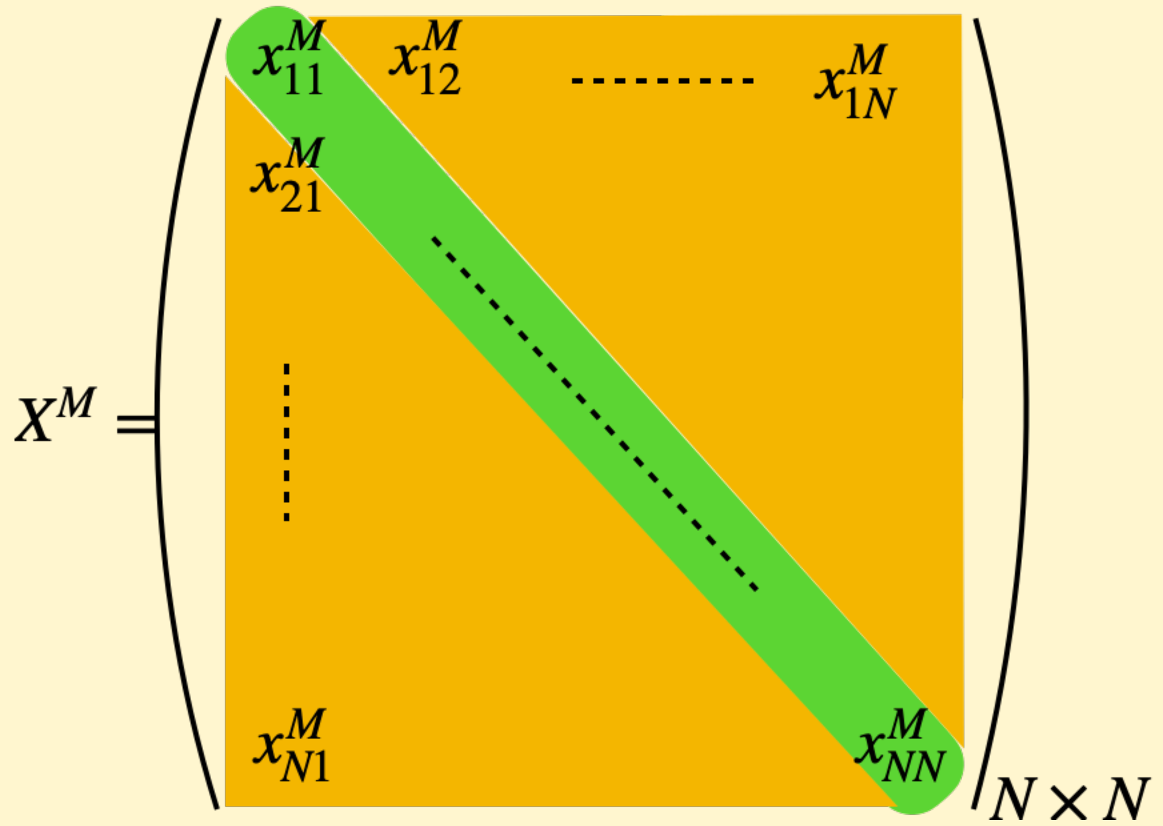
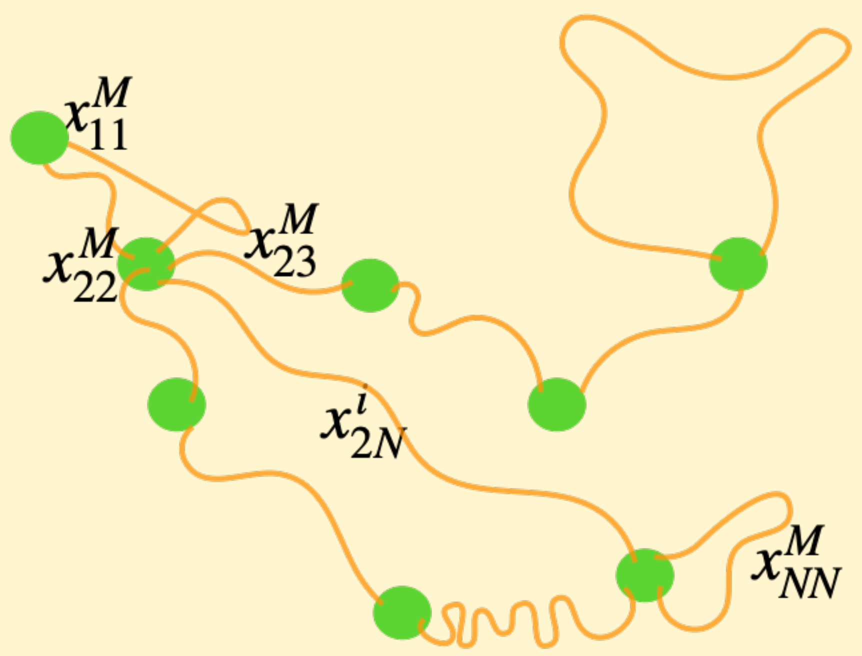
large N limit

$\lambda \sim N^0$, $T \sim N^0$ $\longleftrightarrow$ $\lambda = 1$: **fix**, $\lambda^{-\frac{1}{3}}T \sim N^0$: **fix**



Witten 1995



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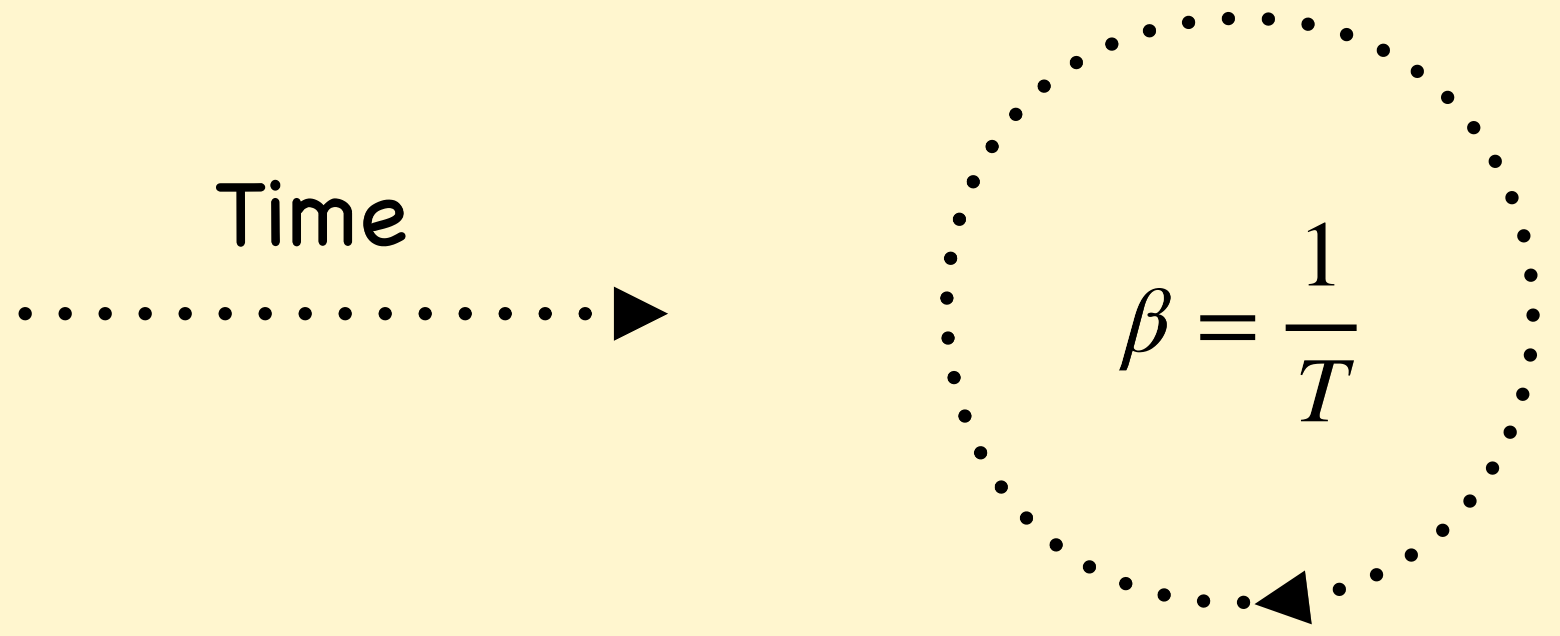

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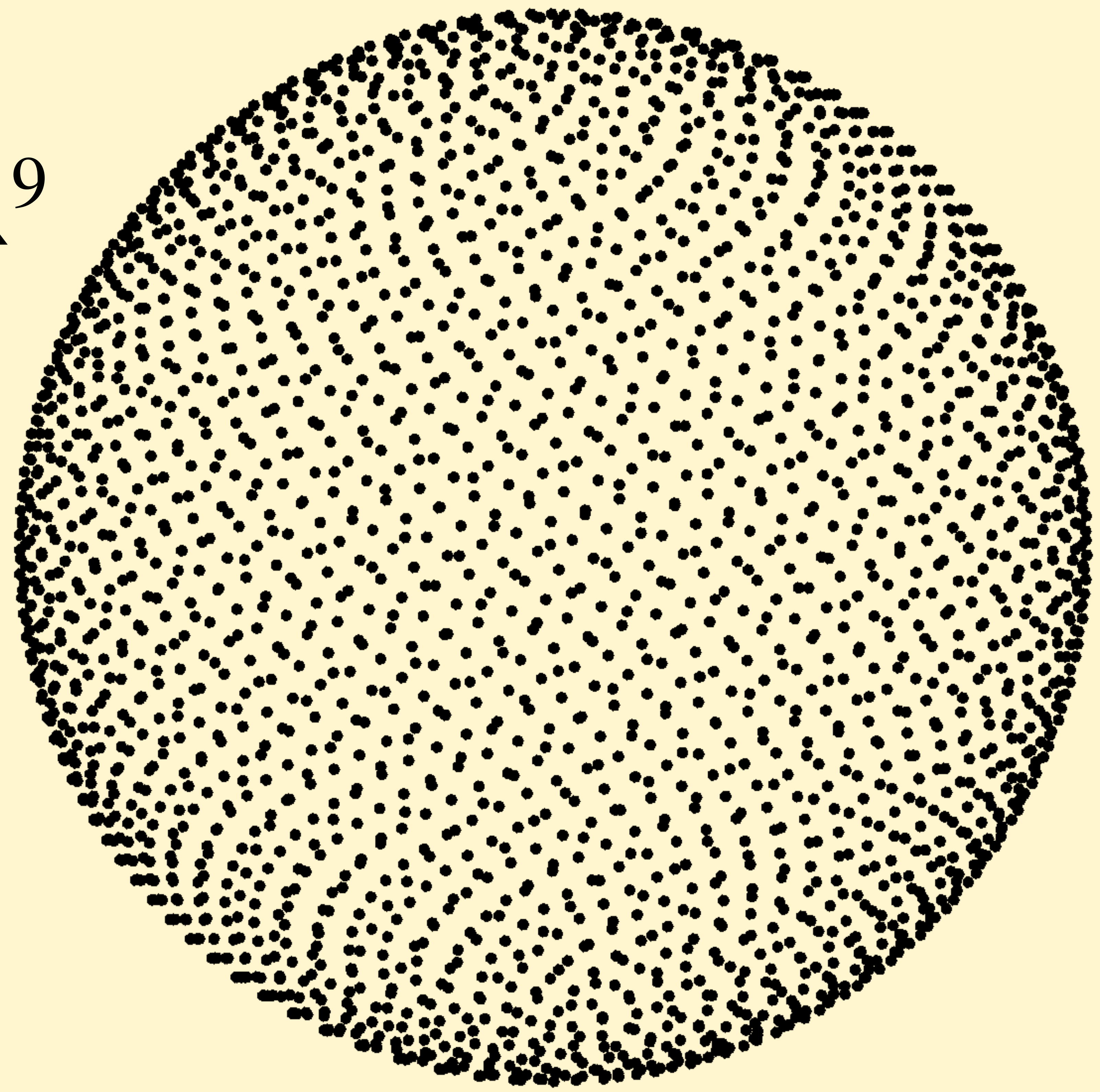
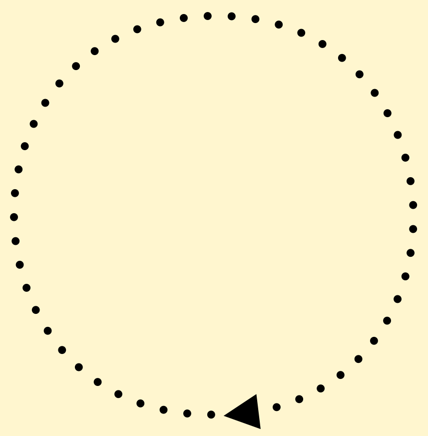


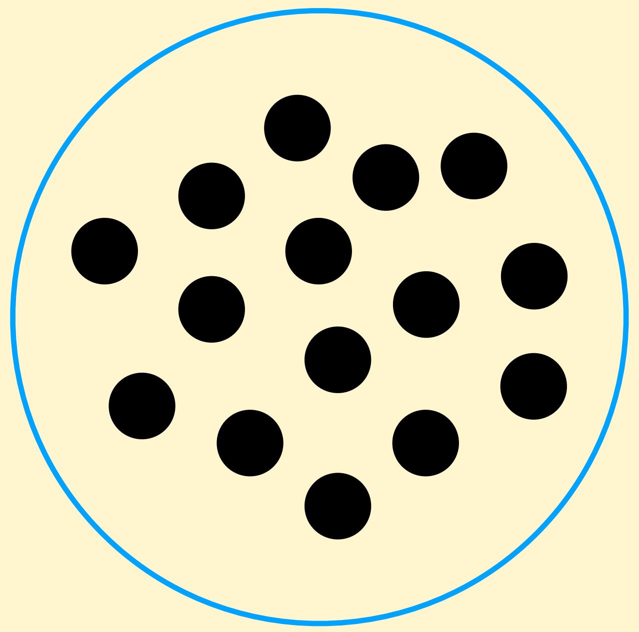
Confinement/deconfinement

- Polyakov loop $P = \frac{1}{N} Tr\left(\mathscr{P}exp\left(i\int_0^\beta dtA_t \right)\right) = \frac{1}{N}\sum_{j=1}^{N} e^{i\theta_j}$ On the lattice

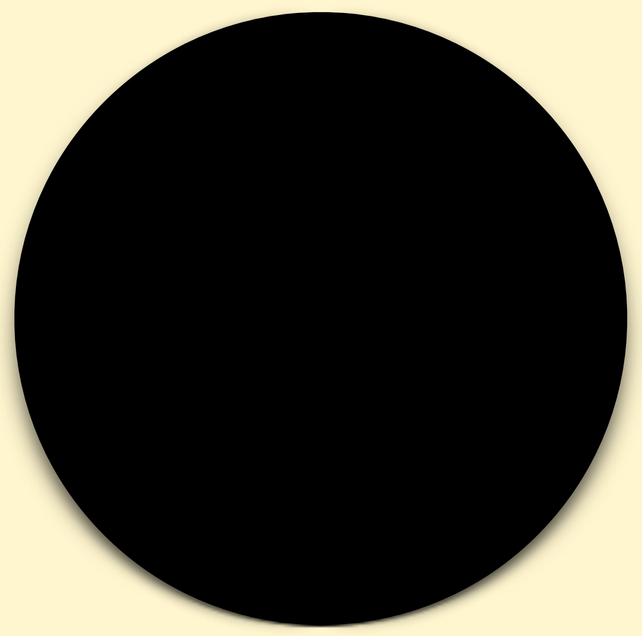


- Restoration/breaking of $U(1)$ symmetry

- Intuition from $AdS_5 \times S^5$



Gravity side



$E, P = 0$

Confinement

Graviton gas

MAGOO, Witten, Sundborg

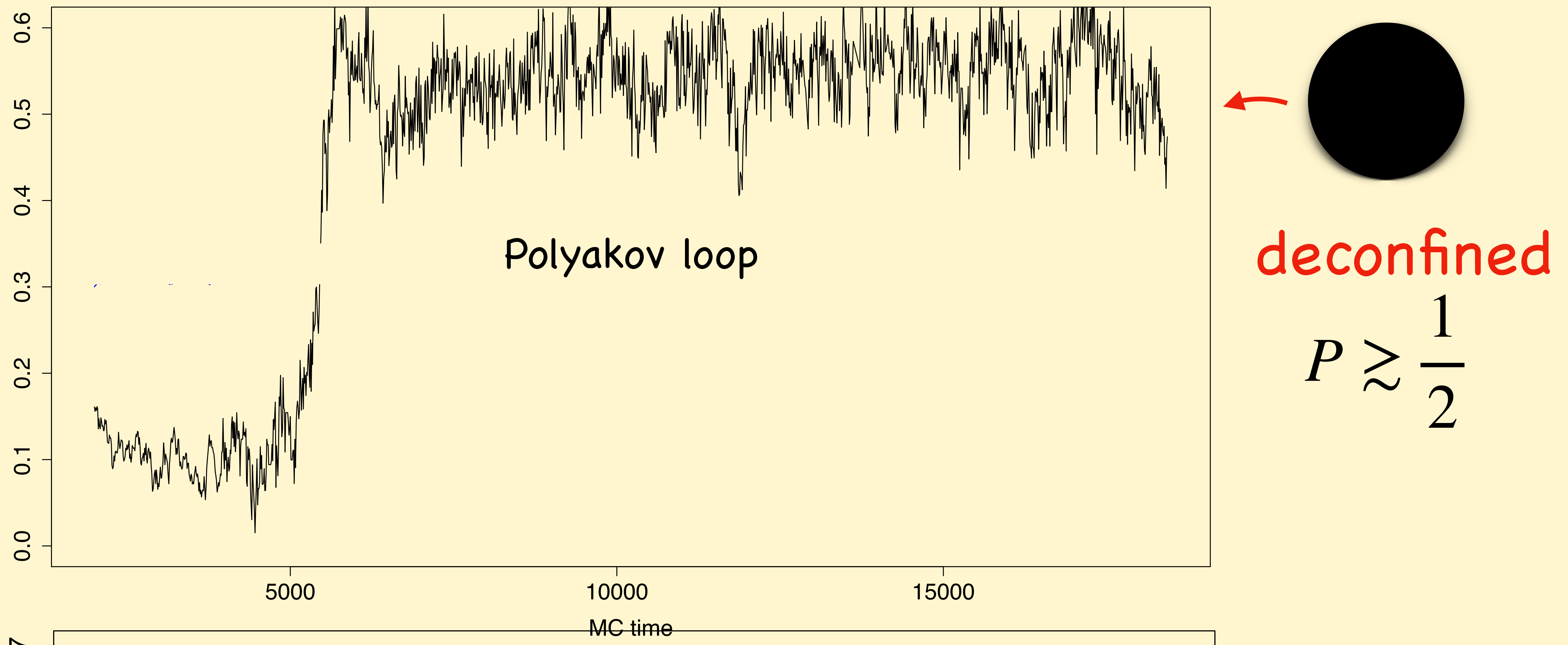
Aharony-Marsano-Minwalla-Papadodimas-Van Raamsdonk, 2003

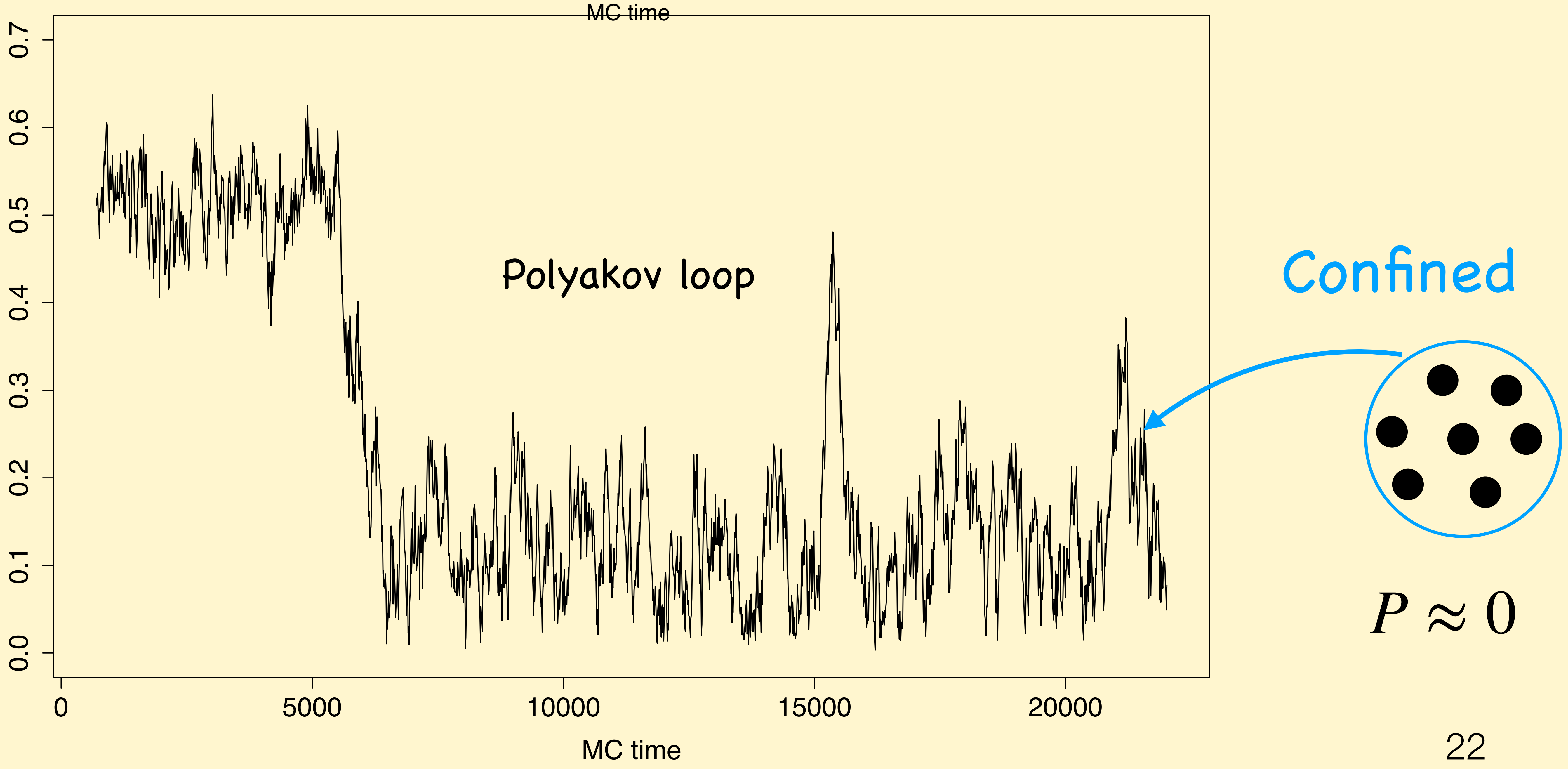
$E \neq 0,\ P \gtrsim \frac{1}{2}$

Deconfinement

Black holes

Confinement/deconfinement





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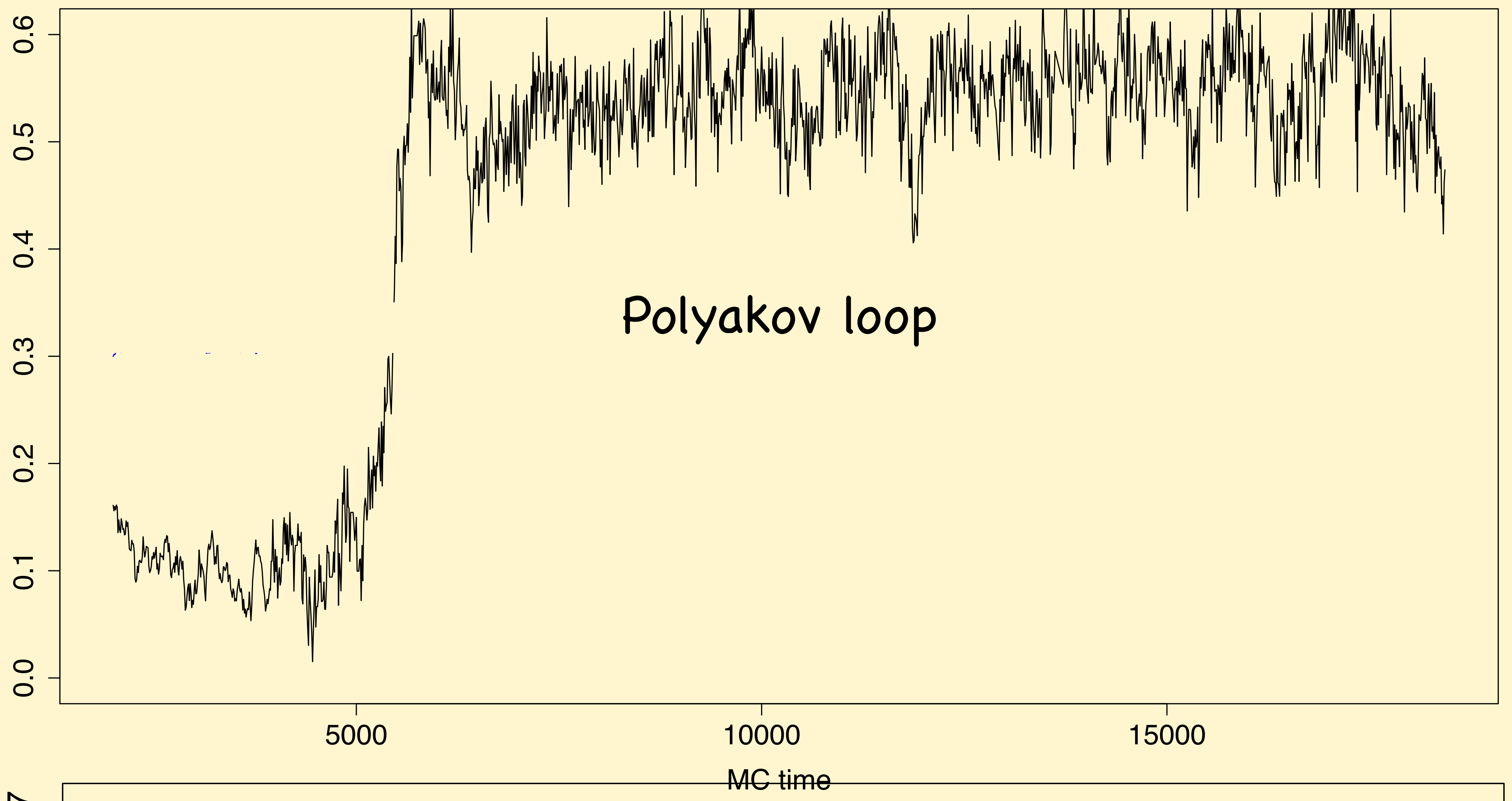
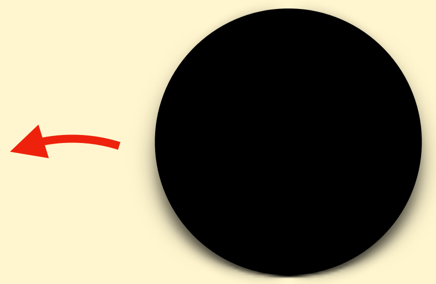
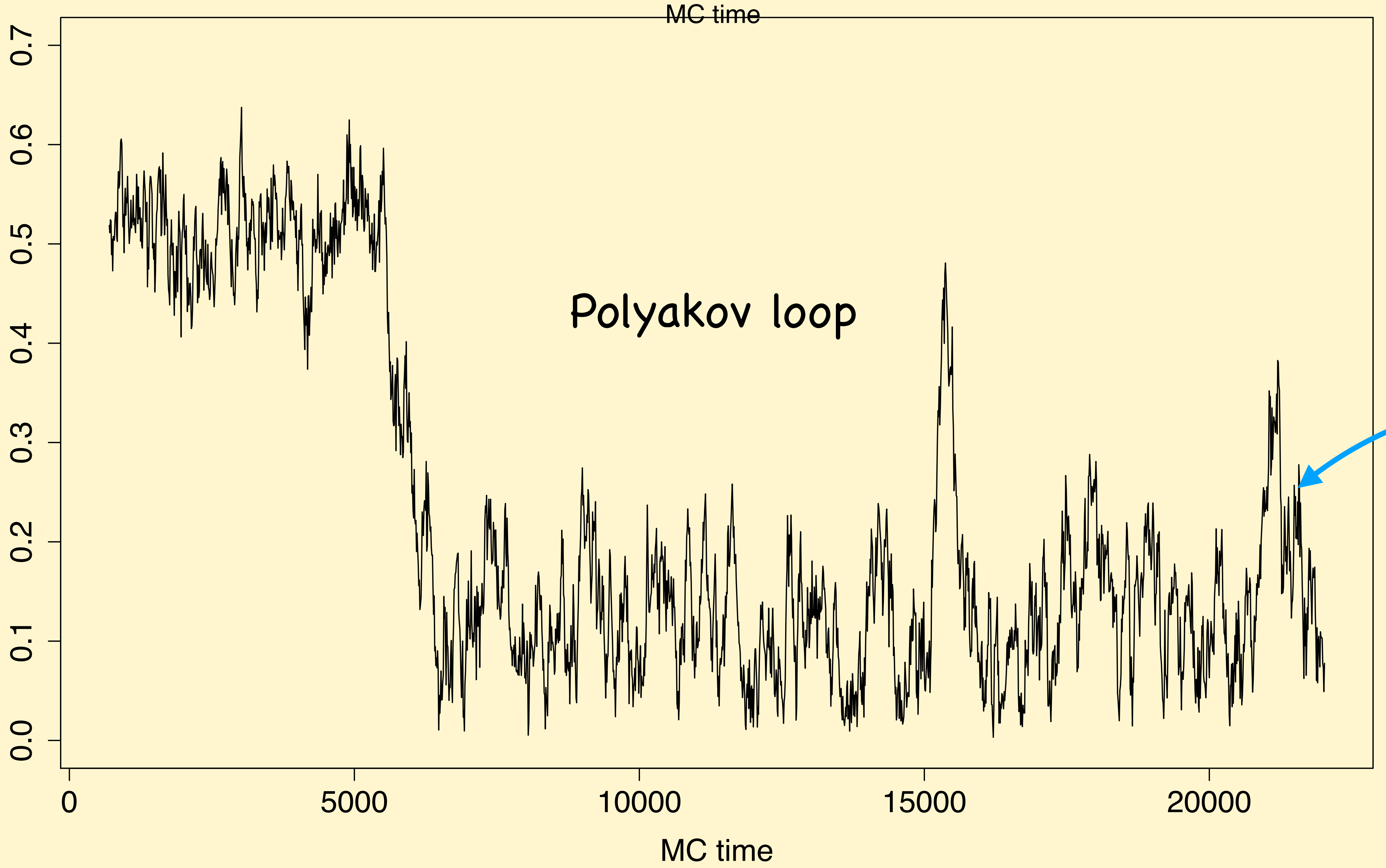


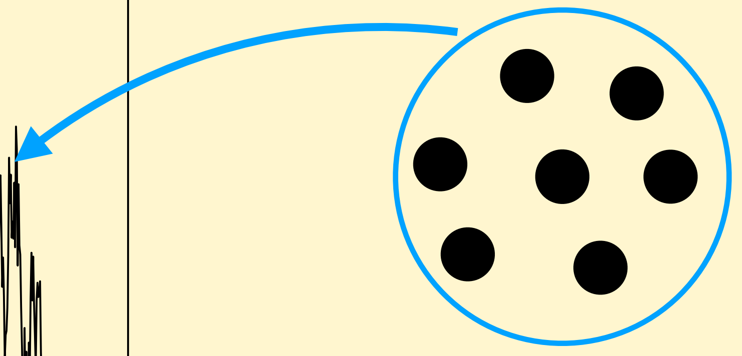


deconfined

$$P \gtrsim \frac{1}{2}$$



Confined



$$P \approx 0$$

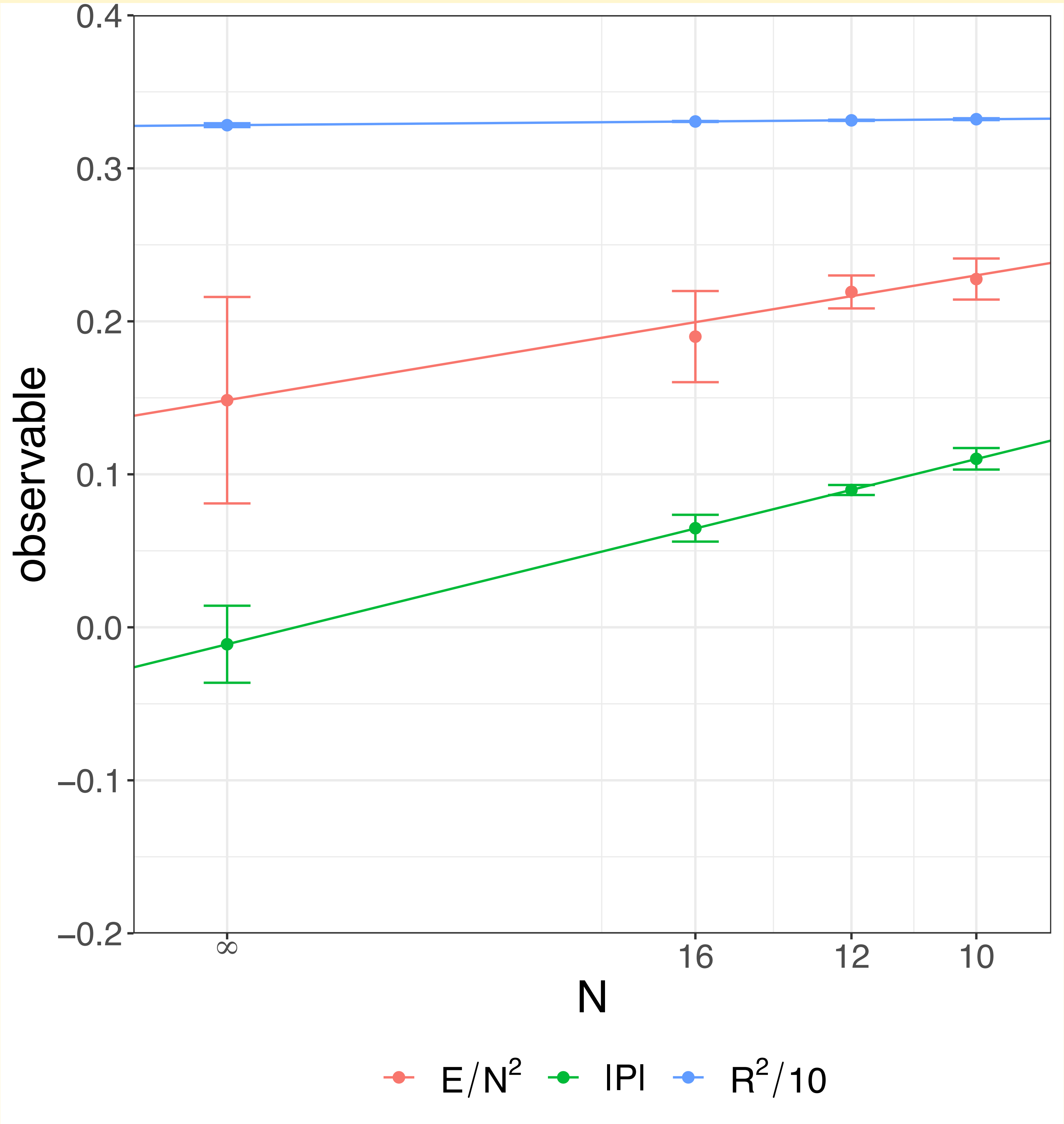
The gravity theory predicts always deconfinement

$$(E = 7.41 N^2 \lambda^{-3/5} T^{14/5})$$

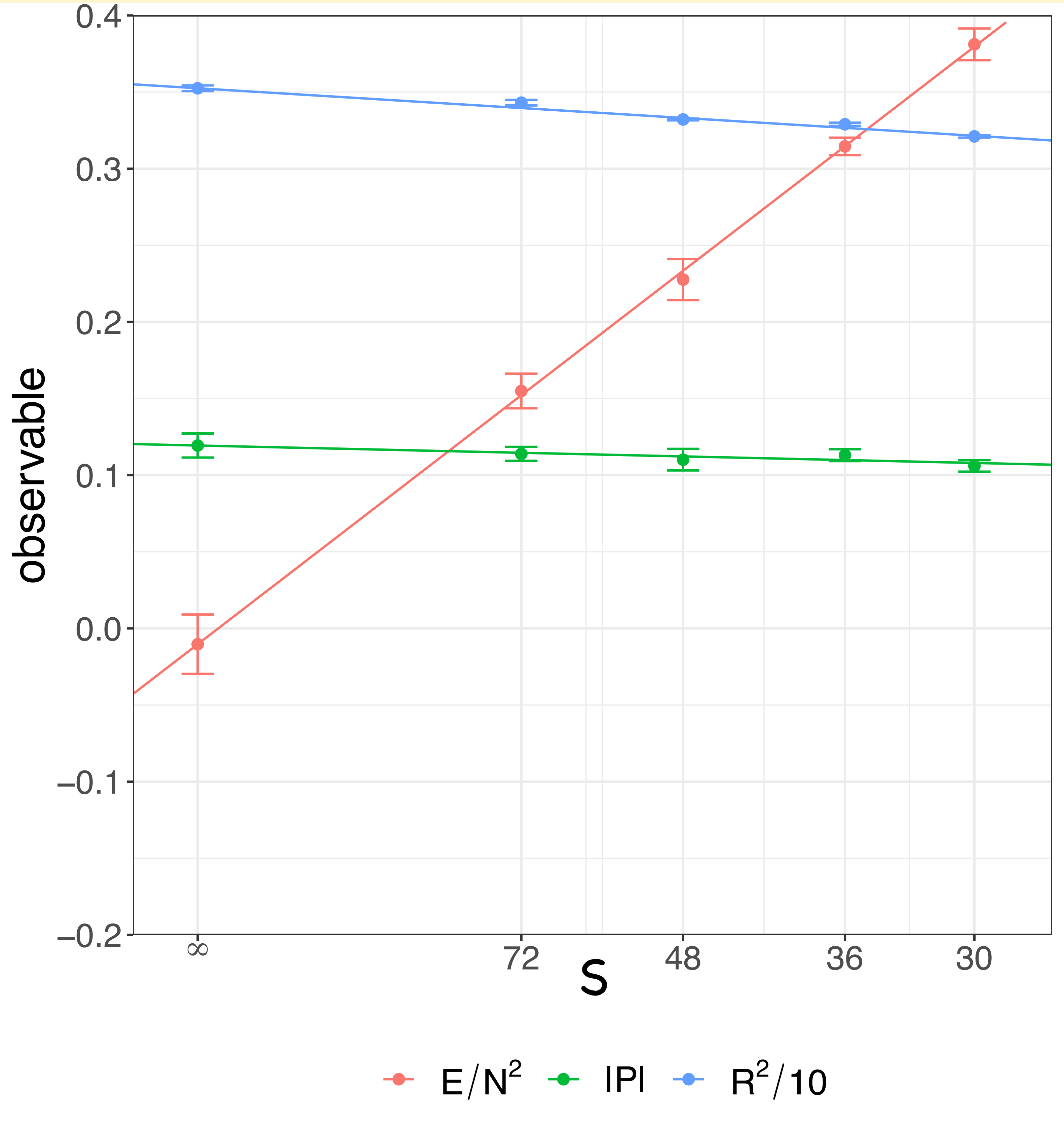
How to understand this?

Confined studies of D0-matrix model

Large N and S=48 @ T=0.2



Continuum and N=10 @ T=0.2



Combine both

Deconfined phase

$$\frac{E}{N^2} \simeq 7.41 T^{\frac{14}{5}} \simeq 0.0818 \quad \text{@ T=0.2}$$

$$P \simeq 0.5 \quad \text{@ T=0.2}$$

Confined phase

$$E \simeq 0$$

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@ large N and continuum

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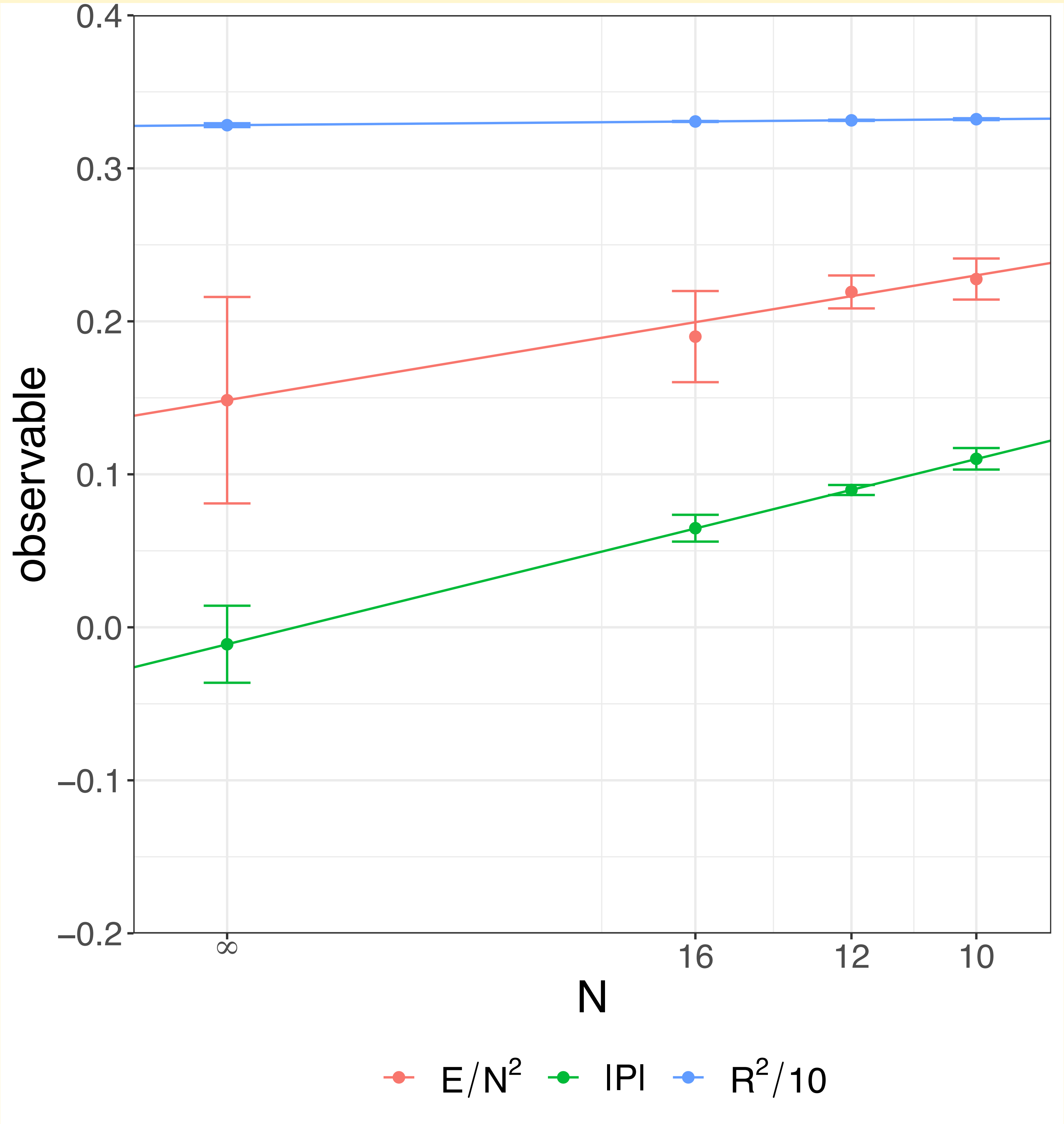

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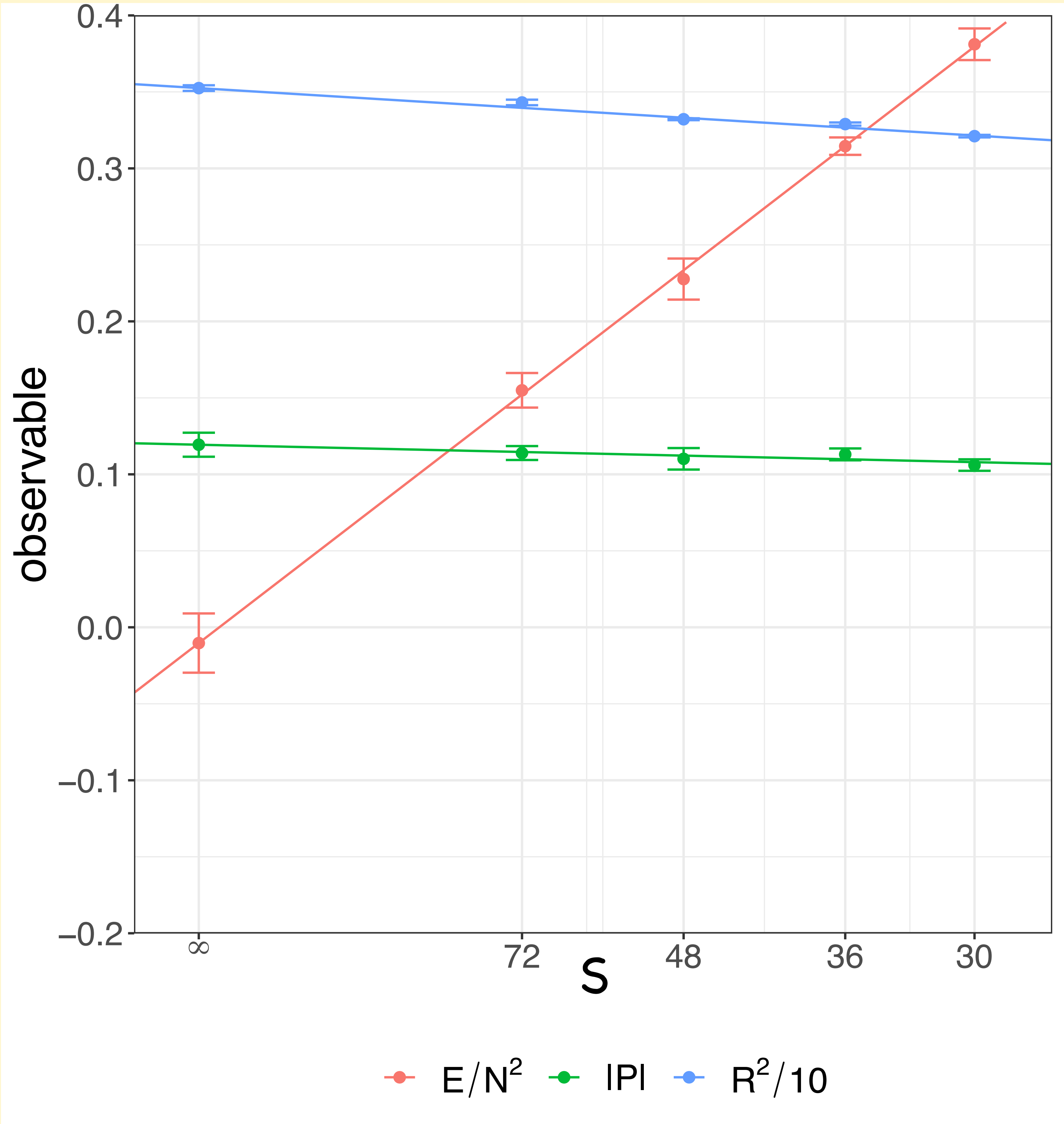

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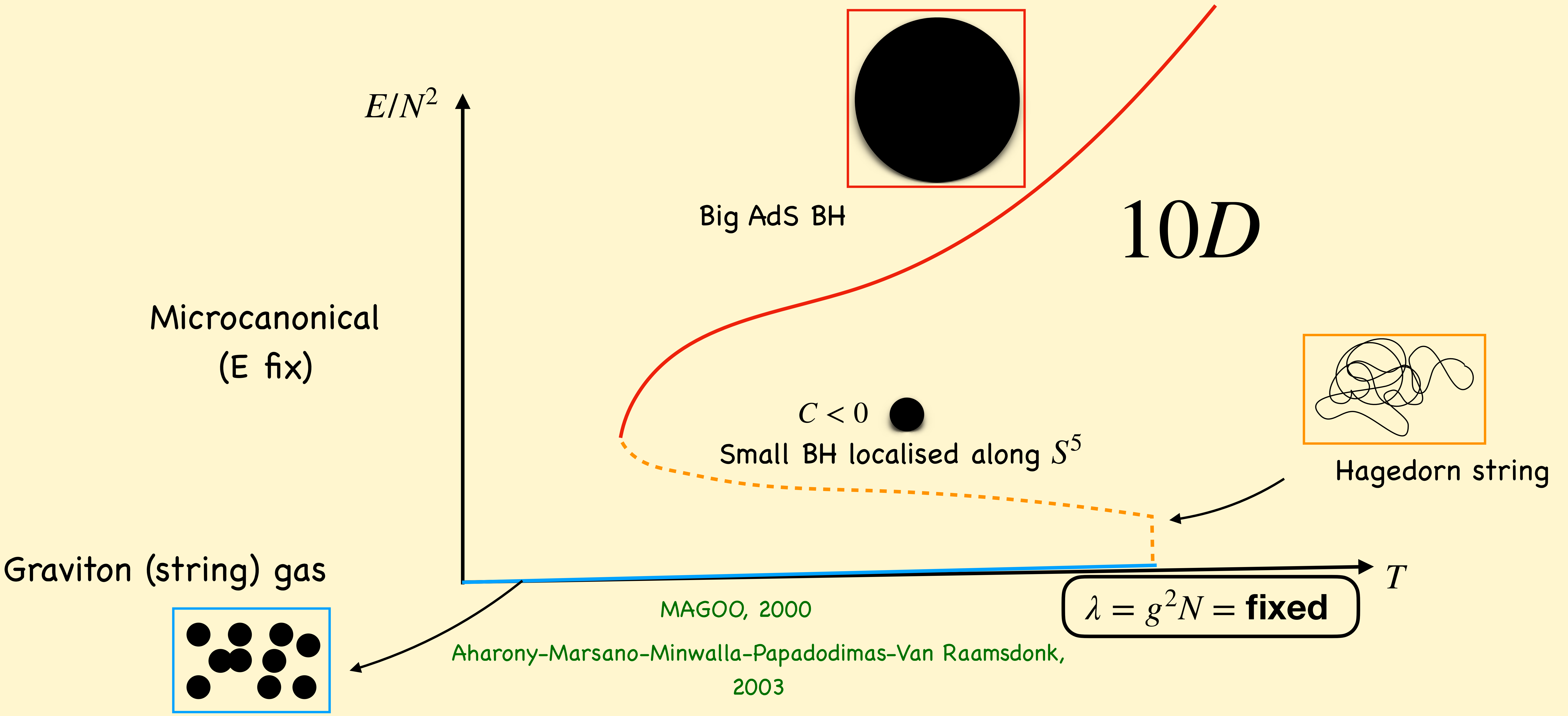
Conventional holography

How to understand confinement?

Motivation from

$AdS_5 \times S^5$



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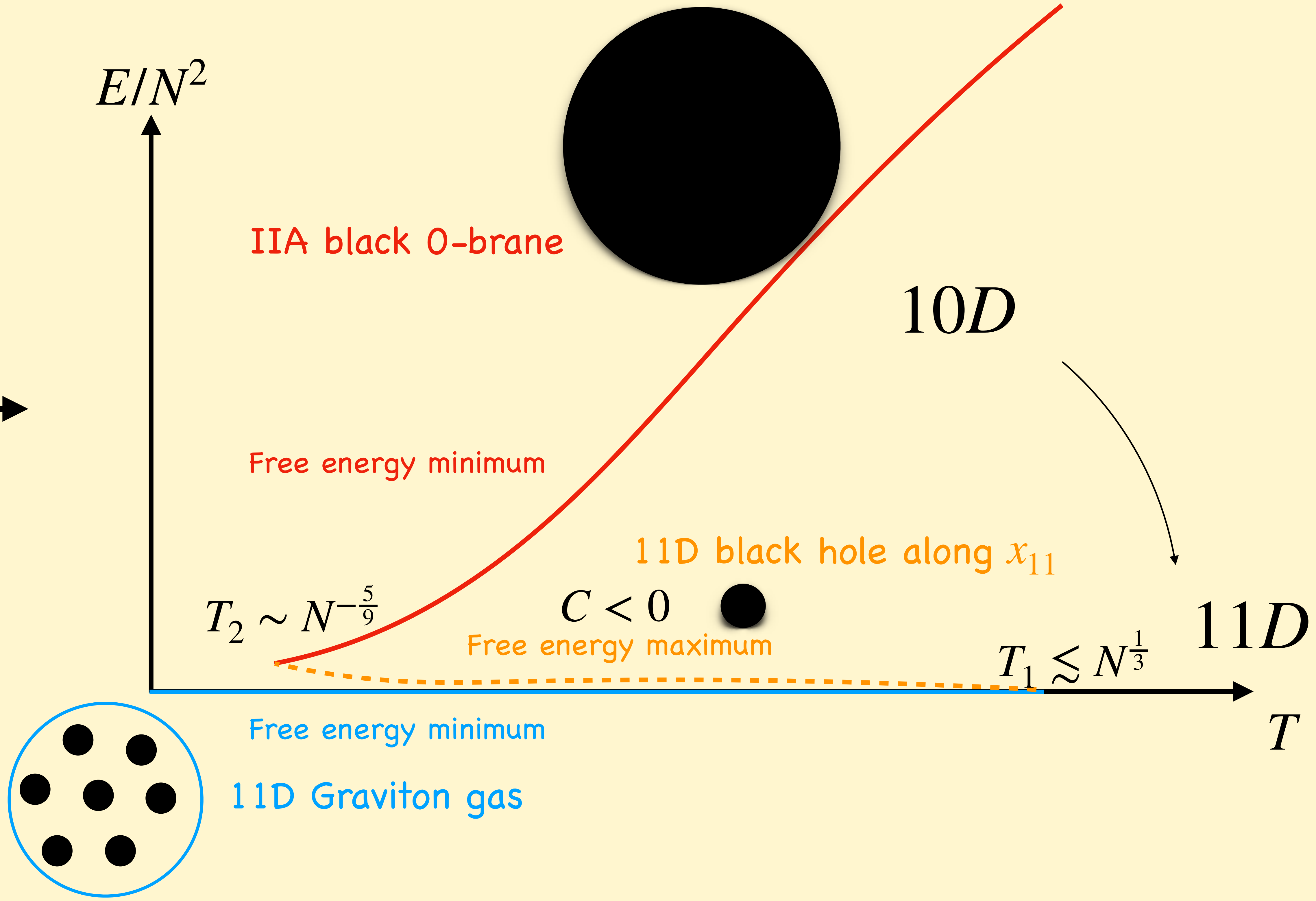
Confinement in the D0-matrix model

How to understand confinement?

Large N

Gravity limit

$$g_{eff} = \frac{1}{E^3}$$



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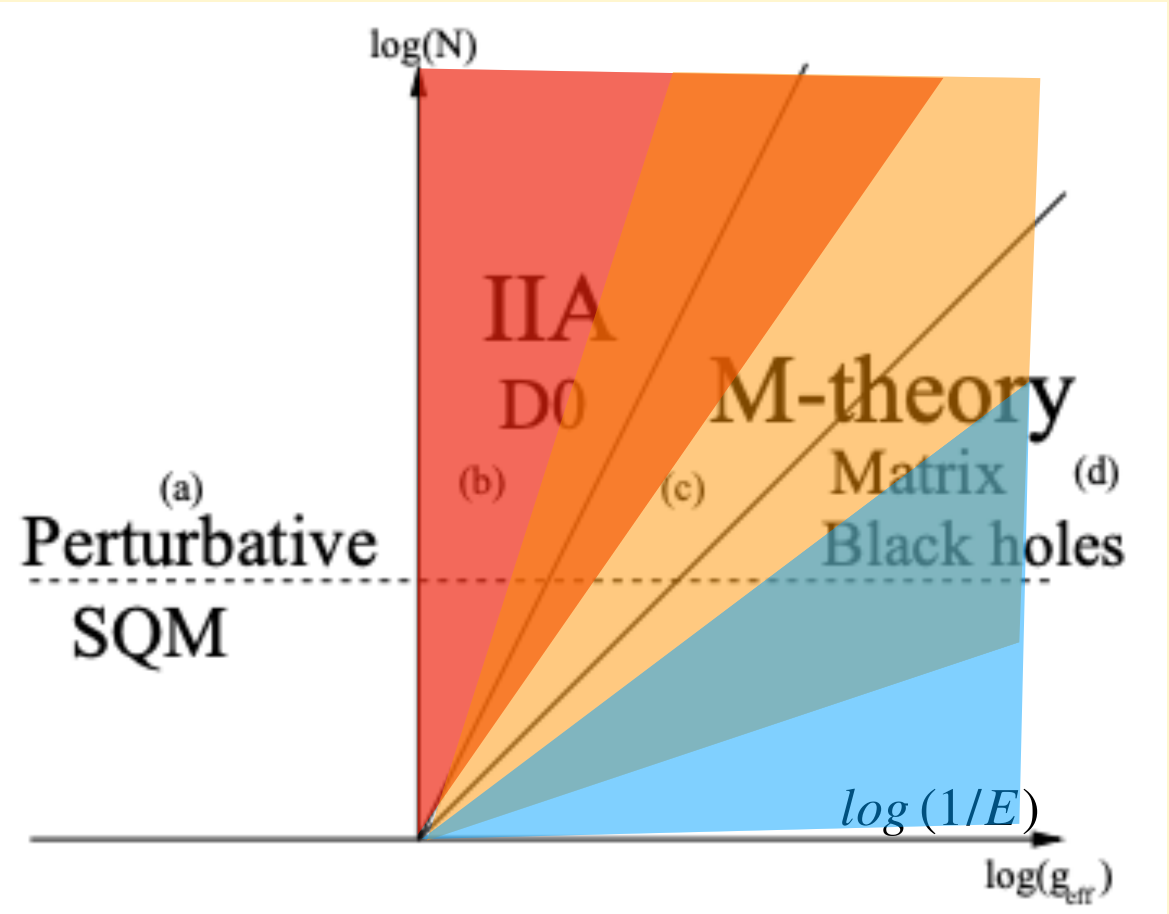

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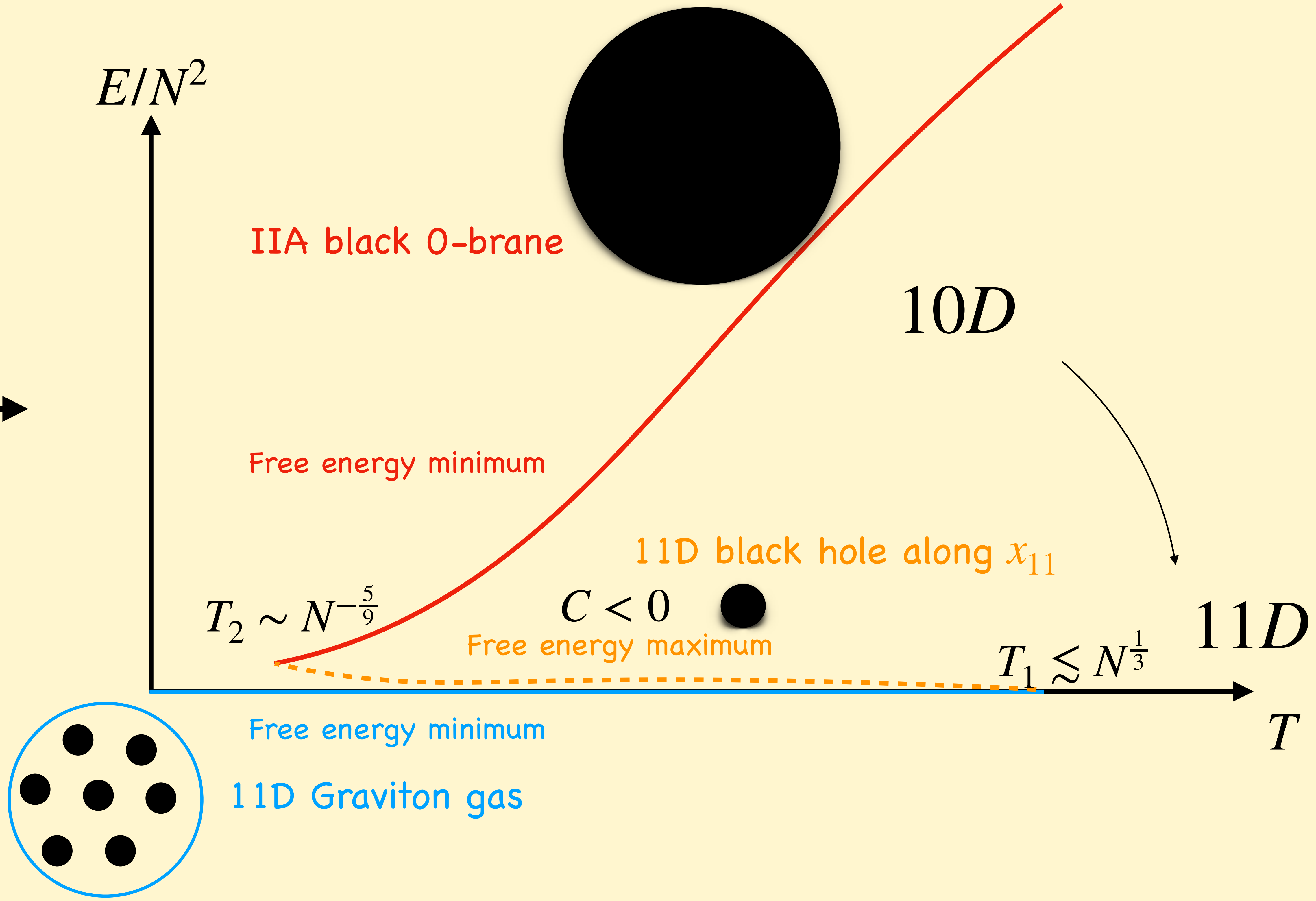

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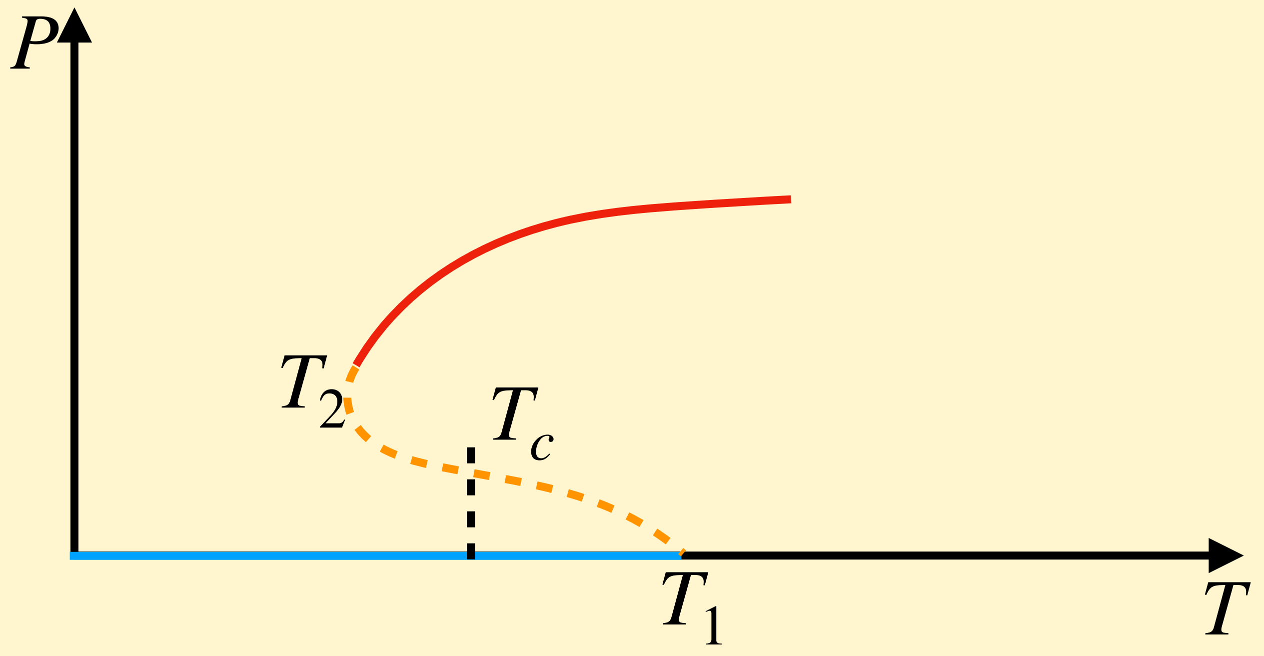
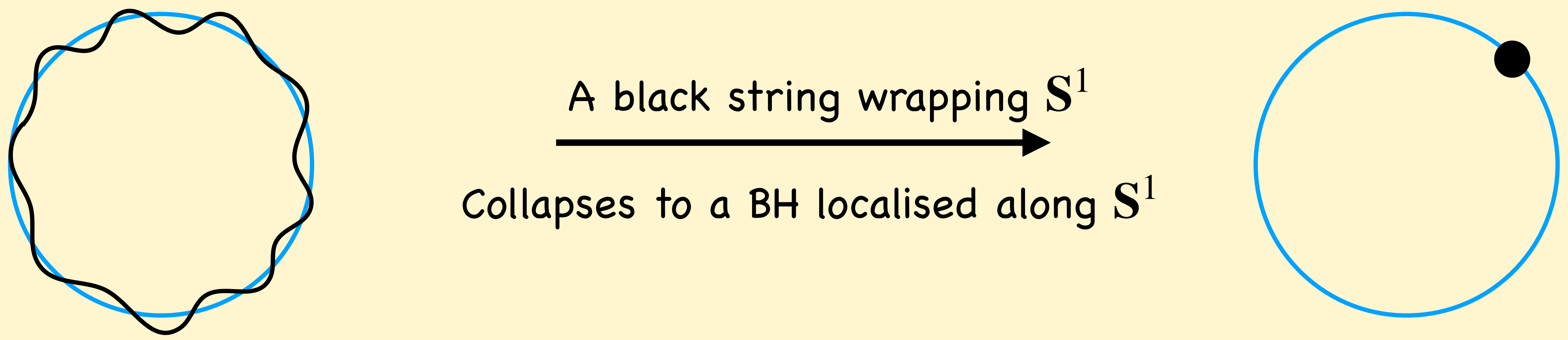


How to determine T_1 and T_2 ?

- T_2 corresponds to Gregory-Laflamme transition (or Gross-Witten-Wadia)

Gregory-Laflamme 1994, Gubser-Mitra 2001, Harmark-Niarchos-Obers 2004, Figueras, Kudoh-Wiseman 2004, Horowitz-Martinec 1997, Gross-Witten, Wadia, ...



$$T_2 \sim N^{-5/9}$$

- T_1 corresponds to maximum/minimum confinement/deconfinement temperature

Schwarzschild BH in 11D with $M = M_{Pl}$ $\longrightarrow$ $T_1 \lesssim N^{1/3}$

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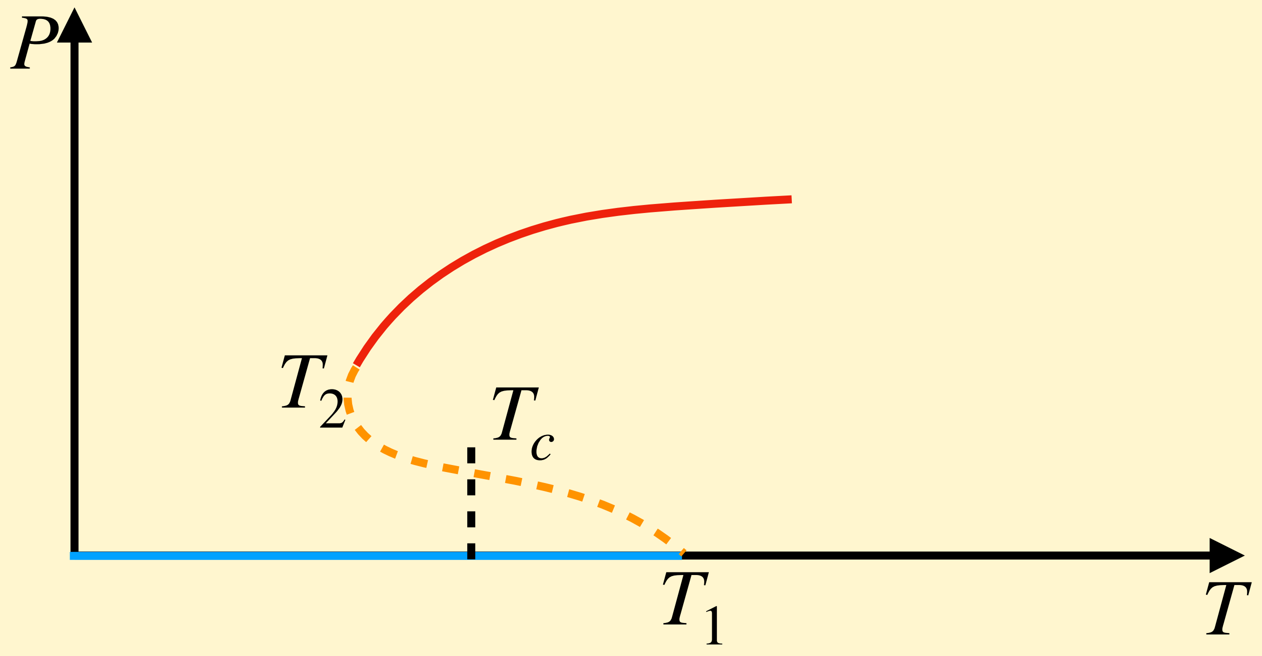

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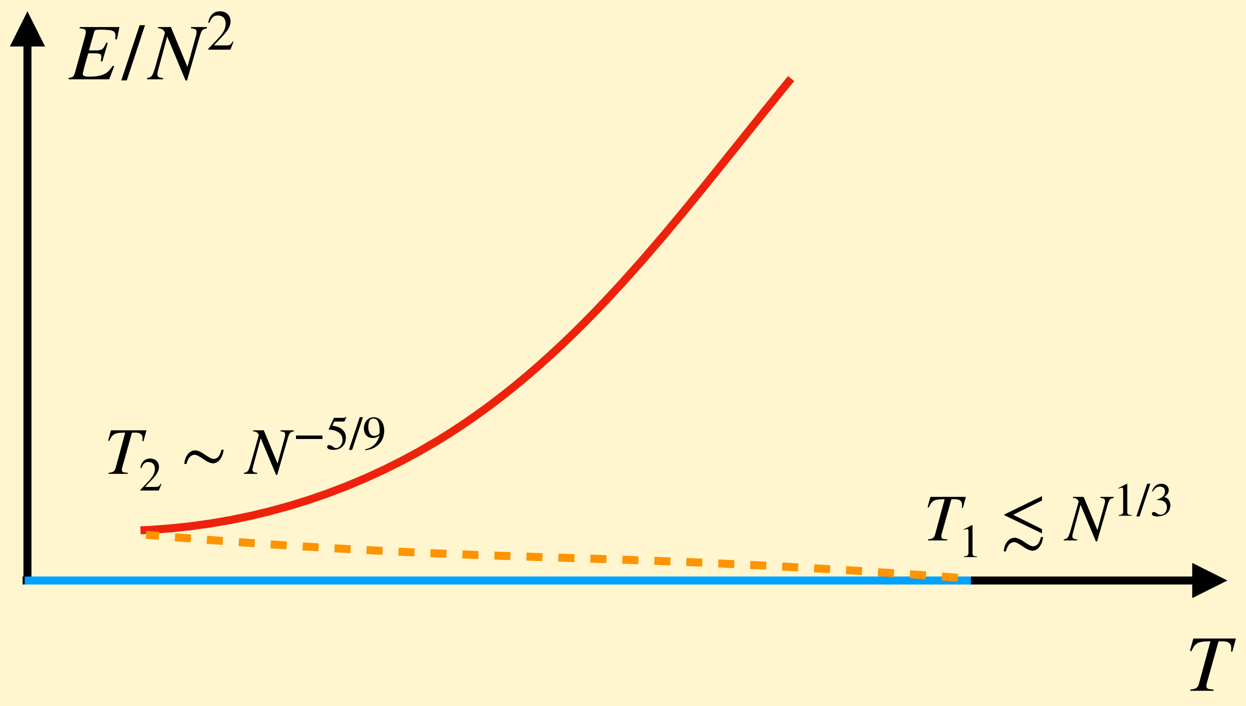

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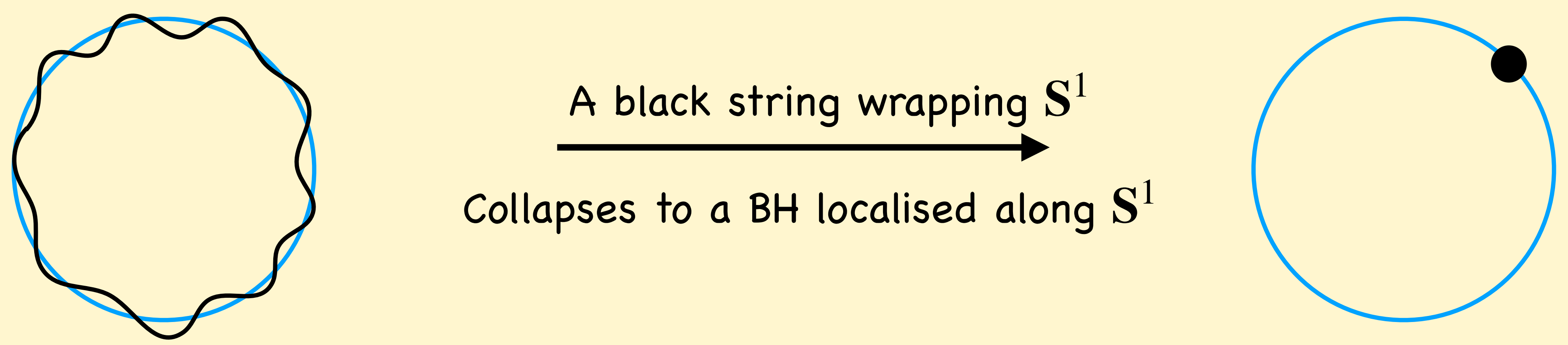

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Deconfined studies in the D0-matrix model

Tests of gauge/gravity duality

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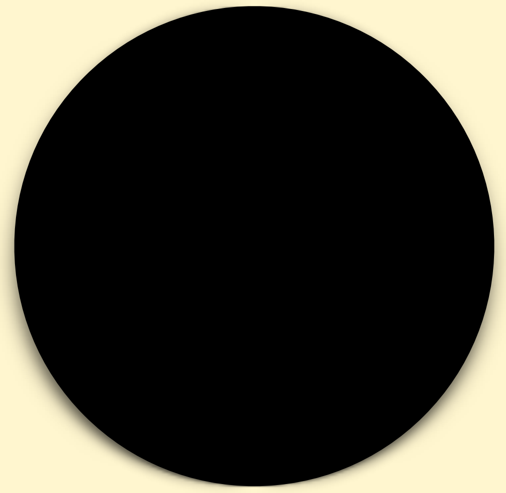

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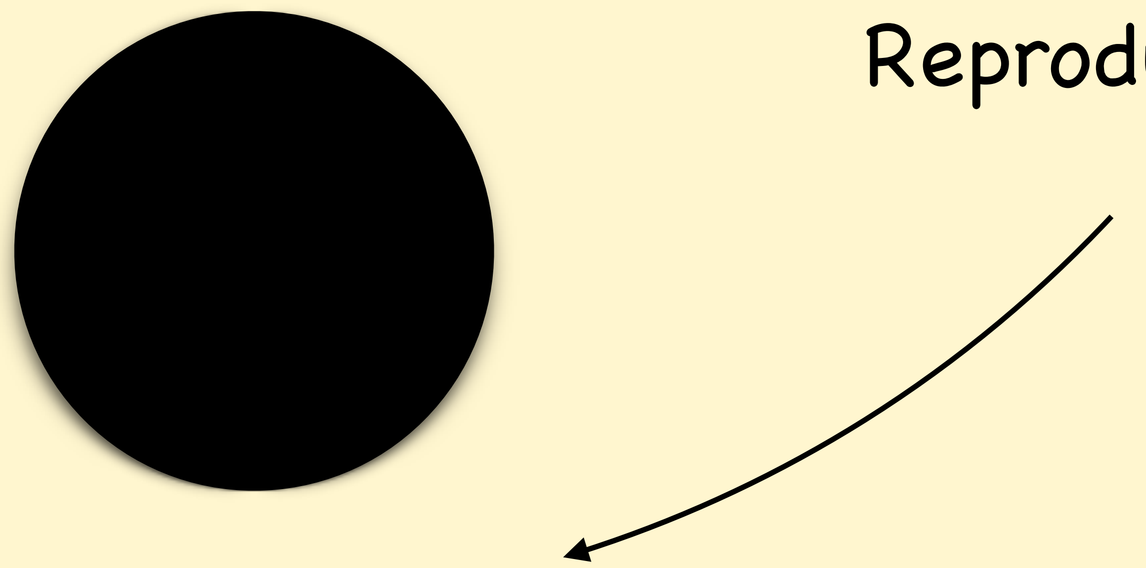

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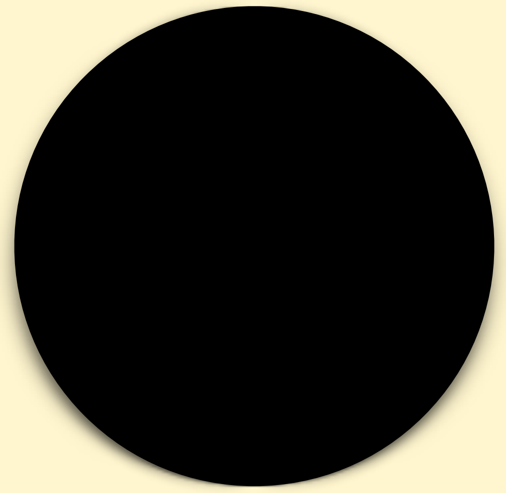

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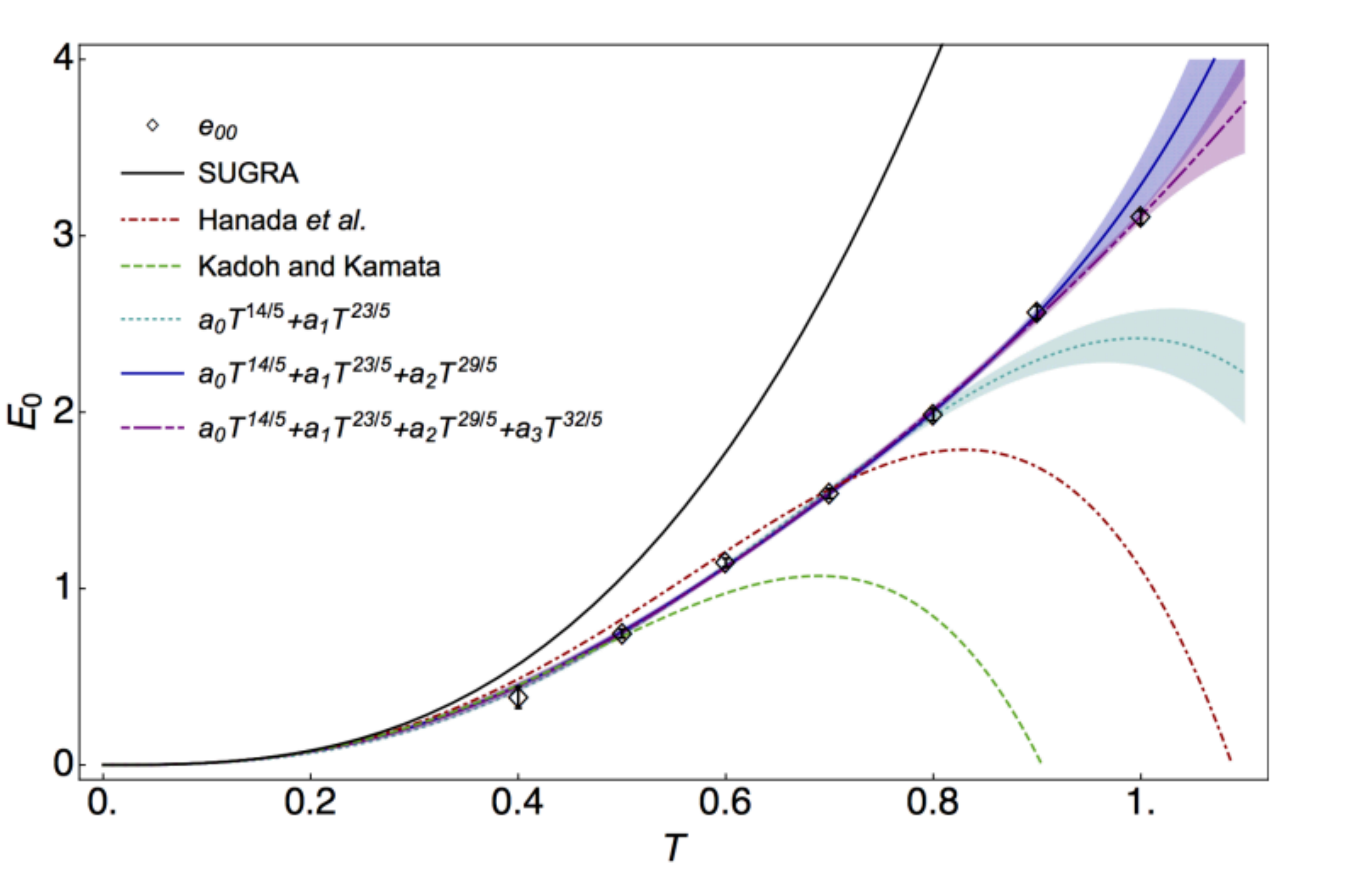

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Finite μ corrections for **BMN supergravity**

$$\frac{E}{N^2} \simeq a_0 T^{\frac{14}{5}} f(\mu)$$

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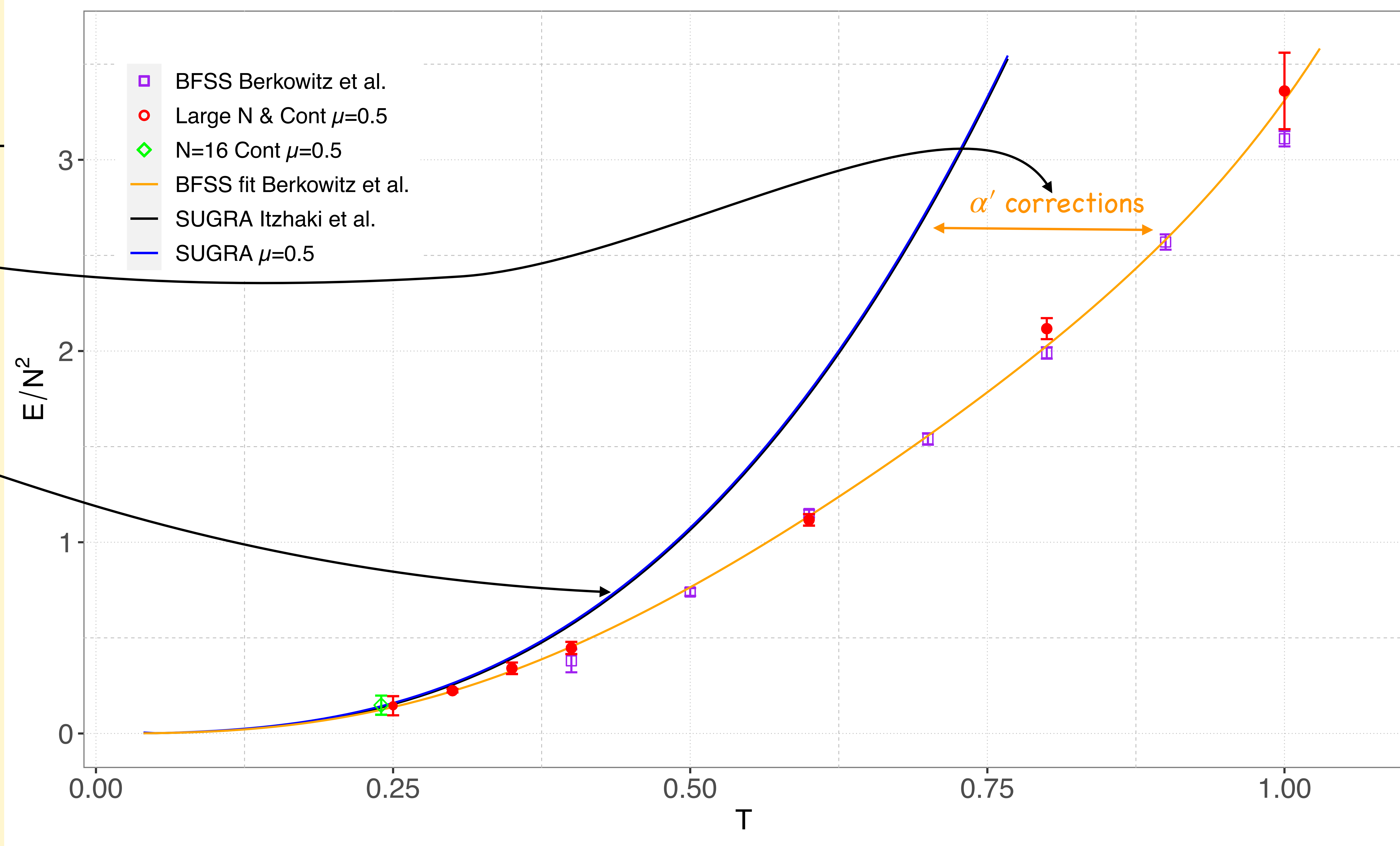
2210.04881

Appeared on 11/10/22



Is the singlet constraint important?

Maldacena-Milekhin conjecture

$$Z_{\text{gauged}} = \int [dX][d\psi][dA_t]e^{-S_{\text{matrix}}[X,\psi,A_t]}$$ → **Gauge** singlet constraint: $\mathcal{G} := \frac{iN}{2\lambda}\left(2[\dot{X}_M, X_M] + [\bar{\psi}_\alpha, \psi_\alpha]\right) = 0$

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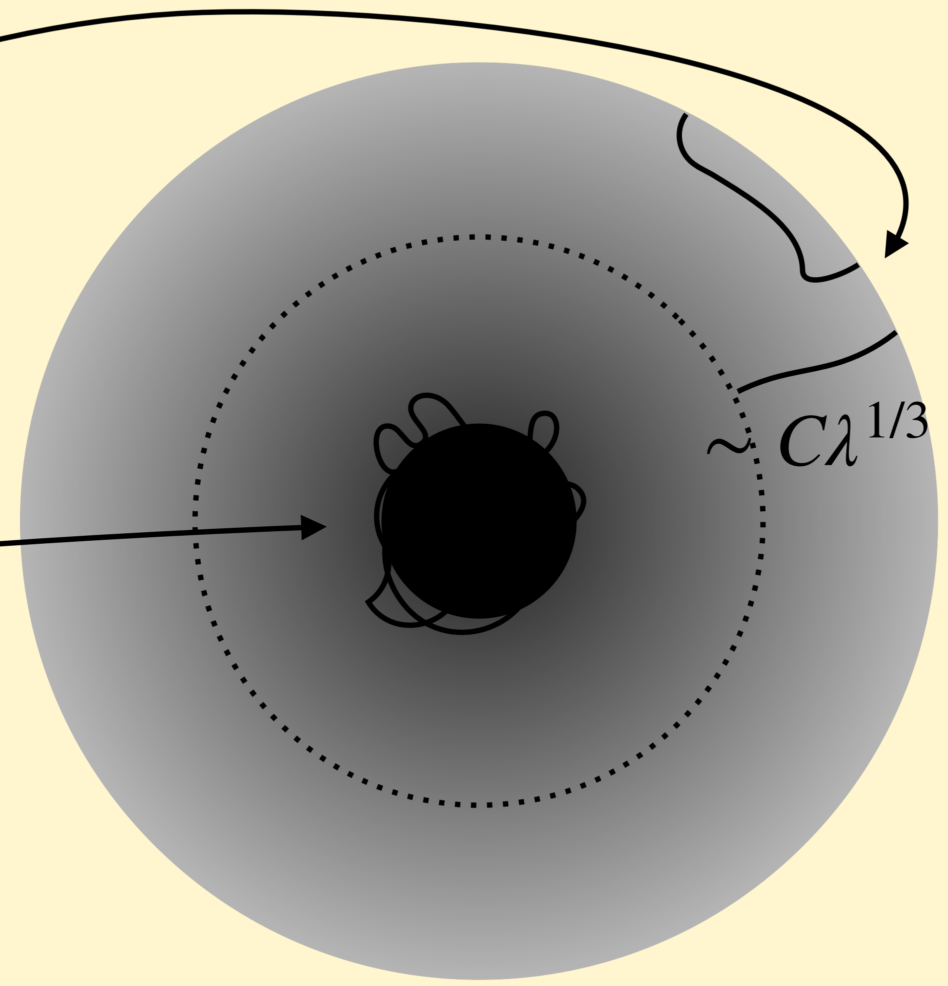
$$\mathscr{H}_{\text{total}} = \mathscr{H}_{\text{singlets}} \otimes \mathscr{H}_{\text{non-singlets}}$$



1802.00428

- Non-singlets do not contribute at low temperatures

$$\Delta Z = Z_{\text{gauged}} - Z_{\text{ungauged}} \simeq e^{-\frac{C\lambda^{1/3}}{T}} \quad , \quad T \to 0$$

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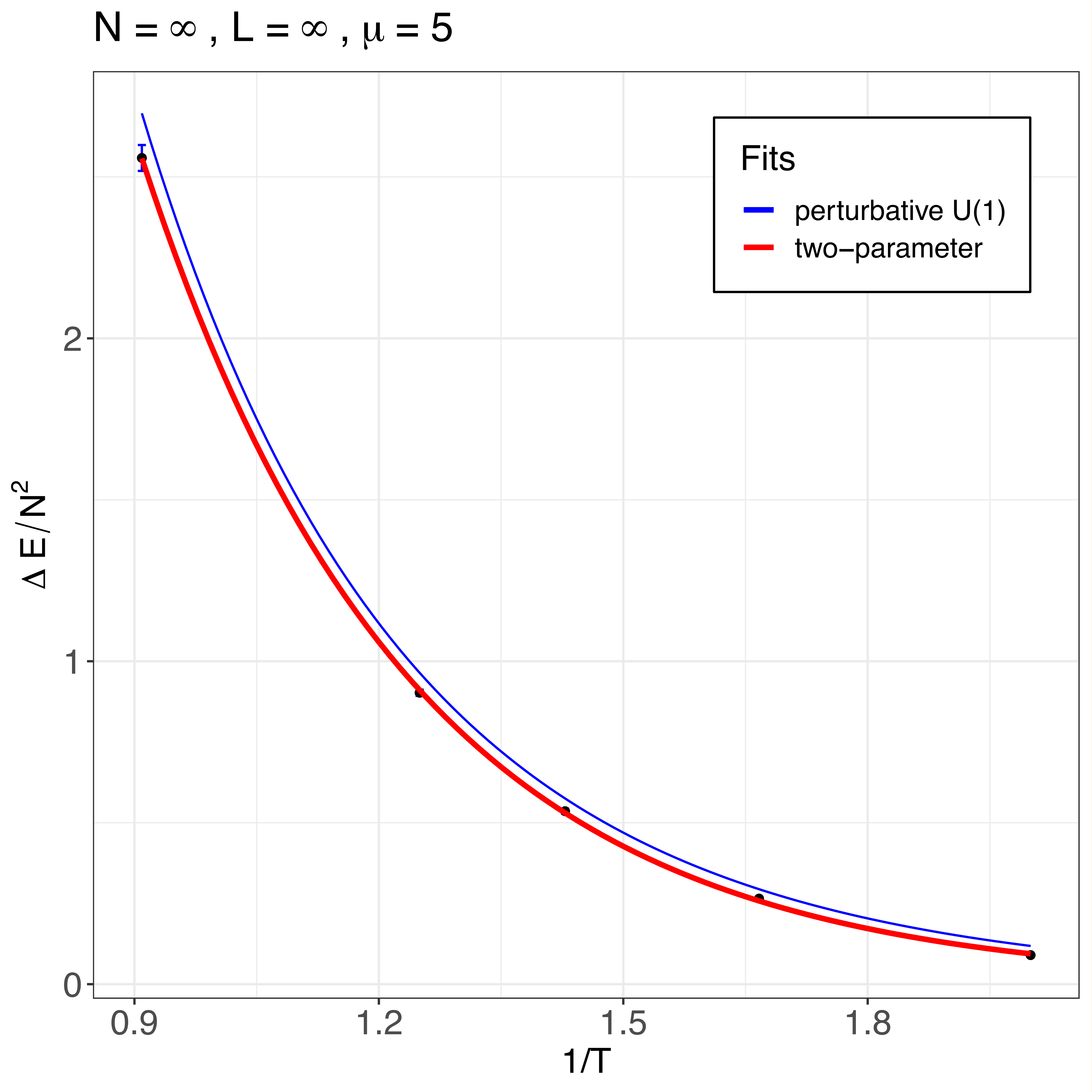
$$\frac{E_{U(1)}}{N^2} = 6 \cdot \frac{\mu}{2} e^{-\frac{\mu}{2}\frac{1}{T}} + 8 \cdot \frac{3\mu}{4} e^{-\frac{3\mu}{4}\frac{1}{T}} + 3 \cdot \mu e^{-\frac{\mu}{T}} \quad , \quad \frac{\lambda}{T^3} \gg 1 \; , \quad \frac{\mu^3}{\lambda} \gg 1$$

$$\frac{\Delta E}{N^2} = E_{\text{ungauged}} - E_{\text{gauged}} = D \cdot e^{-\frac{C}{T}} \quad , \quad \frac{\lambda}{T^3} \gg 1 \quad , \quad \frac{\mu^3}{\lambda} \ll 1$$



What do we learn?

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- A stable confined phase has been observed for the first time
- Interesting possibility to probe contents of M-theory.
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- Low temperature precision test for holography (internal energies)
- Non-AdS/non-CFT, non-gauge/gravity, stringy corrections
- More studies (confined and deconfined) are important
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Thank you